### **Radiation Fundamentals II**

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• Spectral lines of Astronomical interest

• Spectral line formation

• Polarization: application Faraday Rotation

Essential Radio Astronomy, Condon, J. & Ransom, S. https://www.cv.nrao.edu/~sransom/web/xxx.html

### Spectral Lines: HI 21cm line



- •Study of neutral component of ISM -- cold and warm; our Galaxy and external galaxy (cold 60 K, 30 cm<sup>-3</sup>, 5x10<sup>19</sup> cm<sup>-2</sup>; Warm 8000 K, 0.2 cm<sup>-3</sup>, 10<sup>18</sup> cm<sup>-2</sup>)
- •Kinematics of our Galaxy and external galaxy (eg. spiral galaxy)
- •Large scale structure (nearby z<0.16; < 700 Mpc)

### **Spectral Lines: CO line**





Study of cold molecular gas in ISM; (10 K, 10<sup>3</sup> cm<sup>-3</sup>, 20 pc)

Proxy for  $H_2$ ; cloud properties (need to know  $X_{co}$ ).

GBO/AO single dish workshop, Green Bank, August 19, 2019

<sup>12</sup>CO, <sup>13</sup>CO, C<sup>18</sup>O Image of Orion region

#### Spectral Lines: Radio Recombination Lines



FIG. 3. The theoretically computed radio spectra for the Orion Nebula and NGC 2024, together with the observed flux densities.

#### Spectral Lines: Astronomical Masers





Stimulated microwave spectral line emission

OH, H<sub>2</sub>O, CH<sub>3</sub>OH, SiO, HRRL

Structure of Milky Way



Plan view of Milky Way. Dots are the locations of newly formed OB-type stars determined from VLBI trigonometric parallaxes using associated  $H_2O$  or methanol maser emission. The Galactic center is denoted by the plus (+) sign at (0,0) kpc; the Sun is labeled in yellow at (0,8.4) kpc.



Einstein A coefficient (spontaneous emission coefficient)

Classical approximation: A<sub>ul</sub> is average spectral power emitted by a dipole divided by the photon energy.

HI 21cm line: 
$$A_{F=10} = 2.85 \times 10^{-15} \text{ s}^{-1} (\sim 11 \text{ Myr half-life})$$

<sup>12</sup>C<sup>16</sup>O rotational trans:  $A_{J=10} = 7.2 \times 10^{-8} \text{ s}^{-1} (\sim 0.4 \text{ yr}; \propto \text{J}^3)$ 

RRL transitions: 
$$A_{n+1,n} = 5.3 \times 10^9 \left(\frac{1}{n^5}\right)$$
 s<sup>-1</sup> (n=100; ~ 2 sec)

Energy level population

$$\frac{N_u}{N_l} = \frac{g_u}{g_l} e^{-\frac{E_{ul}}{k_b T}}$$
 T is a unique temperature in TD equilibrium  
Otherwise T=T<sub>ex</sub> (excitation temperature)

#### Statistical Equilibrium

For two level system  

$$N_u A_{ul} + N_u C_{ul} + N_u B_{ul} \frac{I_{\nu}}{c} = N_l C_{lu} + N_l B_{lu} \frac{I_{\nu}}{c}$$

. .

Multi-level system

$$N_i \sum_j R_{ij} = \sum_j N_j R_{ji}$$

Solution to this equation gives the level population;  $T_{ex}$  can be different for different levels

#### Energy level population: LTE approximation

$$\begin{array}{l} \displaystyle \frac{N_u}{N_l} &= \displaystyle \frac{g_u}{g_l} \ e^{-\frac{E_{ul}}{k_b T}} & T = T_k \ \text{the kinetic temperature} \\ & \text{(In general for LTE T=T_{ex} is same of all level;} \\ \displaystyle C_{ul} \ \approx \ A_{ul} \end{array}$$

 $n^* \sigma v \approx A_{ul}$   $n^*$  critical density (density of collision partner);  $\sigma$  cross section; v mean velocity

HI 21 cm line: n<sup>\*</sup> << 1 cm<sup>-3</sup>

Condition:

<sup>12</sup>C<sup>16</sup>O J=1-0 line:  $n^* \sim 10^3 \text{ cm}^{-3}$  (T<sub>k</sub> ~ 20 K)

**RRLs**: 
$$n^* \sim 90 \text{ cm}^{-3} (T_k \sim 10^4 \text{ K}; n=100)$$

Energy level population: LTE approximation

$$\frac{N_u}{N_l} = \frac{g_u}{g_l} e^{-\frac{E_{ul}}{k_b T}} \quad T = T_k \text{ same for all levels}$$
$$\frac{N_u}{n_d} = \frac{g_u e^{-\frac{E_u}{k_b T_k}}}{Q}$$
$$Q = \sum_i g_i e^{-\frac{E_i}{k_b T_k}} \quad \text{Partition function}$$

HI 21cm line

$$Q \approx g_0 + g_1 e^{-\frac{h\nu_{10}}{k_b T_k}} - \frac{h\nu_{10}}{k_b} \sim 0.08 \text{ K}$$

 $\approx 1+3=4$  (This is true for other values of T encountered in the ISM)

$$N_u \approx \frac{3}{4}n_d$$

Energy level population: LTE approximation

<sup>12</sup>C<sup>16</sup>O J=1-0 line

$$Q_{rot} \approx \frac{2k_b T_k}{h\nu_{10}} = 0.36T_k$$

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 $N_u \approx n_d \times \frac{3 \ e^{-\frac{h\nu_{10}}{k_b T_k}}}{0.36 T_k}$  for  $\nu_{J=10} = 115.27 \ \text{GHz transition}$ 

(see Mangum & Shirley 2015; Turner 1991)

•Rotation – lowest energy state

•Q has contribution from rotation, vibration, ...

#### Hydrogen RRLs

Number density of atoms in quantum state n

$$N_n \propto n^2 \frac{n_p n_e}{T_e^{3/2}}$$
$$\propto n^2 \frac{n_e^2}{T_e^{3/2}} \quad n_d \approx n_e$$

- Saha-Boltzmann equation in Q
- •Q diverges for hydrogenic atoms
- •For fulling ionized gas in ISM

 $n_{d} \sim n_{p}$  the proton density  $\sim n_{e}$ 

 $T_{k} = T_{e}$  the electron temperature

#### Optically thin case; no background radiation



What we want?  $n_d (cm^{-3}); N \sim n_d 2R_s (cm^{-2})$ T (K), R<sub>s</sub> (pc)

Prediction for obs

 $S_{L}(\nu) = \frac{P_{\nu}}{4\pi R_{d}^{2}} \propto N_{u} A_{ul} 2R_{s} \left(\frac{R_{s}}{R_{d}}\right)^{2} \phi(\nu) h\nu_{0} \quad \text{ergs/s/cm}^{2}/\text{Hz} \quad \text{Flux density of the line}$   $\propto N_{u} 2R_{s} A_{ul} \phi(\nu) h\nu_{0} \Omega_{s} \quad \text{ergs/s/cm}^{2}/\text{Hz}$   $N_{u} \text{ in terms of the total density} \quad \text{Normalized Gaussian}$ 

Not optically thick; background radiation





$$I_{\nu}(\tau_{\nu}) = I_{\nu}(0) e^{-\tau_{\nu}} + S_{\nu}(1 - e^{-\tau_{\nu}})$$

Line flux density

$$S_L(\nu) = S_{\nu} \tau_{\nu} \Omega_s$$
 For  $I_{\nu}(0) = 0; \tau_{\nu} << 1$ 

 $= u_{\nu} 2R_s \Omega_s$  ergs/s/Hz/cm<sup>2</sup>

Source function

$$S_{\nu} = B(T) \approx \frac{2k_bT}{\lambda^2}$$

In TD equilibrium by Kirchhoff's law

LTE approximation

$$\frac{N_u}{N_l} = \frac{g_u}{g_l} e^{-\frac{E_{ul}}{k_b T}} \quad T = T_{ex}$$

$$S_{\nu} \approx B(T_k) \approx \frac{2k_b T_k}{\lambda^2} \quad T_{ex} = T_k \text{ same for all levels}$$

$$S_{\nu} \approx B(T_{ex}) \approx \frac{2k_b T_{ex}}{\lambda^2} \quad T_{ex} \text{ same for all levels}$$

Solution to RT in temperature

$$I_{\nu}(\tau_{\nu}) = I_{\nu}(0) e^{-\tau_{\nu}} + S_{\nu}(1 - e^{-\tau_{\nu}})$$

$$S_{\nu} \approx B(T_{ex}) \approx \frac{2k_b T_{ex}}{\lambda^2} \quad T_{ex} \text{ same for all levels}$$

$$I_{\nu}(0) = \frac{2k_b T_{\nu,b}}{\lambda^2} \quad T_{\nu,b} \text{ brightness temperature}$$

$$T_{\nu}(\tau_{\nu}) = T_{\nu,b} e^{-\tau_{\nu}} + T_{ex}(1 - e^{-\tau_{\nu}}) \qquad \mathsf{K}$$

Line optical depth:LTE case



Peak line optical depth: LTE case

$$au_{
u} \propto N_u \, 2R_s \, A_{ul} \, rac{1}{
u_0 T_{ex}} \, \phi(
u)$$



HI 21cm line

 $N_u = \frac{3}{4} n_d; A_{ul} \text{ is const}$ Freq = 1420.405 MHz

$$\tau_{\nu}(\nu_0) = 5.1335 \times 10^{-19} \frac{N_H (\text{cm}^{-2})}{T_s(\text{K})\Delta v (\text{km s}^{-1})}$$

 $N_H$  column density of HI,  $T_{ex} = T_s$  spin temperature

Peak line optical depth: LTE case

$$au_{
u} \propto N_u \, 2R_s \, A_{ul} \, rac{1}{
u_0 T_{ex}} \, \phi(
u)$$



CO J=1-0 transition

$$\tau_{\nu_{1,0}}(\nu_0) = 3.95 \times 10^{-15} \left( 1 - e^{-\frac{h\nu_{1,0}}{k_b T_{ex}}} \right) \frac{N_{CO}(\text{cm}^{-2})}{T_{ex}(\text{K}) \ \Delta v(\text{km s}^{-1})}$$

 $N_{CO}$  column density of CO,  $T_{ex}$  excitation temperature (did not use hv/k<sub>b</sub>T<sub>ex</sub> << 1 approximation; T<sub>ex</sub> ~ 20 K)

Peak line optical depth: LTE case

$$\tau_{\nu} \propto N_u \, 2R_s \, A_{ul} \, \frac{1}{\nu_0 T_{ex}} \, \phi(\nu)$$



#### **RRL** alpha transition

$$\tau_{\nu}(\nu_{0}) = 1.92 \times 10^{3} \frac{EM(\text{pc cm}^{-6})}{T_{e}^{5/2}(\text{K})\Delta\nu(\text{KHz})}$$
$$EM = \int n_{e}^{2} \text{d}s \qquad \text{Emission measure}$$

For ionized gas  $T_{ex}=T_{k}=T_{e}$  $N_{u} \alpha n^{2} 1/T_{e}^{3/2} n_{e}^{2}$  $A_{ul} \alpha 1/n^{5}$ RRL alpha freq  $\alpha 1/n^{3}$ 

Knowing source function and optical depth we can get  $S_L$  from RT solution

Excitation temperature T<sub>ex</sub>

Consider two level system

$$\begin{split} N_{u}A_{ul} + N_{u}C_{ul} + N_{u}B_{ul}\frac{I_{\nu}}{c} &= N_{l}C_{lu} + N_{l}B_{lu}\frac{I_{\nu}}{c} & \text{Statistical equilibirum} \\ & \\ \frac{N_{u}}{N_{l}} = \frac{g_{u}}{g_{l}} e^{-\frac{E_{ul}}{k_{b}T_{ex}}} &= \frac{C_{lu} + B_{lu}\frac{I_{\nu}}{c}}{A_{ul} + C_{ul} + B_{ul}\frac{I_{\nu}}{c}} \\ \text{Radiation dominates} & \text{Collision dominates} \\ & \\ \frac{B_{lu}\frac{I_{\nu}}{c}}{A_{ul} + B_{ul}\frac{I_{\nu}}{c}} & \\ & \\ T_{ex} \to T_{b} & \\ \end{array}$$

Non-LTE effects

$$\frac{N_u}{N_l} = \frac{g_u}{g_l} e^{-\frac{E_{ul}}{k_b T_{ex}}} \qquad \mathsf{T}_{ex} \text{ different for different level}$$

$$= \frac{b_u g_u}{b_l g_l} e^{-\frac{E_{ul}}{k_b T_k}} \qquad \mathsf{LTE level population with T=T}_k$$

$$\mathsf{b}_u, \mathsf{b}_l \text{ are departure coefficient}$$

Solve for b<sub>1</sub>, b<sub>1</sub> using statistical equilibrium equation.

#### Maser emission

 $N_u > N_l$  Upper level population greater than lower level

$$\frac{N_u}{N_l} = \frac{g_u}{g_l} e^{\frac{E_{ul}}{k_b T_{ex}}} \quad \mathsf{T}_{\mathsf{ex}} \text{ is negative } !!$$

 $I_{\nu}(\tau_{\nu}) \approx I_{\nu}(0) e^{\tau_{\nu}}$  Optical depth is negative  $\rightarrow$  exponential amp.

#### Polarization of galactic background emission



WMAP image of polarized galactic synchrotron emission at 3° scale

DRAO image of galactic polarized synchrotron Emsission smoothed to 3° scale. The difference in emission structure is dominated by Faraday depolarization.

#### **Measurement**



General elliptical pol from two (orthogonal) linear pol

$$\vec{E} = \left( E_x \ e^{j\phi_x} \ \hat{x} + E_y \ e^{j\phi_y} \ \hat{y} \right) e^{j\omega t}$$

$$E_x, E_y, \delta = \phi_x - \phi_y$$

p, fractional pol



Stokes parameters (incoherent source)

 $I = \langle E_x^2 \rangle + \langle E_y^2 \rangle$ 

 $Q \; = \; < E_x^2 > - < E_y^2 >$ 

$$U = 2 \operatorname{Re} \langle E_x E_y^* \rangle = 2 \langle E_x E_y \cos(\delta) \rangle$$

$$V = 2 \operatorname{Im} \langle E_x E_y^* \rangle = 2 \langle E_x E_y \sin(\delta) \rangle$$

**Measurement** 

Stokes parameters (incoherent source)

$$I = \langle E_x^2 \rangle + \langle E_y^2 \rangle$$

$$Q = \langle E_x^2 \rangle - \langle E_y^2 \rangle$$

$$U = \langle E_{45}^2 \rangle - \langle E_{-45}^2 \rangle$$

$$V = < E_l^2 > - < E_r^2 >$$



Receiving sense Hamaker & Bregman (1996)



Angle of linear polarization

 $P = Q + iU = |P| e^{2i\chi}$  Linear polarization in complex notation; |P| lin. pol flux density

#### **Measurement**



General elliptical pol from two (orthogonal) circular pol





•Linear pol light  $\rightarrow$  sum of right and left circular pol

•Right and left circular pol velocities are different (due to coupling with the cyclotron motion of thermal electrons)

$$\begin{split} 2\Delta\psi &= \left(\frac{2\pi}{\lambda_r} - \frac{2\pi}{\lambda_l}\right) \Delta s & \Delta \Psi \text{ is pol rotation angle} \\ \Delta\psi &\approx \lambda^2(\text{m}^2) \times 0.8125 \int_{obs}^{source} B_{||}(\mu\text{G})n_e(\text{cm}^{-3})\text{d}s(\text{pc}) & \text{n}_e \sim 0.1 \text{ cm}^{-3} \text{ f}_p \sim 2.8 \text{ KHz} \\ & \text{For } \omega >> \omega_p \& \omega >> \omega_c \end{split}$$

 $pprox \lambda^2(m^2) RM \ (rad \ m^{-2})$  RM is the rotation measure

#### RM measurement



 $\chi = \frac{1}{2} \tan^{-1} \left( \frac{U}{Q} \right)$  Angle of the measured linear pol

 $P = Q + iU = |P| e^{2i\chi}$  |P| is the flux density of linear polarization

$$\Delta \psi \approx \lambda^2 (m^2) RM (rad m^{-2})$$
 Diff between the incident and emerged pol angle

Combining RM and DM one could get  $B_{\parallel}$  field (Haverkorn et al 2015)

Faraday tomography or RM synthesis

Burn (1966) Brentjens & de Bruyn (2005)



 $\phi(r) = 0.8125 \int_{obs}^{r} B_{||} n_e \mathrm{d}s$  Faraday depth (not RM; RM is the integral over the whole source)

 $\Delta\psi(r)~=~\lambda^2\phi(r)~$  Change in pol angle due to Faraday depth

$$P_{obs}(\lambda^2) = \int P(r) \ e^{2i\phi(r)\lambda^2} d\phi(r)$$
Fourier transform ( $\Phi$ ,  $\lambda^2$ )  

$$A$$

$$Q_{obs} + iU_{obs}$$

$$Q(r) + iU(r) \text{ linear pol at r average over the source}$$

$$GBO/AO \text{ single dish workshop, Green Bank, August 19, 2019}$$

Thank You