A Stroll Through Radio Astronomy

Things it took me the longest time to comprehend...

Jay Lockman August 2019



- The importance of the Planck Law
- Using the ideas of radiative transfer
- Reciprocity and antenna design
- Thinking about interferometers
- Nyquist sampling and convolution
- Noise and the propagation of errors



$$I = I_0 \ e^{-\tau} + I_s (1 \ - \ e^{-\tau})$$

W m⁻² Hz⁻¹ sr⁻¹

The Planck Law for Blackbody radiation

$$B_{\nu}(\nu,T) = \frac{2h\nu^3}{c^2} \frac{1}{\exp\left(\frac{h\nu}{kT}\right) - 1}.$$





$$I = I_0 \ e^{-\tau} \ + I_s (1 \ - \ e^{-\tau})$$

For Blackbody $\tau = \infty$













$$T_b = T_0 \ e^{-\tau} \ + T_s (1 \ - \ e^{-\tau}) \quad \mathsf{K}$$



$$T_b = T_0 \ e^{-\tau} + T_s (1 \ - \ e^{-\tau})$$
 K $T_b \approx T_0 \ e^{-\tau} + T_s \ \tau \ (\tau << 1)$



 $T_b \approx T_0 \ e^{-\tau} + T_s \ \tau \quad (\tau << 1)$ $T_b = 0.99 T_0 + 3 (K)$



 $=T_0e^{T_0}$ $T_s \tau$

f(t) = f(-t)









1, 2 Spillover



4. What are the effects of blockage?



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4. What are the effects of blockage?
3. What does it mean that the surface looses efficiency? What is the effect of a hole in the dish?
Can you reduce sidelobes by wrapping legs in absorber?

The Parabolic Reflector







Blockage

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Load switch



A MEASUREMENT OF EXCESS ANTENNA TEMPERATURE AT 4080 Mc/s

Measurements of the effective zenith noise temperature of the 20-foot horn-reflector antenna (Crawford, Hogg, and Hunt 1961) at the Crawford Hill Laboratory, Holmdel, New Jersey, at 4080 Mc/s have yielded a value about 3.5° K higher than expected. This excess temperature is, within the limits of our observations, isotropic, unpolarized, and free from seasonal variations (July, 1964–April, 1965). A possible explanation for the observed excess noise temperature is the one given by Dicke, Peebles, Roll, and Wilkinson (1965) in a companion letter in this issue.

The Big Bang!

Thinking about interferometers

Radio image of the Moon





angular resolution = λ /Diam



angular resolution = λ /Diam



$$Power = S A_e \quad (w \ Hz^{-1})$$
$$S = \int T_b \ d\Omega = \langle T_b \rangle \ \Omega$$

Flux Density

$$Power = S \ A_e \ (w \ Hz^{-1})$$
$$S = \int T_b \ d\Omega = \langle T_b \rangle \ \Omega$$
$$\Omega = \frac{\lambda^2}{Diam^2}$$
Solid Angle

$$Power = S A_e \quad (w \ Hz^{-1})$$
$$S = \int T_b \ d\Omega = \langle T_b \rangle \ \Omega$$
$$\Omega = \frac{\lambda^2}{Diam^2}$$

Effective Area $A_e \propto n_{ant} D_{ant}^2$

$$Power = S A_e \quad (w \ Hz^{-1})$$

$$S = \int T_b \ d\Omega = \langle T_b \rangle \ \Omega$$

$$\Omega = \frac{\lambda^2}{Diam^2}$$

$$A_e \propto n_{ant} \ D_{ant}^2$$

$$Power \propto \langle T_b \rangle \ n_{ant} \ (\frac{D_{ant}}{Diam})^2$$



Diam = 100 m

Eler

angular resolution = λ /Diam



$$Power = S A_e \quad (w \ Hz^{-1})$$

$$S = \int T_b \ d\Omega = \langle T_b \rangle \ \Omega$$

$$\Omega = \frac{\lambda^2}{Diam^2}$$

$$A_e \propto n_{ant} \ D_{ant}^2$$

$$Power \propto \langle T_b \rangle \ n_{ant} \ (\frac{D_{ant}}{Diam})^2$$

point source

 $t \propto$ $\overline{A^2}$

extended source

 $t \propto f^2 \propto rac{Diam^4}{A_e^2}$

A digression on the sensitivity of radio telescopes

Instrument	f ²	21cm HPBW
GBT	1	9.1'
Arecibo	1	3.2'
VLA-D	~104	46"
VLA-C	~10 ⁶	14"
VLA-B	~10 ⁸	4.3"
ASKAP	~10 ⁶	

$$t \ \propto f^2 \propto \frac{Diam^4}{A_e^2}$$

A digression on the sensitivity of radio telescopes

Instrument	f ²	21cm HPBW
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For a given angular resolution, the brightness sensitivity is always greatest for a filled aperture

This is not related to the issue of missing short spacings

VLBA Limited to $T_b > 10^5$ K







If we know the beam, why can't we deconvolve?



Convolution of two Gaussians

FWHM(obs)² \approx FWHM(source)² + FWHM(beam)²

FWHM(beam)=9'

FWHM(obs) = 10'

FWHM(source) = $(100-81)^{1/2} = 4.36$

FWHM(obs)² \approx FWHM(source)² + FWHM(beam)²

FWHM(beam)=9'

 $FWHM(obs) = 10'(\pm 0.5)$

FWHM(source) = $4.36^{+1.04}_{-1.22}$

FWHM(obs)² \approx FWHM(source)² + FWHM(beam)² FWHM(beam)=9' FWHM(obs) = 10'(±1')

FWHM(source) = $4.36^{+1.96}_{-4.36}$

Deconvolution is not robust!

Nuclear winds are common in galaxies



THE LARGE-SCALE BIPOLAR WIND IN THE GALACTIC CENTER

Joss Bland-Hawthorn

Anglo-Australian Observatory, P.O. Box 296, Epping NSW, Australia

AND

MARTIN COHEN

Radio Astronomy Laboratory, 601 Campbell Hall, University of California, Berkeley, CA 94720 Received 2002 July 2; accepted 2002 September 4



γ -ray Data From Fermi










γ -ray Data From Fermi



Fermi Bubbles Artist's Conception







HI column density $(10^{18} \text{ cm}^{-2})$

McClure-Griffiths et al. 2013 Australia Telescope Compact Array





*More Area *More Sensitive *Further out in Fermi Bubble

GBT Hydrogen Clouds in the Fermi Bubble Wind



GBT Hydrogen Clouds in the Fermi Bubble Wind





Model of the Hydrogen Clouds being Ejected from the Galactic Center



Vexp = 330 km/s







