



Observing Techniques and Calibration 2

David Frayer (thanks to Ron Maddalena, Jim Condon, and Ryan Shannon for various slides)

Topics for Talk 1 (Larry)

- Blackbody radiation
- Basic radiative transfer
- Spectral Line Data Calibration
- Calibration of Ta via noise diodes and definition of Tsys: Ta= Tsys*(ON-OFF)/OFF
- ON-Ref observing techniques: PS, FSW, BSW, NOD, SubBeamNod
- Mapping Methods: RALong/DecLat and Daisy OTF, Point-Map
- Nyquist Sampling



Topics for Talk 2 (David)

- Continuum Observations
- 1/f noise and confusion
- Pulsar Observations
- Calibration into physical units Jy and K
- Definition of the various telescope temperature scales and efficiencies
- The effects of weather on data at highfrequency





The spectrum of the universe







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very isotropic, and hence very distant





"Confusion"

Profile plot of 45 deg² near the NGP imaged with 12' resolution at 1.4 GHz. Strongest source S \approx 1.5 Jy.







5σ confusion source detection limits at GBT and Arecibo





Continuum Radiometer Equation

$$\sigma_{Ta} = T_{SYS} \sqrt{\frac{1}{\Delta \nu \cdot t} + \left(\frac{\Delta G}{G}\right)^2}$$

Note (need to include confusion noise floor): (sigma_Ta_total)^2 = (sigma_ta)^2 + (sigma_confusion)^2

- Receiver delta-G/G gain stability results in "1/f" noise.
- Baseline subtraction removes the 1/f noise for spectral-line observations.
- The 1/f noise dominates for single-dish continuum observations, except for receivers built with fast electronic beam-switching (e.g., CCB for Ka-band on the GBT).
- Need multiple scans over the source for continuum observations to decrease the noise (increasing the bandwidth does not help for single-dish continuum observations, unlike interferometer continuum observations where the gain instability does not correlate between antennae).





Electronic 1/f noise



Example power spectrum for GBT 340 GHz receiver.

- The "white" noise level is ~10^-5 in power spectrum for frequencies >4Hz (time < 0.25 s).
- □ The 1/f noise dominates for time >1 sec.
- ❑ At 100sec the 1/f noise is ~1000x higher than the white noise.
- ❑ Would need to slew across point-sources at ~6deg/s to remove the 1/f noise for this Rx on the GBT
 and the GBT cannot slew nearly this fast.... → would need a new Rx to measure continuum (which is actually under construction)





Effective Continuum Band-width limits for the GBT (GBT Memo#282 Mason 2013)

Receiver	$\frac{\Delta G}{G}(5\mathrm{Hz})$	$\Delta \nu_{eff.}$
	2	[MHz]
PF(300 MHz)	$\sim 2.3 imes 10^{-3}$ *	2
PF(850 MHz)	$\sim 1.2 \times 10^{-3}$ *	7
L	$(3.7 \pm 2.4) \times 10^{-4}$	73
\mathbf{S}	$(7.9 \pm 2.8) \times 10^{-4} *$	16
\mathbf{C}	$(1.5 \pm 0.8) \times 10^{-3}$	4
Х	$(6.9 \pm 1.5) \times 10^{-4}$	21
Ku	$(1.6 \pm 1.2) \times 10^{-4}$	390
Κ	$(5.2 \pm 2.0) \times 10^{-4}$	37
Ka	$(5.0 \pm 2.1) \times 10^{-4}$	40
Q	$(4.1 \pm 1.1) \times 10^{-4}$	59
W	$(3.6 \pm 1.3) \times 10^{-4}$	77

Do not blindly assume the 1.5 GHz VEGAS Mode-1 bandwidth for continuum noise estimates. Use these tabulated bandwidths for sensitivity estimates based on Rx.

Example:

Ka-band and VEGAS/DCR

- limited to 40MHz. CCB
- (12GHz) with fast

electronic switching is about

(12000MHz/40MHz)^0.5 =

17x more sensitive based

on effective bandwidth

considerations....





Pulsar Observations





Dynamic Spectra: Pulsar Radiation as a Function of Frequency and Time



Time

It is necessary to remove the effect of dispersion 1) Increase SNR to detect pulsar; and 2) Best study pulsar

CSIRO

Ryan Shannon, Pulsar Observations @ Parkes Radio School

Dispersed Pulse Coherently dedispersed pulse





Calibration





Background Material:

$$P(\nu, \theta, \varphi, t) =$$

$$G_{RF}(\nu, t)G_{IF}(\nu, t) \cdot$$

$$[T_{Rc\nu r}(\nu, t) + T_{A}(\nu, \theta, \varphi, t)$$

$$\eta_{l}(E) \left(T_{CMB}e^{-\tau(\nu, t)A(E)} + T_{ATM}(\nu, t) \cdot \left(1 - e^{-\tau(\nu, t)A(E)}\right)\right) +$$

$$\left(1 - \eta_{l}(E)\right) \cdot T_{Spill}(E)]$$

 $T_{SYS}(v, \theta, \varphi, t) = [Everything between the square brackets]$

$$T_A(\nu,\theta,\varphi,t) = \frac{\eta_R e^{-\tau(\nu,t)A(E)}}{\iint_{4\pi} f d\Omega} \iint_{4\pi} f(\theta - \theta',\varphi - \varphi') T_B(\theta',\varphi') \ d\theta' \ d\varphi'$$

$$S = \frac{2k}{\lambda^2} \iint_{\Omega_{SRC}} T_B d\Omega \qquad \qquad S = \frac{2kT_A e^{\tau(\nu,t)A(E)}}{\eta_A A_G}$$
(Point Source Only)

$$P_{Off}(\nu) = G_{RF}(\nu)G_{IF}(\nu) \cdot [T_{Rc\nu r}(\nu) + etc.]$$

$$P_{On}(\nu) = G_{RF}(\nu)G_{IF}(\nu) \cdot [T_{Rc\nu r}(\nu) + T_{Diode}(\nu) + etc.]$$

$$G_{RF}(\nu)G_{IF}(\nu) = \frac{P_{On}(\nu) - P_{Off}(\nu)}{T_{Diode}(\nu)}$$

Slight differences in calibration methods for continuum and spectral-line data:

Continuum:

$$T_{A} = T_{Diode} \frac{P_{SIG} - P_{REF}}{P_{On} - P_{Off}}$$

Spectral Line:

$$T_A(\nu) = \frac{P_{SIG}(\nu) - P_{REF}(\nu)}{P_{REF}(\nu)} \left\langle \frac{T_{Diode}(\nu)P_{REF}(\nu)}{P_{On}(\nu) - P_{Off}(\nu)} \right\rangle_{BW}$$

Continuum - Point Sources On-Off Observing



Daisy Scan Map: maximizes the number of passes over small source for continuum observations (i.e. avoid overheads of turning the telescope around as done in RA/Dec mapping)



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Calibration

Noise Diodes



All GBT receivers besides 4mm, Argus, and Mustang use noise diodes.







Calibration

Hot & Cold Loads



Gain: g =(Thot – Tcold)/(Vhot –Vcold) [K/Volts]

Tsys = g Voff

Example GBT 4mm Rx





Absolute Calibration on known astronomical sources (point sources)

→ Corrects for any errors in the adopted Tdiode/gains measured in the lab and corrects for the telescope response

Observe and process target source and known calibrator (3cX) in the same way, then the flux density of the source S(source) is simply:

S(source)/S(3cX) = T(source)/T(3cX),

where S(3cX) is known and T(source) and T(3cX) are observed.

Absolute calibration typically known to 5-15%





Absolute Calibration tied to Mars via WMAP

VLA calibration (0.05-50 GHz):

Perley & Butler 2013, 2017

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VLA Stable Calibrators







Do not blindly assume the GBT Noise Diode Calibration

- Noise diodes are recommended to be sent back for recalibration every 6 months to meet laboratory specs – we never do do this – we could expect drifts on time scale of 1-2 years.
- The KFPA has variable noise diodes.
- The noise diodes were last calibrated empirically for the GBT 10+ years ago.....
- There are significant variations in the noise diodes as a function of frequency.
- You should calibrate your data.





Flux Density vs Antenna Temp

$$P_{rec} = \frac{1}{2} A_e S_v \Delta v = k T_a \Delta v$$

 $A_e = \eta_a (\pi/4) D^2$
 $\Rightarrow S_v = 3520 T_a / (\eta_a [D/m]^2)$
i.e., $T_a / S_v = 2.84 \eta_a$ for the GBT ($\eta_a = 0.71$ at low v)

$$\eta_a = 0.352 \frac{T'_A}{S_\nu}.$$

Aperture efficiency for the GBT







History of GBT Performance [Microns] 2003 600 500 2009 Errors 400 2009b 300 Surface 2014 200 20 40 60 80 0 Elevation [deg]

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Effective surface errors determines the aperture efficiency.

Surface errors of 235+/-15 microns were measured for 2017-2019 observing seasons



Gain-Curve: Improvements to Surface Models over time







Telescope Efficiencies

$$\eta_{rss} = \frac{\iint_{2\pi} P_n(\Omega) d\Omega}{\iint_{4\pi} P_n(\Omega) d\Omega}, \quad \begin{array}{l} \text{Rear-spillover and scattering} \\ \text{efficiency -- fraction of power in the} \\ \text{forward (on sky) 2pi direction} \end{array}$$
$$\eta_l = \eta_r \eta_{rss}, \text{ where the ohmic efficiency is } \eta_r \simeq 1.0.$$

$$\eta_{fss} = \frac{\iint_{\Omega_D} P_n(\Omega) d\Omega}{\iint_{2\pi} P_n(\Omega) d\Omega},$$

Forward-spillover and scattering efficiency -- fraction of power in diffraction pattern including the error beam which is in the forward direction.

$$\eta_{mb} = \frac{\iint_{\Omega_{mb}} P_n(\Omega) d\Omega}{\iint_{4\pi} P_n(\Omega) d\Omega},$$

Main-beam efficiency – fraction of power in main-beam lobe (useful for small sources)





Antenna Temperature Scales

 $T'_{A} = T_{A} \exp(\tau_{o} A)$, Measured antenna temperature corrected for ATM

 $T^*_A = \frac{T'_A}{\eta_l}, \qquad \text{Measured antenna temperature in forward 2pi direction}$



 $T_R^* = \frac{T_A'}{\eta_l \eta_{fss}}$, Measured antenna temperature associated with antenna diffraction pattern (including error beam) – corrected for both forward and rear-ward scattering and spillover

$$T_{mb} = \frac{T'_A}{\eta_{mb}}.$$

Measured antenna main-beam temperature

$$T'_A = \eta_{mb} T_{mb} = \eta_l T^*_A = \eta_l \eta_{fss} T^*_R.$$





Point Source Calibration

$$\eta_a = 0.352 \frac{T'_A}{S_\nu}.$$

Assuming a Gaussian beam, the main-beam efficiency is related to the aperture efficiency by the following relationship:

$$\eta_{mb} = 0.8899 \eta_a \left(\theta_{FWHM} \frac{D}{\lambda}\right)^2,$$





Results from ~200 quasar measurements







Extended Source Calibration

$$\eta_{\text{Source}} = \eta_c \eta_{fss} \eta_l,$$

$$\eta_c = \frac{\iint_{\Omega_{\text{Source}}} P_n(\Psi - \Omega) B_n(\Psi) d\Psi}{\iint_{\Omega_D} P_n(\Omega) d\Omega},$$

For extended sources larger than the main-beam, need to know the detailed beam error pattern and source structure to compute source "coupling" efficiency (eta_c)

Planet calibration (disk convolved to Gaussian beam). Works best for planets smaller than mainbeam to avoid eta_c uncertainties.

$$\eta_{mb} = \frac{T'_a}{J_\nu(T_b)[1 - \exp(-x^2)]},$$
$$x = \sqrt{\ln 2}(\theta_S/\theta_{mb})$$





Extended Sources: T_{mb} vs T_{source}

$$T'_{A} = \frac{1}{\Omega_{A}} \iint P_{n}(\theta, \phi) T_{s}(\theta, \phi) d\Omega$$

$$\uparrow$$
source
$$\frac{\text{compute using } T_{mb}}{\Omega_{mb} - n} + T' - n + T$$

$$\frac{\Omega_{mb}}{\Omega_A} = \eta_{mb} ; T'_A = \eta_{mb} T_{mb}$$
$$T_{mb} = \frac{T'_A}{\eta_{mb}} = \frac{1}{\Omega_{mb}} \iint P_n(\theta, \phi) T_s(\theta, \phi) d\Omega$$





case 1: If T_s is uniform, $\theta_s < \theta_{mb} \Rightarrow \Omega_s << \Omega_{mb}$

 $\Rightarrow P_n \approx 1$ over source:







Gaussian Source and Gaussian Beam approximation (with negligible error pattern)

More general, assume T_s is Gaussian and beam Gaussian $\Omega_{Gaussian} = 1.133 \ \theta_{FWHM}^2$ $T_{mb} = T_s \left[\frac{\theta_s^2}{\theta_{mb}^2 + \theta_s^2} \right]$ $\theta_{mb} >> \theta_s \qquad T_{mb} = T_s \left(\frac{\theta_s^2}{\theta_{mb}^2} \right) \checkmark$ $T_{mb} = T_s \checkmark$





Calibration with Two Loads

Two loads

"Direct" Calibration, e.g. 4mRx

$$g = \frac{T_{amb} - T_{cold}}{V_{amb} - V_{cold}}$$
$$T_A = T_{sys} \frac{ON - OFF}{OFF} = g(ON - OFF)$$
$$T_A' = T_A e^{\tau_0 A} \qquad \underline{\text{need } \tau}$$

$$T_{cold} = 54.0 \,\mathrm{K} - 0.6 [\mathrm{K/GHz}](\nu - 77.0 \,\mathrm{GHz}),$$

Effective cold temperature for 4mm Rx GBT Memo #302





Calibration with One Load, T_A*

With a calibration vane with a temperature sensor, one can calibrate the approximate Ta* scale without any knowledge of the sky!!!

$$T_A^* = T_{cal}(C_{ON} - C_{OFF}) / (C_{amb} - C_{OFF}),$$

$$T_{cal} = (T_{amb} - T_{OFF}) \exp(\tau_o A) / \eta_l.$$

$$(T_{amb} - T_{OFF}) = T_{amb} - (\eta_l T_{sky} + (1 - \eta_l) T_{spill})$$

$$T_{sky} = T_{ATM} (1 - \exp(-\tau_o A)) + T_{bg} \exp(-\tau_o A),$$

Doing some algebra Tcal reduces to:

$$T_{cal} = (T_{ATM} - T_{bg}) + (T_{spill} - T_{ATM}) \exp(\tau_o A) + (T_{amb} - T_{spill}) \exp(\tau_o A) / \eta_l.$$

$$T_{cal} \simeq (T_{ATM} - T_{bg}) + (T_{amb} - T_{ATM}) \exp(\tau_o A).$$

 $T_{cal} \approx T_{amb}.$

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Average Tcal/Tamb ratio measured to be 0.997+/-0.018 for Argus observing sessions except for high ATM opacity.

Assuming either Tamb=Tspill or eta_l =1 which are both good approximations for the GBT. We adopt the Tcal equation in the box for Argus using the Weather database for Tatm and tau_o.



The atmosphere is important at high frequency (>10 GHz)

Opacity attenuates the signal and adds to the Tsys:
 Tsys = Trcvr + Tspill +Tbg * exp(-tau*A) + Tatm * [exp(-tau*A) – 1]
 Air Mass A~ 1/sin(Elev) (for Elev > 15°)

Stability

Tsys can vary quickly with time Worse when tau is high

 Atmosphere is in the near-field so the tau observed is similar for all beams for multi-beam receivers



GBT site has many days with low water vapor per year (<10mm H_2O are ok for 3mm, 50% of time)



Opacity vs Frequency



Total Opacity

Opacities from Various Atmosphere Components



Opacities from Various Atmosphere Components



Opacities from Various Atmosphere Components



GBT 86 GHz Performance

Table 2: 86GHz GBT Efficiency and Calibration Parameters

D	100 m
ϵ	$235\pm15\mu\mathrm{m}$
κ	1.20 ± 0.02
η_a	0.347 ± 0.032
η_{mb}	0.442 ± 0.043
η^*_M	0.465 ± 0.035
η_{Jupiter}	0.53 ± 0.05
η_{Moon}	0.814 ± 0.029
η_l	0.985 ± 0.015
η_{fss}	0.965 ± 0.020
	$\begin{array}{c} {\rm D} \\ \epsilon \\ \kappa \\ \eta_a \\ \eta_{mb} \\ \eta_M^* \\ \eta_{\rm Jupiter} \\ \eta_{\rm Moon} \\ \eta_l \\ \eta_{fss} \end{array}$

^aPower in the forward 2π direction. ^bFactional power in the forward direction inside the $\sim 1^{\circ}$ diameter error pattern.







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