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JPL PHYSICAL OPTICS SCATTERING PROGRAM

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I. INTRODUCTION

This report describes a FORTRAN program developed at JPL for the calculation of the scattered pattern from a reflector of arbitrary shape. The program written by A. Ludwig is contained in a report entitled "Calculation of Scattered Patterns from Asymmetrical Reflectors" (Technical Report 32-1430; 2/15/70). The JPL report stresses the following: the use of a Fourier expansion representation of the reflector surface, the possible inaccuracy of a far field assumption for the fields incident on the reflector, and the use of a fast integration technique developed by Ludwig. Some good examples of the accuracy of the program are included. A copy of the JPL report may be obtained from Rick Fisher in Green Bank or from Sarah Martin (microfiche) in Charlottesville.

The program is actually two programs: one which calculates a scattered pattern and one which computes, using a spherical wave expansion, a completely general representation of the incident field pattern. The programs will be referred to as SCAT and SWE, respectively.

The method of surface specification has been altered to that of tabular form and the integration routine has been slightly adjusted to better handle reflector edge contributions. No major changes have been made to the SWE program.

SCAT employs the technique of physical optics to calculate the scattered field. The program takes the Fourier transform, using a fast integration technique, of an array of induced current dipoles.

The incident field representation provided by the SWE program allows accurate calculations to be made even with the reflector well into the near field of the source. Input into the SWE program is the far field source pattern, such as those typically measured on an antenna range. The far field form of

spherical waves is matched to the (far field) incident pattern. Backing up into the near field, the SWE accurately approximates the transverse and radial components of the electric field.

(An attempt has been made towards the end of this report to describe in physical terms some of the theory behind the SCAT and SWE programs.)

One application of this program is in checking the taper at the edge of the main dish of the 140-ft Cassegrain telescope. For this problem the scattering surface is the subreflector, which has an asymmetrical edge, and the incident field is the antenna range (or computed) pattern of the appropriate Cassegrain feed.

A slight variation of this situation involves the re-design of the subreflector to include a conical vertex plate and a flange about the reflector edge.

Another application is to aid in the design of beam polarization splitters which may be added to the 140-ft. To the SCAT program, the beam splitter would be a tilted plane reflector placed very near the feed aperture.

II. GENERAL PROGRAM CHARACTERISTICS

A. Summary of What Program Does

1. SCAT program.

Given a scattering surface specified by $\rho(\theta, \phi)$ (i.e., spherical coordinates) and the magnetic fields incident on this surface, the scattered pattern is computed over a grid of observation points. The scattering surface is assumed to be perfectly conducting; therefore, the electric field is zero at the surface.

The values of ρ are either input in tabular form or calculated from an equation in θ and/or ϕ .

The E and H plane values (dB or volts, and degrees) of the incident magnetic field are read in by the SWE program. The SWE program computes the coefficients of the expansion and stores them. These stored values are read in by the SCAT programs.

Ludwig has shown that assuming far field conditions for the source field for the trivial case of scattering from an infinite plane reflector located within or near the traditional far field boundary line ($2D^2/\lambda$, D = source diameter) can result in rather strong backlobes in the scattered pattern. The strength of the backlobe and distortion of the main (reflected) lobe increases as the reflector further penetrates the near field; in other words, the magnitude of the error depends on frequency. (See JPL report, figures 3-5 on pages 6-8.) Ludwig further shows (Figure 7) that use of the spherical wave representation almost entirely eliminates the problem.

The scattered fields are, as mentioned above, computed using the method of physical optics. It is assumed that the induced currents are zero on shadowed portions of the reflector and that on directly illuminated areas the induced surface current value is twice the value of the tangential (to the surface) incident H field. These two assumptions constitute the physical optics approximations.

As a means of reducing computation time the far field form of the scattered pattern is calculated; the near field form can be found by submitting the scattered field values to the SWE program.

Summing the product of the induced surface currents and phase delay (path-length) for each Δs of the surface, the scattered field as seen from a particular point on the output grid is determined. This integration is performed for each point on the output grid. Finally, the far field source pattern is added to the scattered pattern yielding what is termed the "total fields". For the

situation of an infinite plane reflector, the total fields would be the incident fields pointed in the opposite direction on the output grid and with a 180° phase change.

From comparisons of computed and measured subreflector patterns, Ludwig has shown the computed pattern to be accurate down to -35 dB (relative power) and even through the first side lobe. Fairly reliable spillover efficiency calculations should be possible, since the major spillover contribution is generally from fields within 25 dB of the pattern maximum. Proper handling of the side lobes requires the use of the Geometrical Theory of Diffraction (GTD) which does not make the assumption of zero induced currents on shadowed reflector surfaces.

2. SWE program.

The SWE program finds the coefficients of expansion of the incident (far zone) magnetic field pattern in terms of transverse electric (TE) and transverse magnetic (TM) spherical waves. The TE and TM spherical waves are the general vector solutions to Maxwell's equations for an electromagnetic wave travelling in a source-free region, V. For the case of representing a magnetic field, the TM vector solution has no radial components, whereas the TE solution has components in all three coordinate directions.

A spherical wave expansion operates under the same mathematical principles as the Fourier expansion of a function. In a Fourier expansion, a function is represented by a summation of sines and cosines each multiplied by a coefficient particular to the order of sine and cosine variation; this order of variation could be called the mode order. To represent the function, the two sets of coefficients (one set multiplying the sines, the other the cosines) must be determined. This is done by evaluating integrals involving sines and cosines and specific function values. The integrals are simplified through the use of the orthogonality relations between sines and cosines.

In a spherical wave expansion the sines and cosines of the Fourier expansion are replaced with the TE and TM spherical wave solutions of Maxwell's equations. The integrals for the mode coefficients contain the corresponding TE and TM mode functions and the far field incident pattern values. Thus, the procedure for determining the coefficients of the spherical wave expansion is basically the same as for that of the Fourier expansion.

The expansion is in three variables, in this case the spherical coordinates, and contains two mode orders. Thus, the expansion is a double summation; one summation for each mode variable. One mode order, generally called the mode order and designated by a "n", specifies the degree of ρ variation. The other mode order, called the order of azimuthal variation and designated by an "m", specifies the degree of ϕ variation. The degree of θ variation depends on both mode orders.

In principle, the expansion scheme outlined above can represent any input pattern at any distance from its source. However, because the case of $m = 1$ is of particular importance, the double summation has been reduced in the program to a single summation over n. Also, one set of spherical wave solutions has been neglected, which means that the incident radiation is assumed to be linearly polarized.

B. Coordinate System Used.

The origin of the system will be at the phase center of the source fields. The reflecting surface is specified by the vector $\bar{\rho}$ which is a function of the angles θ and ϕ ; ρ , θ , ϕ represent a point on the surface in spherical coordinates. Since $\bar{\rho}$ is defined by θ and ϕ , these two angles will specify the point of integration.

The output grid over which the observer views the scattered pattern is also designated by spherical coordinates (R , θ , Φ). Since the far field form of the scattered pattern is computed, the output grid is defined entirely by θ and Φ .

The Z-axis is the reflector axis. The three coordinate sets — (ρ, θ, ϕ) , (R, θ, Φ) and (X, Y, Z) — all have the same origin. The coordinate system is shown in Figure 1.

C. Program Structure and Order of Calculation.

1. SCAT program.

For storage reasons, the reflector surface over which the integration is performed is divided into integration grids. The scattered fields from each grid are superimposed and added to the incident fields to yield the total (scattered) pattern.

The reflector is divided along θ and/or ϕ (Figure 2), depending on the size of the reflector, the choice of $\Delta\theta$ and $\Delta\phi$ and storage requirements. This partitioning introduces no appreciable error and is quite useful since the total integration may cover several thousand points.

The output grid is not segmented because it generally includes only several hundred points at most.

For each integration grid, $\rho(\theta, \phi)$ is either read in or calculated for every (θ, ϕ) point on the grid. Since the normal to the surface is required at each integration point, $\frac{\partial\rho}{\partial\theta}(\theta, \phi)$ and $\frac{\partial\rho}{\partial\phi}(\theta, \phi)$ are computed over the grid. The derivatives are computed using the numerical method of backward and forward differences or by an analytic expression in θ and ϕ which is inserted in subroutine SURF.

In general the outermost (in θ) integration grids will contain points that are beyond the edge of the reflector. This is so because in the program, θ and ϕ are independent of each other; that is, the integration grid is represented by two vectors rather than by a matrix of points. However, at the reflector edge, the boundary is defined by θ_{edge} , which will generally be a function of ϕ . To correctly represent the reflector, the edge values of θ are read in or calculated for each ϕ of the integration grid.

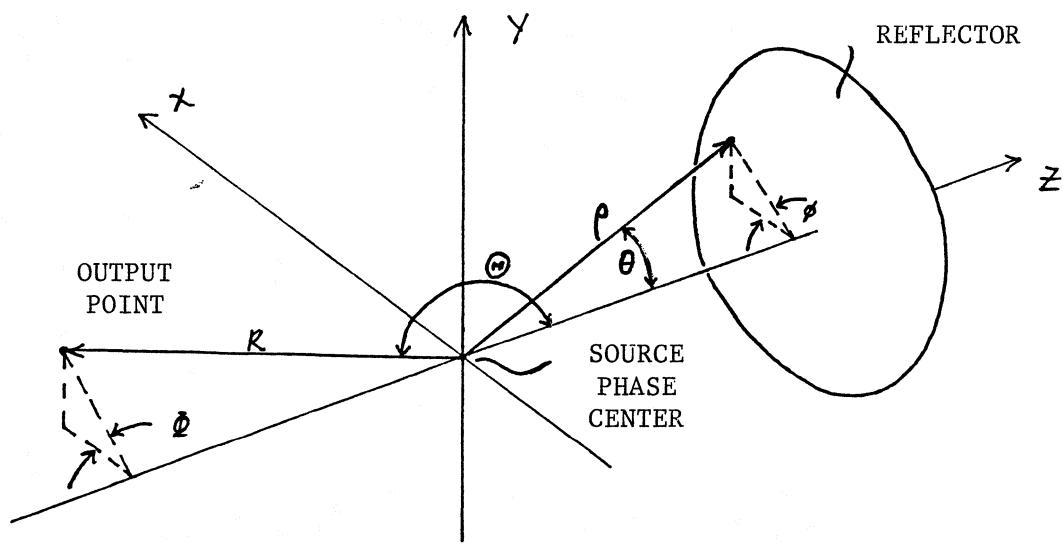


FIGURE 1

Coordinate system (from JPL report, page 3).

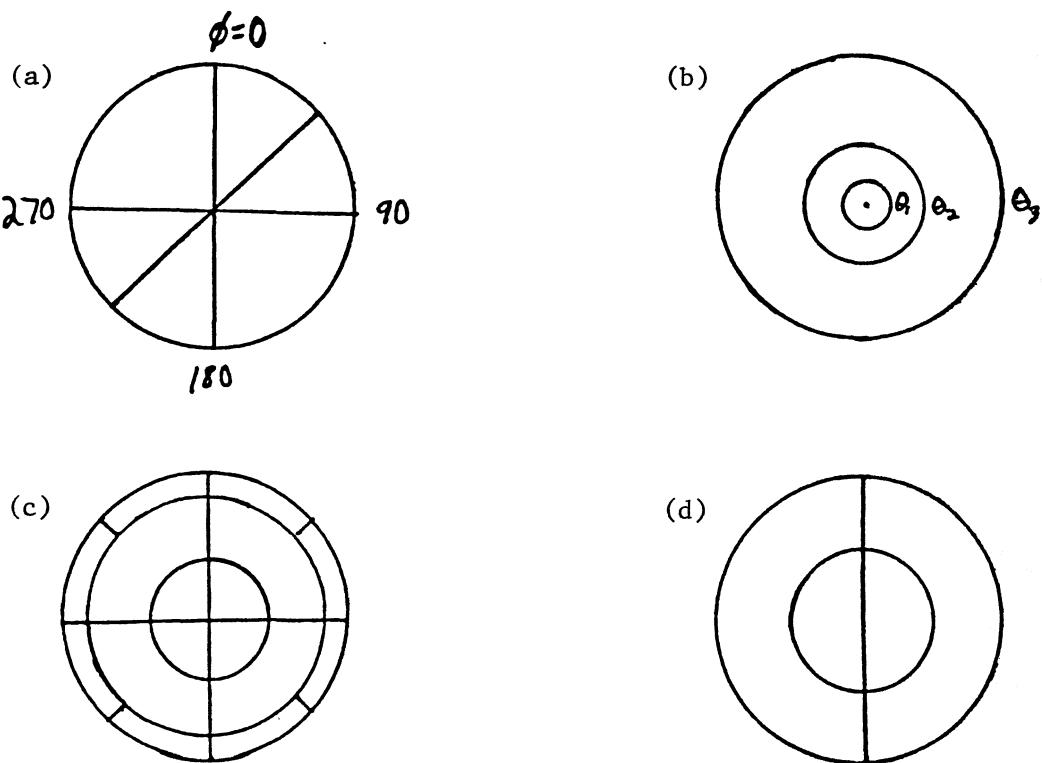


FIGURE 2

Sample integration grids.

An integer parameter, IEDGE, is used to determine if the edge will be encountered for the present integration grid and, if it is encountered, how the edge values will be obtained. IEDGE is read in by the MAIN subroutine of the SCAT program for each integration grid.

As mentioned above, the edge integration grids will contain points beyond the actual reflector. Since such values may not be known, they are not required inputs. When the values of $\rho(\theta, \phi)$ are in tabular form, the user can set to 999.99999 the first ρ value that corresponds to a point beyond the reflector edge; the ρ values remaining for that given value of ϕ can (and must) be set to some arbitrary value (e.g., zero).

If all the values of ρ for an edge integration grid are known, as is the case when ρ is calculated analytically, then including them will result in slightly better accuracy.

Note that the SCAT program is at its maximum user-convenience when ρ (in particular) and the edge values of θ can be computed from equations inserted in the program.

An informal flowchart of the SCAT program is shown in Figure 3. Subroutine names are written next to most flowchart components.

2. SWE program.

The organization of the SWE program is considerably simpler; therefore, no flowchart of it is shown. A mode order term and the far field E and H plane pattern values from the appropriate feed comprise the inputs. Details of input considerations will be given below.

After reading in the input pattern, the program establishes a matrix of θ mode weights (forms of Legendre polynomials) and a matrix of input pattern differentials. The evaluation of the coefficient integrals is then done by matrix multiplication. The real and imaginary values of the TE and TM coefficients are normalized and stored on the computer disk. As a check on the validity of the coefficients, the far field form is computed allowing a comparison to be made with the input pattern.

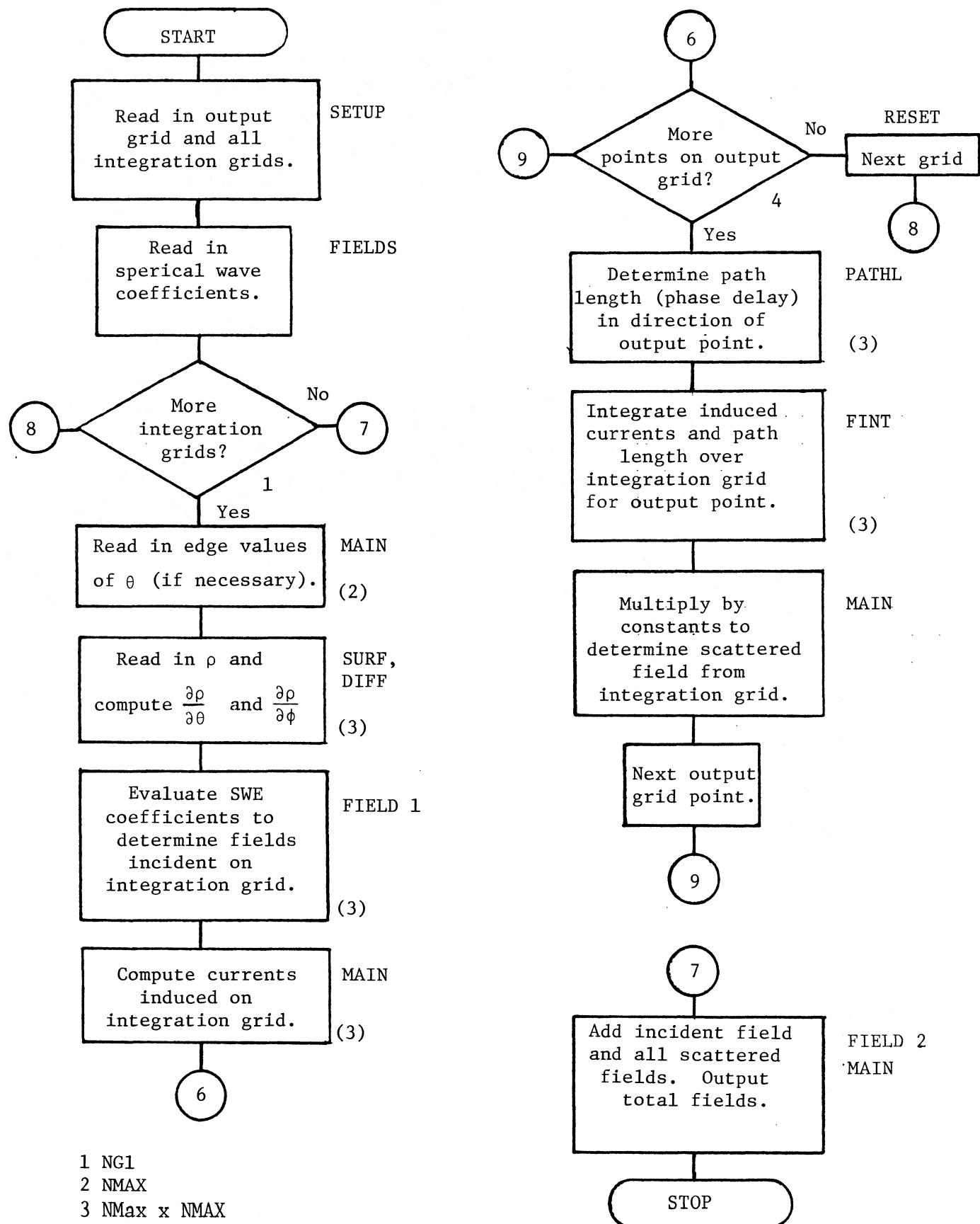


FIGURE 3

SCAT program flowchart.

III. DETAILS OF USING THE PROGRAM

Both programs are stored on the disk of the IBM 360 computer located in Charlottesville. The user command system is called Pandora. Each subroutine is stored as a Pandora member and has a member name either identical to or similar to the subroutine name. The Pandora member name (PMN) will be given for each subroutine.

For both the SCAT program and the SWE program this section provides the following information:

For each subroutine of the program:

Subroutine description, input variables (that READ data), input considerations, and a table of inputs and their formats.

Discussion of program output.

A. Subroutines of Scattering Program

1. MAIN (PMN MAIN1SC)

Subroutine description:

This routine coordinates the efforts of the program. The procedure followed by MAIN has been outlined by the SCAT program flowchart (Figure 3). Note the SUBROUTINE names written by the flowchart blocks.

Input variables and their tasks:

TITLE - any alphanumeric statement \leq 72 characters.

PC - phase constant, $k = 2\pi/\lambda$ (λ in meters).

XT, YT, ZT - translations to expected phase center of scattered pattern (affects only output phase data).

SCALE - scale factor for output fields. If input is zero, program sets SCALE = 1.0.

(The following variables are read in for each integration grid.)

IEDGE - an integer; determines if and how program obtains θ at the reflector edge. There are four cases:

IEDGE (continued):

IEDGE < 0 - Program calculates values for that integration grid using equations inserted into subroutine EDGEEQ.

IEDGE = 0 - θ at edge is constant. Program reads in just one value.

IEDGE = 11 - Edge will not be encountered on present integration grid. No values read in.

IEDGE > 0 - θ at edge is a function of ϕ ($\neq 11$) and is read in for each value of ϕ in the integration grid (NMAX values).

TEDGE (NMAX) - edge values of θ in degrees. Either NMAX values, one value or no values are read in depending on IEDGE.

Input considerations:

TITLE usually describes reflector and frequency. The scattered field phase center translations may be set to zero if the output phase is of no interest.

The choice of $\Delta\theta$ and $\Delta\phi$ (which determines the number of integration grids) is covered in SETUP.

Table of inputs and their formats:

<u>Card</u>	<u>Input Variable(s)</u>	<u>Format</u>
1	TITLE	18A4
2	PC, XT, YT, ZT, SCALE	5F10.4
[SETUP DATA]		
[FIELDS DATA]		
Read in for each integra- tion grid.	IEDGE TEDGE (Either NMAX, 1 or no values.)	I5 7(1X, F9.5)
[SURF DATA]		

2. Setup (PMN SETUP)

Subroutine description:

This routine sets up the output grid and all integration grids. The integration grid parameters are reset for each new grid by ENTRY RESET within SETUP.

The input parameters of the output grid consist of the number of values (JMAX and KMAX), initial value and increment which define the θ and ϕ points over which the fields are viewed. A typical grid could be $\phi = 0^\circ$ and 90° and $\theta = -90^\circ$ to $+90^\circ$ with a step of 2° . Thus, KMAX = 2, JMAX = 91.

The input parameters of each input grid are identical to those of the output grid but define θ and ϕ instead. As an example, consider Figure 2(d). The first grid might be $\theta = 0^\circ$ to 6° with $\Delta\theta = 0.2^\circ$ and $\phi = 0^\circ$ to 180° with $\Delta\phi = 4^\circ$. Thus, MMAX = 31 and NMAX = 46. The next grid might be $\theta = 0^\circ$ to 6° and $\phi = 180^\circ$ to 360° . Two more grids would follow. The order of the four grids is of no consequence.

Input variables and their tasks:

JMAX	-	number of θ values on output grid, ≤ 181 .
TT1	-	initial θ value.
DDT	-	θ increment.
KMAX	-	number of ϕ values on output grid, ≤ 5 .
PP1	-	initial ϕ value.
DPP	-	ϕ increment.
NG1	-	number of integration grids, ≤ 21 .
MM(I)	-	number of θ values on Ith integration grid, ≤ 36 .
TI	-	initial θ value.
DT	-	θ increment ($\Delta\theta$).
NN(I)	-	number of ϕ values on Ith integration grid, ≤ 91 .
PI	-	initial ϕ value.
DP	-	ϕ increment ($\Delta\phi$).

(All initial values and increments in degrees.)

Input considerations:

$\Delta\theta$ and $\Delta\phi$ are chosen such that the dimensions of the largest ΔS on the reflector surface are roughly one wavelength by one wavelength. This ensures that the pathlength term does not vary by more than 2π over any ΔS . Thus, $\Delta\theta$ can be found from $\Delta\theta \approx \lambda/\rho_{\max}$, where ρ_{\max} is the greatest distance to the reflector and $\Delta\theta$ is in radians. The corresponding equation for $\Delta\phi$ is $\Delta\phi \approx \Delta\theta/\sin(\theta_{\max})$, where θ_{\max} is the largest θ angle subtended by the reflector edge.

The values obtained for $\Delta\theta$ and $\Delta\phi$ should be rounded to the nearest convenient value (tenths or hundredths of a degree). An upper limit of $\Delta\theta = 1^\circ$ is necessary, since larger values sample the incident field too sparsely. Decreasing $\Delta\theta$ and/or $\Delta\phi$ offers a good check on the validity of the chosen increments. However, halving either increment nearly doubles the computation time. To give some idea of the CPU time needed for a given JMAX, KMAX, TMMAX, TNMAX combination (define TMMAX and TNMAX as total number of θ and ϕ points on reflector surface), the following table is provided.

KMAX	JMAX	TMMAX	TNMAX	CPU TIME (minutes)	
1	91	35	92	18	
1	91	35	184	36	IBM
1	181	35	184	50	360/65
2	181	35	184	89	

To ensure accurate edge contributions, the user should be certain that for any value of ϕ the edge value of θ does not fall within the θ range of the first two values of θ (inclusive) on the integration grid. For most cases, the edge value of θ could be within this range and would produce no noticeable differences in the total scattered pattern. It is, however, absolutely necessary that all edge θ values be greater than or equal to the first θ of integration.

The integration routine and SURF subroutine set the limit of MMAX and NMAX ≥ 4 .

Table of inputs and their formats:

<u>Card</u>	<u>Input Variable(s)</u>	<u>Format</u>
1	JMAX, TT1, DTT	I5, 2F10.2
2	KMAX, PPI, DPP	I5, 2F10.2
3	NG1	I5
4	MM(1), T1, DT	I5, 2F10.2
5	NN(1), P1, DP	I5, 2F10.2
:	:	:
2NG1 + 2	MM(NG1), T1, DT	I5, 2F10.2
2NG1 + 3	NN (NG1), P1, DP	I5, 2F10.2

3. FIELDS (PMN FIELDS)

Subroutine FIELDS reads, off the computer disk, the spherical wave coefficients of the expansion of the incident fields. As mentioned previously, these coefficients are written on the disk by the SWE program.

FIELDS contains two entry points, ENTRY FIELD1 and ENTRY FIELD2. FIELD1 evaluates the SWE coefficients to find the near field form of the pattern incident on the reflector.

FIELD2 determines the electric field at an infinite distance from its source. FIELD2 is used in calculation of the far field form of the incident pattern.

Both FIELD1 and FIELD2 use the SWE coefficients in their computations.

Input variables:

- TITLE - alphanumeric statement, \leq 70 characters.
- LMAX - maximum mode order, \leq 70.
- MCOMP - order of azimuthal variation.
- A(N,1), A(N,2) - real and imaginary components of $TE_{MCOMP,N}$ ($N = 1$ to LMAX) spherical wave coefficient.
- B(N,1), B(N,2) - real and imaginary components of $TM_{MCOMP,N}$ spherical wave coefficient.

Input considerations:

The user need not include the above information with the other data for the SCAT program. Instead, FIELDS reads this data directly off the computer disk. The first three variables are inputs of the SWE program; the coefficients are calculated by it. The input considerations and table of formats for these variables are given in the SWE program description.

4. SURF (PMN SSURF).

Subroutine description:

For each integration grid, this subroutine reads in all $\rho(\theta, \phi)$ values and calculates $\frac{\partial \rho}{\partial \theta}(\theta, \phi)$ and $\frac{\partial \rho}{\partial \phi}(\theta, \phi)$. There are four input cases that SURF considers; these are determined by the integer input parameter ISURF. The four cases are the following: ρ is a function of θ and ϕ and is read in from a table of values; ρ is a function of θ and is read in from a table of values; ρ is known in analytic forms as a function of θ and/or ϕ ; ρ , $\frac{\partial \rho}{\partial \theta}$ and $\frac{\partial \rho}{\partial \phi}$ are all known in analytic form. In the first situation, MMAX x NMAX values are read in (every point on the integration grid); in the second case, MMAX values are read in (every θ point on integration grid); in the remaining cases, no values are read in. When $\frac{\partial \rho}{\partial \theta}$ and $\frac{\partial \rho}{\partial \phi}$ are not analytically specified, they are numerically determined using the techniques of forward and backward differences. (See Computer Methods for Science and Engineering, LaFara.) This method of computing derivatives at the points of a table of data has been found to often be more accurate than is necessary for the SCAT program. For differentiable surfaces and not overly varying integration grids (as outlined in SETUP), this accuracy is maintained down to tables of four values. Thus, MMAX and NMAX should be greater than or equal to four.

To compute ρ and/or its partial derivatives, equations in θ and/or ϕ are inserted into SURF. The values of $\sin \theta$, $\cos \theta$, $\sin \phi$, $\cos \phi$, θ and ϕ (radians)

are stored for each integration grid in the variables SIT(M), COT(M), SIP(N), COP(N), T(M), P(N).

As previously mentioned, when ρ is input in tabular form, a value of 999.99999 indicates the reflector edge has just been passed. For a given ϕ , the remaining θ points on the integration grid must be assigned some arbitrary ρ value, which does not figure in any computations. If ρ is known for points beyond the edge, it would be slightly more accurate to include the value immediately beyond (on the integration grid) the edge value. Doing this would allow ρ at the edge to be interpolated rather than extrapolated.

Input variables:

ISURF - integer parameter that determines how ρ is obtained.

There are four cases as follows:

ISURF < 0 - ρ , $\frac{\partial \rho}{\partial \theta}$ and $\frac{\partial \rho}{\partial \phi}$ all determined analytically.
(# -2)
Read no values.

ISURF = -2 - ρ determined analytically; $\frac{\partial \rho}{\partial \theta}$ and $\frac{\partial \rho}{\partial \phi}$ numerically computed (backward and forward differences). Read no values.

ISURF = 0 - ρ is a function of θ only and is read
(MMAX values) from a table.

ISURF > 0 - ρ is a function of θ and ϕ and is read
(MMAX x NMAX values) from a table.

F (M, N) - $\rho(\theta, \phi)$

FT (M, N) - $\frac{\partial \rho}{\partial \theta}(\theta, \phi)$

FP (M, N) - $\frac{\partial \rho}{\partial \phi}(\theta, \phi)$

Input considerations:

If the partial derivatives of ρ are to be determined numerically, MMAX and NMAX can at minimum be two but it is suggested that they be no less than four.

A value of ρ is required for every (θ, ϕ) or θ (depending on ISURF) on the integration grid.

For the case of reading in ρ for each (θ, ϕ) point on the integration grid, the θ DO loop is within the ϕ DO loop. Thus, for each ϕ value a new block of data is started.

The above two restrictions must be kept in mind when setting a value of ρ to 999.99999.

If at all possible, ρ should be represented by an equation. For this situation the input procedure is at its simplest.

Table of inputs and their formats:

	<u>Card</u>	<u>Variable(s)</u>	<u>Format</u>
Read in for each integra- tion grid.	1	ISURF (For ISURF < 0, no values read.)	I5
	2	F(1, 1) to F(MMAX, 1)	7(1X, F9.5)
	.	or to F(MMAX, NMAX)	
	.		

5. DIFF (PMN DIFF)

Subroutine description:

This subroutine is called by SURF to numerically approximate the derivative of a dependent variable with respect to the independent variable at each point in a table of data. The technique of forward and backward differences is employed; forward near the front of the table and backward elsewhere. By reaching back or forth into the table of data and taking the differences of the dependent variables along the way, the method is able to accurately compute derivatives of rapidly changing data within the table. The spacing of the independent variable ($\Delta\theta$, $\Delta\phi$) must be constant throughout the table (integration grid). The table should contain a minimum of four values.

There are no input variables.

6. FINT (PMN EFINT, FINT)

This subroutine performs the integration of the product of the induced surface currents and the pathlength over each integration grid for each point on the output grid. It is here in the program where the great majority of number crunching is performed. This integration routine was developed by Ludwig specifically for the form of integral encountered in scattering problems.

(The fast Fourier transform is not applicable.) Ludwig's method reduces by a factor of 16, relative to a Simpson's rule integration, the number of integration points needed.

In this integration routine the θ loop is embedded within the ϕ loop.

The reflector edge is handled in the following way: If IEDGE indicates an edge is present, then for a given azimuthal angle FINT compares each $\theta + \Delta\theta$ value with the edge value of θ (which is a function of ϕ). When the edge is encountered, the edge θ for ϕ and for $\phi + \Delta\phi$ are averaged. The contribution from this smaller (in some cases, slightly larger) ΔS is then evaluated by extrapolating or interpolating the pathlength and surface currents to the averaged edge θ . (See Figure 4.). The reflector, in effect, has a discontinuous edge.

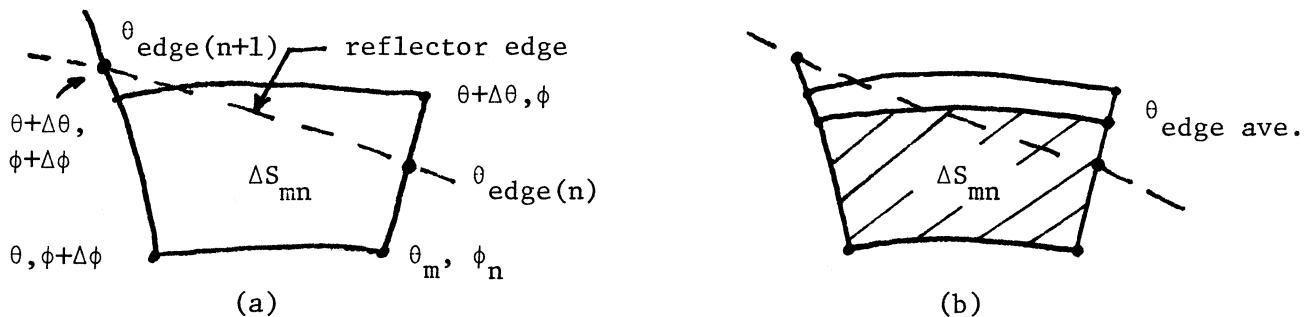


FIGURE 4

Another routine has been developed with the intent of better representing the edge. The ΔS_{mm} in Figure 4(a) is divided along ϕ , the number of divisions being determined by the difference between the two edge values of θ . θ_{edge} is then interpolated to each ϕ division. The same technique of θ averaging is applied over the reduced edge section, yielding smaller edge discontinuities. This integration routine is stored in PMN EFINT.

7. PATHL (PMN PATHL).

Subroutine description:

This member computes the pathlength from every point on the integration surface in the direction of each point on the output grid. The pathlength is the dot product between the vector ρ and the unit vector in the direction of the output variables.

8. SPHANK (PMN SPHANK).

Subroutine description:

SPHANK computes the term $h_n^{(2)}(k\rho) \cdot (-j)^{n+1} \cdot k\rho \cdot e^{jk\rho}$, where $h_n^{(2)}(k\rho)$ is the spherical Hankel function of the second kind, n is the mode order, and $j = \sqrt{-1}$.

This subroutine is called by FIELD1 (magnetic field at finite distance from source). The value SPHANK determines provides the near field ρ dependence. For far field calculations, this factor must be removed. Otherwise, SPHANK would provide only the value of $h_n^{(2)}(k\rho)$.

The spherical Hankel function represents an outward travelling spherical wave. $h_n^{(2)}(\chi)$ is defined as $j_n(\chi) - j y_n(\chi)$ where j_n and y_n are spherical Bessel functions. $j_n(\chi)$ and $y_n(\chi)$ are analogous to $\sin \chi$ and $\cos \chi$, and $h_n^{(2)}(\chi)$ is analogous to $e^{-j\chi}$.

SPHANK provides the only (known) numerical program limitation. For $\chi < 100$ (roughly), $y_n(\chi)$ blows up as n increases. As χ decreases, the blow up

occurs for smaller n. This is not a computing problem but is a characteristic of the spherical Bessel function, y_n . (See Handbook of Mathematical Functions, Abramowitz and Stegun; p. 438, Fig. 10.1 to 10.3 and pp. 465-466, Table 10.5.)

To avoid computer overflow SPHANK stops calculation when $y_n \sim 10^7$. For a wavelength of 5 cm ($k = 125.7 \text{ m}^{-1}$) and a plane reflector at a distance of 1 m, this truncation produced no changes in the scattered field phase or amplitude. It is suggested that the problem be such that $k\rho > 100$. For $k\rho < 100$, the above tables should be consulted to determine the maximum "safe" mode order.

There are no input variables.

9. LEGEND (PMN LEGEND)

Subroutine description:

LEGEND computes the value of the associated Legendre polynomial, $P_n^m(\cos \theta)$, where m is the order of azimuthal variation and n is the mode order.

The Legendre polynomial provides the correct polar angle (θ and Θ) variation of the spherical wave field representation. Therefore, it is called by both FIELD1 and FIELD2.

LEGEND has no numerical difficulties analogous to those of SHPANK. (See Tables of Functions, Johnke and Emde; pp. 112-113.)

There are no input variables.

10. VECTOR (PMN VECTOR)

Subroutine description:

This subroutine converts complex numbers in rectangular form (real and imaginary) to polar form (amplitude and phase). Most program computations are done with numbers in rectangular form.

There are no input variables.

11. ADJUST (PMN ADJUST)

Subroutine description:

ADJUST normalizes phase angles to the range of -180° to $+180^\circ$.

There are no input variables.

12. EDGEEQ (within PMN SSURF)

Subroutine description:

Equations for the edge value of theta are inserted into this subroutine. The same values of $\sin \phi$, $\cos \phi$, etc., used in SURF are available. The edge equations will be functions of ϕ only.

If necessary, the equations can be divided up for the different ϕ integration grids.

There are no input variables.

B. Output of Scattering Program.

The printout of the program is fairly self-explanatory. Briefly, the bulk of the output is the following:

<pre>_____ Output for each inte- gra- tion grid.</pre>	<ul style="list-style-type: none"> -- Output grid of all integration grids. -- Spherical wave coefficients. -- If edge is to be encountered, all edge values of theta. -- A selection of the values of ρ, $\frac{\partial \rho}{\partial \theta}$ and $\frac{\partial \rho}{\partial \phi}$. -- Scattered fields (E and H planes of far-electric field) <ul style="list-style-type: none"> from integration grid. -- Far field incident pattern. -- Superposition of incident fields and all grid scattered <ul style="list-style-type: none"> fields (i.e., yields total scattered fields).
--	---

C. Subroutines of Spherical Wave Expansion Program.

The SWE program determines the coefficients of the spherical wave expansion of an input far field pattern, such as those typically measured on an antenna range. In principle, this expansion is completely general and provides one with the near field values, including radial components, of the input pattern. (For a mathematical account of spherical wave theory, see Electromagnetic Theory, Stratton; Chapter 7.)

1. MAIN (PMN MAINOR)

Subroutine description:

This routine is essentially the entire program. It reads in the far field pattern, computes the coefficients, writes the coefficients into a disk data set, and computes the far field form of the spherical wave expansion (for comparison with the input pattern).

Input variables:

- TITLE - Alphanumeric statement, \leq 72 characters (identify program).
- MCOMP - Order of azimuthal variation.
- LMAX - Maximum mode order (θ), \leq 80.
- TITLE - Same as above (identify input pattern).
- JIN - Number of input field points, \leq 121.
- IC1 - If \leq 0, convert incident field from dB to volts.
- IC2 - If $>$ 0, neglect incident field phase.
- IC3 - If $>$ 0, compute incident field amplitude from equations inserted in MAIN.
- IC4 - If $>$ 0, $E_\theta = E_\phi$ in phase and amplitude; input each phase and amplitude once.
- PSI - Polar angle θ .
- E - $E_\theta(\theta, \phi)$ amplitude as a function of θ of input pattern; volts or dB.
- EP - $E_\theta(\theta, \phi)$ phase; degrees.
- H - $E_\phi(\theta, \phi)$ amplitude; volts or dB.
- HP - $E_\phi(\theta, \phi)$ phase.

Input variables (continued):

JMAX0 - $180/\Delta\theta + 1$ where $\Delta\theta$ is the desired output increment of the far field SWE pattern.

JOUT - Number of output values starting with $\theta = 0^\circ$.

Input considerations:

As mentioned previously, only one azimuthal expansion component can be handled by the SWE program. This component is usually $m = 1$, although it can be zero or > 1 . The $m = 1$ mode typically corresponds to a pattern amplitude which varies as a full cycle sinusoid in total azimuth. It is for this case that one refers to the "E and H planes" of a pattern. If more than one azimuth component is needed, the components can be run separately and the resulting coefficients superimposed.

The number of n modes necessary to accurately represent the input pattern depends on the complexity of that pattern. For the E plane (E_θ) pattern of the C-band Cassegrain feed at 5 cm, roughly $n = 70$ modes were needed. Since the SWE program requires only about one minute of CPU time, the user can easily test various maximum mode orders and check which yields the most accurate far field pattern.

Comparing the far field spherical wave expansion pattern to the input pattern, one can see that the former tends to oscillate about the latter. Frequent sampling of the input pattern will help to minimize this. For the 5 cm case above, feed pattern values were input every 0.25° until the input pattern was about 35 dB down ($\theta = 20^\circ$) and every 0.5° out to -45 dB ($\theta \approx 30^\circ$).

An oscillation of several hundredths of a volt in the approximation pattern values is of little concern since the input pattern is usually not reliable to such resolutions. A more important consideration is that the patterns contain roughly the same power through the average angle out to the edge of the reflecting surface.

Table of inputs and their formats:

<u>Card</u>		<u>Format</u>
1	TITLE	18A4
2	MCOMP, LMAX	2I5
3	TITLE	18A4
4	J1N, IC1, IC2, IC3, IC4	5I5
5	IC4 ≤ 0 : T(1), E(1), EP(1), H(1), HP(1)	5(1X, F9.5)
	IC4 > 0 : T(1), E(1), EP(1)	3(1X, F9.5)
:	:	:
JIN+4	IC4 ≤ 0 : T(JIN), E(JIN), EP(JIN), H(JIN), HP(JIN)	5(1X, F9.5)
	IC4 > 0 : T(JIN), E(JIN), EP(JIN)	3(1X, F9.5)
JIN+5	JMAX0, JOUT	2I5

2. MULT (PMN MULT)

Subroutine description:

MULT multiplies two matrices. The matrix dimensions are specified in the CALL MULT statement. If any variable dimensions are changed in MAIN, the CALL MULT statement must be changed accordingly.

3. LEGEND (PMN LEGEND)

4. VECTOR (PMN VECTOR)

LEGEND and VECTOR described previously.

D. Output of the SWE Program.

-- E and H planes of input pattern in volts and degrees.

-- Real and imaginary values of TE and TM wave coefficients for each mode.

-- Fraction of total mode power of the coefficients for each mode.

- Total mode power of coefficients.
- Far field summation of spherical modes;
E and H planes in volts and degrees.

E. Submitting a Job.

It will be assumed that the user is already familiar with the Pandora command system of the IBM 360. If the user is not familiar with the system, the Pandora Guide (assembled by the Charlottesville Computer Division) nicely explains it. Some essential member-oriented commands are the following: ENTER, FETCH, SAVES, CLEAR, SCRATCH, CWS, CONCAT, SEQUENCE and SUBMIT. Some essential line editing commands are the following: CHANGE, INSERT, DELETE, MOVE, COPY, EDIT and SEEK.

The necessary JCL (Job Control Language) parameters for the programs are contained in the Pandora members JCLSCAT and JCLSWE (some message suppression). Only the TIME parameter is set by the user. This specifies a maximum CPU time in minutes. In JCLSWE, TIME = 1 and need not be adjusted.

To submit the SCAT program, the following statement is entered:

SUBMIT_JCLSCAT_MAIN1SC_SSURF_FINT_SFPVSALD_DATANAME

(PMN SFPVSALD contains eight subroutines.)

In this statement, JCLSCAT must be first and SFPVSALD and DATANAME must be last and in the shown order. The three other members can be arbitrarily shuffled. DATANAME is specified by the user.

To submit the SWE program, the following statement is entered:

SUBMIT_JCLSWE_MAINOR_LVM_DATANAME

(PMN LVM contains three subroutines.)

This statement must be entered with the order as shown.

It must be kept in mind that whenever the incident field is to be changed in the SCAT program, the SWE program must be run so that the proper coefficients are stored on the disk.

The following Pandora members may prove useful:

DSCATSUB - Contains input data for scattering the 140-ft subreflector; $\Delta\theta = 0.2^\circ$, $\Delta\phi = 2^\circ$.

DSWE5G - Input data for SWE program. For 5 GHz C-band Cassegrain feed pattern.

DSWE10G - For 10 GHz X-band Cassegrain feed.

Shown in Figure 5 is the scattered pattern from the 140-ft subreflector at 5 GHz.

IV. SOME OF THE THEORY UNDERLYING THE PROGRAMS

A. Determining the Scattered Fields.

1. Problem definition:

Given a perfectly conducting reflecting surface and the magnetic fields incident on this surface, the far field scattered pattern is to be found. The incident electric field does not contribute to the fields scattered from a perfect conductor.

In the coordinate system of Figure 1, the following variables are defined:

- $\bar{E}_s(R, \theta, \phi)$ = far zone scattered electric field = $[E_\theta(\theta, \phi) \hat{\theta} + E_\phi(\theta, \phi) \hat{\phi}] \frac{e^{-jkR}}{R}$, $R \rightarrow \infty$.
- $\hat{R}, \hat{\theta}, \hat{\phi}$ = unit vectors.
- $\bar{H}_i(\rho, \theta, \phi)$ = incident magnetic field.
- \hat{n} = outward unit normal to scattering surface.
- $\hat{\rho}, \hat{\theta}, \hat{\phi}$ = unit vectors.
- $\bar{K}(\rho, \theta, \phi)$ = induced surface current.
- $\bar{\rho}$ = $\rho \hat{\rho}$.
- $\hat{i}, \hat{j}, \hat{k}$ = Cartesian unit vectors.
- k = $2\pi/\lambda$.
- j = $\sqrt{-1}$

140-FT SUBREFLECTOR SCATTERED PATTERN

8/21/81

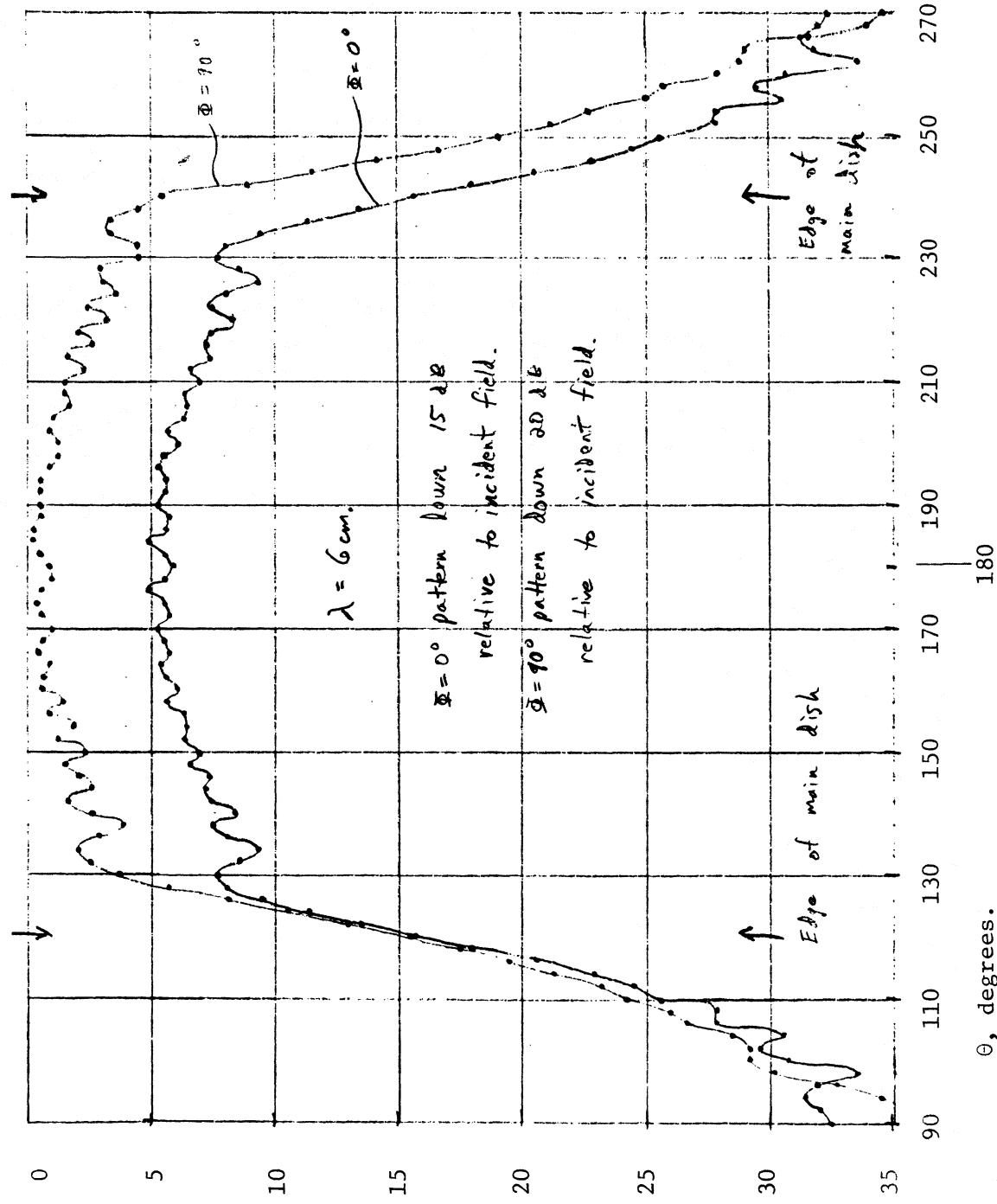


FIGURE 5

The program computes E_θ and E_ϕ in volts. If near field values are needed, then E_θ and E_ϕ can be input to the spherical wave expansion (SWE) program.

To find the total scattered field, the far field incident pattern is added to E_s .

2. Induced currents and the physical optics approximations:

The scattered field, E_s , is determined by taking the 2-dimensional Fourier transform (mapping from $\theta, \phi \rightarrow \theta, \Phi$) of the array of current dipoles. Diffraction effects will be accounted for since the wavefronts of each current dipole are summed. In principle, the accuracy of the technique is limited by the precision with which the induced currents are known.

To determine \bar{K} , the following assumptions are made: On directly illuminated sections of the reflector, the induced currents are those of an optically reflected \bar{H}_i ; on shadowed portions, there are no induced currents. These two assumptions comprise the physical optics approximations.

An expression for \bar{K} can be found by applying the field boundary conditions to the scattering surface. For two arbitrary mediums,

$$\bar{K} = \hat{n} \times (\bar{H}_1 - \bar{H}_2).$$

Since there are no fields beneath the reflector surface, $\bar{H}_2 = 0$ and since \bar{H}_i is optically reflected, $\bar{H}_1 = 2\bar{H}_i$. Thus

$$\bar{K} = 2\hat{n} \times \bar{H}_i \quad (1)$$

Note that \bar{K} gives rise to all fields (incident and scattered) in free space (medium 1).

Ludwig notes that this expression is, in some cases, a poor approximation to the true induced currents. Generally, the calculated currents oscillate about the true currents, making little net contribution to the scattered

fields. Equation (1) can be accurately used for reflectors as small as several wavelengths in diameter.

3. Fields due to the induced currents.

To determine the scattered field at some far field output point (θ, Φ) , the amplitude and phase delay (pathlength) of the radiation from each current dipole is vectorially summed. In principle, the current dipoles are of infinitesimal length; for the numerical integration this, of course, is not true.

For each dS , $d\bar{E}_s$ is proportional to the currents induced on dS . Thus, the radiation amplitude is \bar{K} .

Since the far field of \bar{E}_s is desired, it is instructive to view the output grid, over which the scattered pattern is calculated, as lying on a sphere of infinite radius centered on the source phase center. If the source fields travel directly to the output grid without striking a reflector, the pathlength covered is defined to be zero. When intercepted first by a reflector, the pathlength is defined as the additional distance travelled due to the reflection. This additional distance is that covered by the source fields to the reflector minus the extent to which the incident fields have already travelled in the direction of the output point (extent of colinearity). If R and vector $\hat{\rho}$ are collinear, the pathlength is, by definition, zero. A measure of the colinearity of two vectors is the cosine of the angle between them (dot product). Thus, the pathlength is

$$\gamma = \rho(1 - \hat{\rho} \cdot \hat{R}).$$

To evaluate $\hat{\rho} \cdot \hat{R}$, define the following:

$$x = \rho \sin \theta \cos \phi$$

$$X = R \sin \theta \cos \Phi$$

$$y = \rho \sin \theta \sin \phi$$

$$Y = R \sin \theta \sin \Phi$$

$$z = \rho \cos \theta$$

$$Z = R \cos \Phi$$

and

$$\hat{\rho} = (\hat{x}\hat{i} + \hat{y}\hat{j} + \hat{z}\hat{k})/\rho, \quad \hat{R} = (\hat{X}\hat{i} + \hat{Y}\hat{j} + \hat{Z}\hat{k})/R.$$

The dot product is then

$$\hat{\rho} \cdot \hat{R} = [\sin \theta \sin \theta (\cos \phi \cos \Phi + \sin \phi \sin \Phi) + \cos \theta \cos \theta].$$

The pathlength must be normalized to the wavelength. For an outward travelling spherical wave, the phase delay from source phase center to output point is

$$\text{Phase delay} = e^{-jk\gamma}.$$

At this point, our knowledge of the scattered fields can be summarized by

$$\bar{E}_s \propto \frac{e^{-jkR}}{R} \int_s \bar{K} e^{-jk\gamma} ds,$$

where \bar{E}_s and \bar{K} are expressed in volts/meter and amperes/meter, respectively.

The variable pathlength is neglected in computing the spacial attenuation $(1/R)$ of \bar{E}_s , since R goes to infinity. Therefore,

$$E_{\theta}^{\hat{\theta}} + E_{\phi}^{\hat{\phi}} \propto \int_s \bar{K} e^{-jk\gamma} ds. \quad (2)$$

To get the two sides of equation (2) compatible, we must convert the current-distance of the right side to voltage. Recalling the phase quadrature of associated currents and voltages, the constant of proportionality is $-2\pi j Z_0/\lambda$, where Z_0 = the impedance of free space. The usual expression for this constant is $-j\omega\mu_0$, where μ_0 is the magnetic permeability of free space and $\omega = 2\pi$ times the frequency. Define

$$\bar{I}(\theta, \phi) = [E_{\theta}^{\hat{\theta}} + E_{\phi}^{\hat{\phi}}]/-j\omega\mu_0 \quad (\text{ampere-meters}).$$

Then

$$\bar{I} = \frac{1}{4\pi} \int_s \bar{K} e^{-jk\gamma} ds.$$

The $\frac{1}{4\pi}$ is a normalization constant which arises from integrating over a solid angle. A sphere has 4π square radians (steradians).

Since \bar{H}_i will often be evaluated within the source's near field, it will have a relatively strong radial (ρ) dependence. However, this dependence is predominantly of the form $1/\rho$. (The $e^{-jk\rho}$ phase term has already been included in the expression for γ .) Define:

$$\bar{H}_i = \frac{1}{\rho} [H_\rho \hat{\rho} + H_\theta \hat{\theta} + H_\phi \hat{\phi}],$$

where H_ρ , H_θ and H_ϕ are complex values, thus allowing for phase deviating from the $e^{-jk\rho}$ form. Since H_ρ , H_θ and H_ϕ are slowly varying with ρ , the numerical integration is easier.

Finally, some results from differential geometry concerning normals to surfaces must be considered. We have defined $\bar{\rho}$ as $\bar{\rho} = \bar{\rho}(\theta, \phi)$; specifying θ and ϕ describes the surface. Figure 6 shows the vectors $\bar{\rho}$, $\frac{\partial \bar{\rho}}{\partial \phi}$ and $\frac{\partial \bar{\rho}}{\partial \theta}$. The latter two are tangent to the surface at (θ, ϕ) and are normal to each other. As such, their cross product defines a normal to the surface at (θ, ϕ) . After some staring at Figure 6, one can see that $\frac{\partial \bar{\rho}}{\partial \theta}$ and $\frac{\partial \bar{\rho}}{\partial \phi}$ are the resultants of vectors in the directions $\hat{\rho}$ and $\hat{\theta}$ and the directions $\hat{\rho}$ and $\hat{\phi}$, respectively.

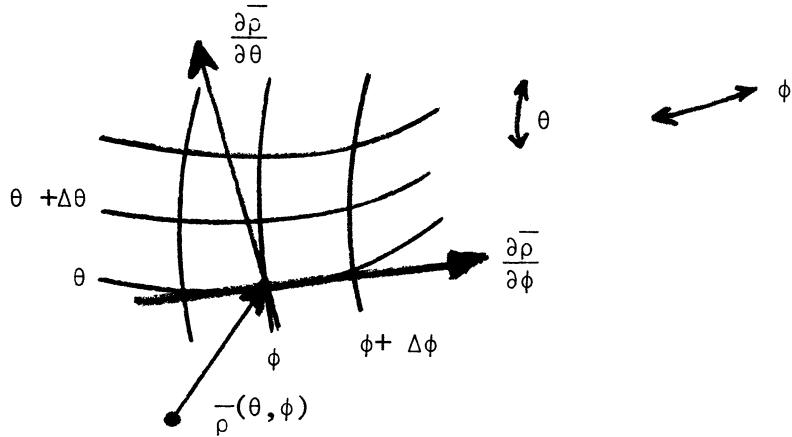


FIGURE 6

Evaluating these derivatives (Figure 1 is helpful), we find:

$$\frac{\partial \bar{\rho}}{\partial \theta} = \left(\frac{\partial \rho}{\partial \theta} \right) \hat{\rho} + \rho \left(\frac{\partial \hat{\rho}}{\partial \theta} \right) = \frac{\partial \rho}{\partial \theta} \hat{\rho} + \rho \hat{\theta}$$

$$\frac{\partial \bar{\rho}}{\partial \phi} = \left(\frac{\partial \rho}{\partial \phi} \right) \hat{\rho} + \rho \left(\frac{\partial \hat{\rho}}{\partial \phi} \right) = \frac{\partial \rho}{\partial \phi} \hat{\rho} + \rho \sin \theta \hat{\phi}$$

and

$$\frac{\partial \bar{\rho}}{\partial \phi} \times \frac{\partial \bar{\rho}}{\partial \theta} = \rho \frac{\partial \rho}{\partial \phi} \hat{\phi} + \rho \sin \theta \frac{\partial \rho}{\partial \theta} \hat{\theta} - \rho^2 \sin \theta \hat{\rho}.$$

The unit normal, \hat{n} , can be found from the relationship

$$\hat{n} dS = \frac{\partial \bar{\rho}}{\partial \phi} \times \frac{\partial \bar{\rho}}{\partial \theta} d\theta d\phi$$

where $dS = (\rho \sin \theta d\phi) (\rho d\theta) = \rho^2 \sin \theta d\phi d\theta$. Substituting for $\hat{n} dS \times \bar{H}_1$ in the equation for \bar{I} , we obtain

$$\bar{I} = \frac{1}{2\pi} \int_s \bar{F}(\theta, \phi) e^{-jk\gamma} d\theta d\phi$$

where

$$\begin{aligned} \bar{F} = & \left(\frac{\partial \rho}{\partial \theta} \sin \theta H_\phi - \frac{\partial \rho}{\partial \phi} H_\theta \right) \hat{\rho} + \left(\frac{\partial \rho}{\partial \phi} H_\rho + \rho \sin \theta H_\phi \right) \hat{\theta} \\ & + \left(-\rho \sin \theta H_\theta - \sin \theta \frac{\partial \rho}{\partial \theta} H_\rho \right) \hat{\phi}. \end{aligned}$$

This is the form in which the radiation integral is evaluated.

B. Technique of Numerical Integration.

1. Form of integral:

Dropping all constants, the form of the integral over a given integration grid and for the output point θ_J, ϕ_K is

$$I(\theta_J, \phi_K) = \int_{\theta_1}^{\theta_M} \int_{\phi_1}^{\phi_N} F(\theta, \phi) \exp[jk\gamma(\theta, \phi, \theta_J, \phi_K)] d\theta d\phi$$

Define $\Delta\Omega_{mn}$ as the piece of solid angle bounded at its four corners by θ_m , θ_{m+1} , ϕ_m , and ϕ_{m+1} . Also, define

$$\begin{aligned}\Delta\theta_m &= \theta_{m+1} - \theta_m \\ \Delta\phi_n &= \phi_{n+1} - \phi_n\end{aligned}$$

$$\rho_{mn}, F_{mn}, \text{etc.} = \rho(\theta_m, \phi_n), F(\theta_m, \phi_n), \text{etc.}$$

The strong sinusoidal variation of the phase delay creates the need for rapid sampling of the scattering surface. Consider a $\Delta\Omega_{mn}$ with physical dimensions on the order of a wavelength; $\rho_{mn} \Delta\theta_m \sim \lambda$ and $\rho_{mn} \sin \theta_m \Delta\phi_n \sim \lambda$. (The dimension of one wavelength is chosen because electromagnetic fields rarely change abruptly over such distances.) For this situation the phase delay can vary up to one full cycle as $\Delta\Omega_{mn}$ is traversed. To successfully apply an integration technique such as Simpson's rule would require further subdivision of the scattering surface, thus creating a monstrous CPU time. Note that because γ is a nonlinear function of θ and ϕ , the fast Fourier transform cannot be applied.

2. The integration technique:

The sampling hassle can be alleviated through use of a linear representation of F and γ (especially) over each $\Delta\Omega_{mn}$. (The procedure to be described below was developed by Ludwig and is shown on page 13 of his JPL report.) Approximate F and γ by

$$\begin{aligned}F(\theta, \phi) &= a_{mn} + b_{mn} (\theta - \theta_m) + c_{mn} (\phi - \phi_n) \\ \gamma(\theta, \phi) &= \alpha_{mn} + \beta_{mn} (\theta - \theta_m) + \xi_{mn} (\phi - \phi_n)\end{aligned}$$

where θ and ϕ assume the corner values of $\Delta\Omega_{mn}$. Applying a least squares plane fit to F and γ at the corners of $\Delta\Omega_{mn}$, the following normal equations are obtained for F :

$$\Sigma F = a_{mn} \sum 1 + b_{mn} \sum (\theta - \theta_m) + c_{mn} \sum (\phi - \phi_n)$$

$$\Sigma F(\theta - \theta_m) = a_{mn} \sum (\theta - \theta_m) + b_{mn} \sum (\theta - \theta_m)^2 + c_{mn} \sum (\phi - \phi_n)(\theta - \theta_m)$$

$$\Sigma F(\phi - \phi_n) = a_{mn} \sum (\phi - \phi_n) + b_{mn} \sum (\theta - \theta_m)(\phi - \phi_n) + c_{mn} \sum (\phi - \phi_n)^2$$

where $\Sigma = \sum_n^{n+1} \sum_m^{m+1}$

Solving for the coefficients yields the following:

$$a_{mn} = \frac{1}{4} [3 F_{mn} - F_{m+1,n+1} + F_{m+1,n} + F_{m,n+1}]$$

$$b_{mn} = \frac{1}{2\Delta\theta_m} [F_{m+1,n} - F_{mn} + F_{m+1,n+1} - F_{m,n+1}]$$

$$c_{mn} = \frac{1}{2\Delta\phi_n} [F_{m,n+1} - F_{mn} + F_{m+1,n+1} - F_{m+1,n}]$$

The results are the same for γ and its coefficients.

Substituting the approximations for F and γ into equation (2) and integrating over $\Delta\Omega_{mn}$, the scattered field contribution ΔI_{mn} can analytically be determined. The expression for ΔI_{mn} is

$$\Delta I_{mn} = \exp jk\alpha_{mn} \left\{ a_{mn} \left[\frac{e_m - 1}{jk\beta_{mn}} \right] \left[\frac{e_n - 1}{jk\xi_{mn}} \right] + b_{mn} \left[\frac{\Delta\theta_m}{jk\beta_{mn}} e_m - \left(\frac{e_m - 1}{(jk\beta_{mn})^2} \right) \right] \left[\frac{e_n - 1}{jk\xi_{mn}} \right] + c_{mn} \left[\frac{e_m - 1}{jk\beta_{mn}} \right] \left[\frac{\Delta\phi_n}{jk\xi_{mn}} e_n - \left(\frac{e_n - 1}{(jk\xi_{mn})^2} \right) \right] \right\}$$

where

$$e_m = \exp jk\beta_{mn} \Delta\theta_m$$

$$e_n = \exp jk\xi_{mn} \Delta\theta_n$$

For each θ_J , ϕ_K of the output grid, the integration subroutine (FINT) computes I from all θ_m , ϕ_n on the integration grid. At every integration point the subroutine computes the least squares coefficients and evaluates the above expression for ΔI_{mn} . To avoid numerical catastrophe an altered expression of ΔI_{mn} is used for β_{mn} and/or ξ_{mn} near zero.

3. Some comments on the method:

Analytic evaluation of the radiation integral can be thought of as summing the (relatively simple) patterns due to infinitesimal current dipoles. The numerical technique sums the patterns, given by ΔI_{mn} , for surface elements a wavelength on a side.

Ludwig claims that for scattering from a hyperboloid, a $\Delta\Omega_{mn}$ 2/3 square wavelengths in size results in errors 40 dB below the scattered pattern maximum.

C. A Few Details of Spherical Wave Expansions.

1. The SWE representation of an electromagnetic field

Define the following variables:

\bar{E} , $\bar{H}(\rho, \theta, \phi)$	= description everywhere in a source-free region V of an electromagnetic field.
TE_{mn} , TM_{mn}	= transverse electric and transverse magnetic fields used to describe \bar{E} and \bar{H} .
m	= order of azimuthal variation of TE, TM fields.
n	= mode order of TE, TM fields.
\bar{m} , \bar{n}	= spherical wave solutions to Maxwell's equations; define TE and TM fields.
a_{mn} , b_{mn}	= TE and TM expansion coefficients.
$z_n(k\rho)$	= any solution to spherical Bessel (differential) equation.
$h_n^{(2)}(k\rho)$	= spherical Hankel function (a particular z_n).
$P_n^m(\cos \theta)$	= associated Legendre function (solution to a form of the Legendre differential equation).

The SWE is used (in this program) to represent a far field input pattern, \bar{E} and \bar{H} , in both the near and far fields. An expansion in TE and TM spherical waves is used as shown:

$$\bar{E} \text{ or } \bar{H} = \sum_m \sum_n a_{mn} TE_{mn} + b_{mn} TM_{mn}.$$

For \bar{E} , TE_{mn} has no radial components and for \bar{H} , TM_{mn} has no radial components.

As discussed earlier, the SWE program calculates the complex-valued wave coefficients. Input for \bar{H} is the far-zone magnetic field of, for example, a feed; $\bar{E} = 0$. The program reduces the flexibility of the expansion in two ways: one, by requiring that m assume a single integer value (usually $m = 1$) and eliminating the summation over m ; and, two, by discarding the odd solutions of \bar{m}_{mn} and \bar{n}_{mn} , thus requiring that \bar{H} be linearly polarized. (The existence of even and odd solutions has been, for convenience, left unrecognized in the notation.)

The general spherical wave expansion is as follows:

Let a sphere of radius ρ_0 contain all field sources. Then the electromagnetic field in the (source-free) region V that includes all space outside of the sphere of radius ρ_0 is

$$\bar{E}(\rho, \theta, \phi) = -\sum_m \sum_n a_{mn} \bar{m}_{mn} + b_{mn} \bar{n}_{mn}$$

$$\bar{H}(\rho, \theta, \phi) = \frac{k}{j\omega\mu} \sum_m \sum_n a_{mn} \bar{n}_{mn} + b_{mn} \bar{m}_{mn}$$

where $\omega\mu/k$ is Z_0 = free space impedance and where

$$\bar{m}_{mn} = z_n(k\rho) \frac{\sum_m P_n^m (\cos \theta)}{\sin \theta} \frac{\sin m\phi}{\cos m\phi} \hat{\theta}$$

$$- z_n(k\rho) \frac{\partial}{\partial \theta} P_n^m (\cos \theta) \frac{\cos}{\sin} m\phi \hat{\phi}$$

$$\bar{n}_{mn} = n(n+1) \frac{z_n(k\rho)}{k\rho} P_n^m (\cos \theta) \frac{\sin}{\cos} m\phi \hat{\rho}$$

$$+ \frac{1}{k\rho} \frac{\partial}{\partial \rho} [\rho z_n(k\rho)] \frac{\partial}{\partial \theta} P_n^m (\cos \theta) \frac{\sin}{\cos} m\phi \hat{\theta}$$

$$+ \frac{1}{k\rho} \frac{\partial}{\partial \rho} [\rho z_n(k\rho)] \frac{\sum_m P_n^m (\cos \theta)}{\sin \theta} \frac{\cos}{\sin} m\phi \hat{\phi}$$

(The even solutions contain the upper of the two signs and the upper of the sin-cos $m\phi$ pair. Further discussion of the specifics of the SWE program expansion will be delayed until section C.)

Since we are dealing with travelling spherical waves, $z_n(k\rho)$ should describe such. The spherical Bessel function that describes the radial variation of an outward travelling spherical wave is $h_n^{(2)}(k\rho)$, the Hankel function of the second kind.

Note the roles played by the mode indices, m and n ; two indices describe variation in three (coordinate) variables. Theta pattern variation, as described by $P_n^m(\cos \theta)$ is dependent on both mode orders (m is not an exponent).

2. Spherical waves as solutions to Maxwell's equations:

In a source free medium, the curl relationships of Maxwell's equations completely describe an electromagnetic field, \bar{E} and \bar{H} . That is,

$$\bar{\nabla} \times \bar{E} = -\mu_0 \frac{\partial \bar{H}}{\partial t} = -j\omega\mu_0 \bar{H}$$

$$\bar{\nabla} \times \bar{H} = \epsilon_0 \frac{\partial \bar{E}}{\partial t} = j\omega\epsilon_0 \bar{E}$$

where

$$\bar{\nabla} = \frac{\partial}{\partial \rho} \hat{\rho} + \frac{1}{\rho} \frac{\partial}{\partial \theta} \hat{\theta} + \frac{1}{\rho \sin \theta} \frac{\partial}{\partial \phi} \hat{\phi}, \text{ and}$$

where an $e^{j\omega t}$ time variation has been assumed and where $\omega\mu_0 = kz_0$ and $\omega\epsilon_0 = k/z_0$. These equations state that the change with distance in directions perpendicular to \bar{E} (\bar{H}) is proportional to the time rate of change of \bar{H} (\bar{E}). The constants of proportionality, μ_0 and ϵ_0 , specify the extent per unit distance to which free space can transmit energy in magnetic and electric fields, respectively. (μ_0 and ϵ_0 are, as shown above, directly related to the impedance of free space, z_0 .) The negative sign in the first curl equation is a result of the orientation of \bar{E} and \bar{H} in a right-hand coordinate system.

Another interpretation of the curl equations is that the existence of an electric field ensures the existence of a corresponding magnetic field, and vice versa, for non-static fields.

Eliminating each variable of the curl equations yields the following:

$$\bar{\nabla} \times (\bar{\nabla} \times \bar{H}) = k^2 \bar{H} \quad (1)$$

$$\bar{\nabla} \times (\bar{\nabla} \times \bar{E}) = k^2 \bar{E}$$

Using vector identities in rectangular coordinates

$$(\bar{\nabla} = \frac{\partial}{\partial x} \hat{i} + \frac{\partial}{\partial y} \hat{j} + \frac{\partial}{\partial z} \hat{k})$$

and noting that the divergence of \bar{E} and \bar{H} are zero ($\bar{\nabla} \cdot \bar{E}, \bar{\nabla} \cdot \bar{H} = 0$ in region V since no sources exist there), these equations are transformed to

$$\bar{\nabla}^2 \bar{H} + k^2 \bar{H} = 0 \quad (\text{Same for } \bar{E}.) \quad (2)$$

This is the vector wave or vector Helmholtz equation.

∇^2 is often called the Laplacian. The (negative of the) Laplacian of a function is closely related to the difference between the value of the function at some point and the average value of the function at straddling points. That is, $\bar{\nabla}^2$ determines the concavity or lumpiness of a function. Thus, equation (2) states that the concavity of the magnitude or direction of \bar{H} is proportional to \bar{H} .

Since the curl of a vector function yields the magnitude and direction of the rotation, or vorticity, of the vector field, equation (1) can be understood as requiring that the vorticity lines of \bar{E} and \bar{H} must themselves exhibit vorticity. Further, this second order vorticity is proportional in magnitude and direction to E and H . (See Methods of Theoretical Physics, Morse and Feshback; chapter 1.) Note that in a source-free space, equations (1) and (2) are identical.

The vector solutions \bar{m}_{mn} and \bar{n}_{mn} of equation (2) are most easily found by considering some physical attributes of \bar{E} and \bar{H} . Recall that the divergence of \bar{E} and \bar{H} is zero, allowing them to be represented by the curl of a vector potential \bar{A} . (To this can be added the gradient of a scalar function. However, since the curl of the gradient of a function is identically zero, this term is of no use.) If $\bar{H} = \frac{1}{\mu_0} \text{curl } \bar{A}$, then from Maxwell's curl equations,

$$\bar{E} = - \frac{\mu_0}{Z_0} \frac{\partial \bar{A}}{\partial t} = -jk\bar{A}.$$

Our choice of \bar{A} can greatly reduce the complexity of the problem. To ease the pain of applying boundary conditions in spherical coordinates, the direction of \bar{A} is (if possible) chosen to be normal to a boundary surface. This will create fields tangent to the surface boundary. The magnitude of \bar{A} is $\psi(\rho, \theta, \phi)$, a rectangular component of \bar{E} or \bar{H} , times a coordinate scale factor. For spherical coordinates the direction and magnitude of \bar{A} are $\hat{\rho}$ and $\rho\psi$, respectively. Thus, a solution to equation (2) is

$$\bar{H} = \bar{\nabla} \times \rho \bar{\psi} = \bar{m} \quad (3)$$

where $\bar{\nabla}$ is in spherical coordinates. Then, equation (3) has no radial components and, if ψ is a component of \bar{H} , \bar{m} is a TM field.

We must, of course, have radial electromagnetic field components in a general field representation. Another solution to (2) which will supply the needed components is simply the curl of \bar{m} . Including a $\frac{1}{k}$ constant of proportionality, a second solution to (2) is

$$\bar{H} = \frac{1}{k} \bar{\nabla} \times \bar{m} = \bar{n}.$$

In general, \hat{n} has ρ, θ and ϕ components. (For detailed mathematical accounts of the solution to (2), see Stratton, chapter 7, and Morse and Feshback, chapter 13.)

The problem has been reduced to determining ψ . Recall that ψ is a solution to

$$\bar{\nabla}^2 H_p, E_p + k^2 H_p, E_p = 0 , \quad p = x, y, z.$$

where E_p and H_p are functions of ρ, θ and ϕ . $\bar{\nabla}^2$ can then be expressed in spherical coordinates. The resulting differential equation is evaluated by the separation of variables

$$\psi(\rho, \theta, \phi) = \psi_1(\rho) \psi_2(\theta) \psi_3(\phi).$$

The solutions to the three resulting ordinary differential equations are:

$$\psi_1(\rho) = z_n(k\rho) = h_n^{(2)}(k\rho)$$

$$\psi_2(\theta) = P_n^m(\cos \theta)$$

$$\psi_3(\phi) = \frac{\cos m\phi}{\sin m\phi} \quad (\text{See Stratton, chapter 7.})$$

The previously shown form of \bar{m} and \bar{n} can be found from $\text{curl } \hat{\psi\rho}$ and $\frac{1}{k} \text{curl curl } \bar{\psi\rho}$.

3. What are spherical waves?

To understand, physically, what spherical waves are, it will be helpful to be aware of some mathematical characteristics of the spherical Hankel function and the associated Legendre function.

The spherical Hankel function of the second kind is defined as

$$h_n^{(2)}(\chi) = j_n(\chi) - j Y_n(\chi), \quad j = \sqrt{-1}$$

At $\chi = 0$, j_n is finite and Y_n has a pole. As $\chi_1 = k\rho_1 \rightarrow \infty$

$$h_n^{(2)}(\chi_1) = j^{n+1} \frac{e^{-j\chi_1}}{\chi_1}$$

and

$$\frac{1}{\chi_1} \frac{\partial}{\partial \chi} \left[\chi h_n^{(2)}(\chi) \right]_{\chi=\chi_1} = j^n \frac{e^{-j\chi_1}}{\chi_1}$$

These are closely related to the $e^{-jk\rho/\rho}$ dependence of a far field pattern.

The associated Legendre function provides the polar angle field variation. The order of index n runs from 0 to ∞ , and m runs 0 to n . A few modes of the function are shown below. (See Stratton, chapter 7; Tables of Function, Johnke and Emde, pp. 112-113.)

$$P_0(Z) = 1$$

$$P_1(Z) = Z = \cos \theta$$

$$P_1^1(Z) = (1 - Z^2)^{1/2} = \sin \theta$$

$$P_2(Z) = \frac{1}{2} (3Z^2 - 1) = \frac{1}{4} (3 \cos 2\theta + 1)$$

$$P_2^1(Z) = 3(1 - Z^2)^{1/2} Z = \frac{3}{2} \sin 2\theta$$

$$P_2^2(Z) = 3(1 - Z^2) = \frac{3}{2} (1 - \cos 2\theta)$$

The functions $\cos m\phi P_n^m (\cos \theta)$ and $\sin m\phi P_n^m (\cos \theta)$ are periodic on the surface of a unit sphere. For $m > 0$, the functions are zero at the poles. The number of nodal lines parallel to the equator is $n - m$. These are

orthogonally intersected by the 2 m longitudinal nodes. Since the surface is divided into rectangular sections within which the above functions are alternately positive and negative, these functions are referred to as tesseral harmonics of n^{th} degree and m^{th} order.

Evaluating \bar{m} and \bar{n} at $m = 0$ and $n = 1$, the simplest form of \bar{m} and \bar{n} , will provide the vital clue in determining the nature of these spherical wave solutions.

Recall that \bar{E} and \bar{H} are defined everywhere in space V which surrounds but does not include a sphere of sources. Let this sphere be of radius a . An entirely equivalent representation of the sources is to place surface currents on the sphere and remove all sources within the sphere. From the boundary conditions, the surface currents are

$$\bar{K} = \hat{\rho} \times \bar{H}(a, \theta, \phi).$$

From the expressions for \bar{E} and \bar{H} involving the vector potential \bar{A} ,

$$\bar{E} = -jk\bar{A} \quad (\text{for } e^{j\omega t} \text{ time variation})$$

$$\bar{H} = \frac{1}{\mu_0} \text{curl } \bar{A}$$

where $\bar{A} = \bar{m}$ and \bar{n} , we can determine the induced currents. For $m = 0$, $n = 1$,

$$\bar{m}_{01} = h_1^{(2)}(k\rho) \sin \theta \hat{\phi}$$

$$\bar{n}_{01} = \frac{1}{k} \bar{\nabla} \times \bar{m}_{01}$$

$$= 2 \cos \theta \frac{1}{k\rho} h_1(k\rho) \hat{\rho} - \sin \theta \frac{1}{k\rho} \frac{\partial}{\partial \rho} [\rho h_1(k\rho)] \hat{\theta}.$$

First consider the case of $\bar{A} = \bar{m}_{01}$. Through the use of various recurrence relations (see Stratton, chapter 7) and the relation $h_n(z) = -j(-1)^n \left(\frac{d}{ZdZ} \right)^n \left(\frac{e^{iz}}{Z} \right)$, the terms involving Hankel functions can be reduced to algebraic expressions in k_0 and $e^{jk_0 p}$. Then the surface current creating the $m = 0, n = 1$ mode fields is

$$\bar{K}_{01} \propto j \frac{\omega}{c} \sin \theta \left[\frac{1 - (ka)^2 - jka}{(ka)^3} \right] e^{jka} \hat{\phi}$$

(See Morse and Feshback, chapter 13, p. 1867.)

Thus, we have a current oscillating in time parallel to the equator. The current goes to maximum at the equator and goes to zero at the poles and travels in the same direction at all points on the sphere. There is no charge build up as can be seen from the lack of radial \bar{E} components.

A current of such form characterizes the magnetic dipole and \bar{m}_{01} represents the field from the dipole.

If we are to determine the radial components of \bar{E} , we must turn the problem around and set $\bar{A} = \bar{n}_{01}$. Then, $\bar{E} = -jk \bar{n}_{01}$ and has a radial component. The surface current is

$$\bar{K}_{01} \propto \omega \mu_0 \sin \theta \left[\frac{j + ka}{(ka)^2} \right] e^{jka} \hat{\theta}$$

which oscillates between poles, alternately depositing positive and negative charges at the poles. As can be seen from $\hat{\rho} \cdot \bar{E}$, the charge is concentrated at the poles and is 90° out of phase with the current.

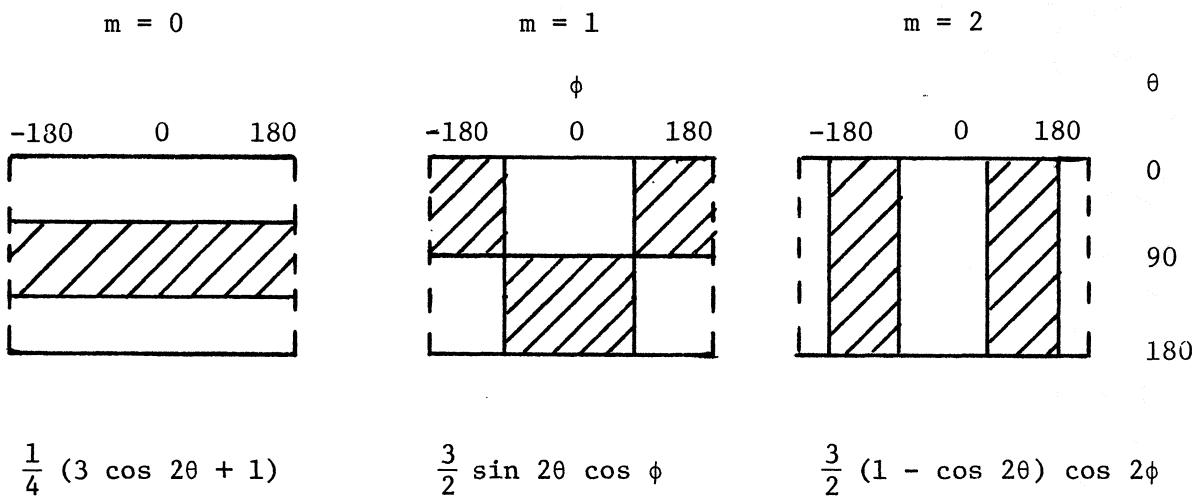
This oscillation of charge is merely a way of describing an electric dipole where \bar{n}_{01} represents the field from this dipole.

For $m = 0, n = 2$, the TE radiation is that of a magnetic quadrupole. The current circulates parallel to the equator going clockwise in one hemisphere

and counterclockwise in the other for half a cycle, then reversing for the next half cycle.

The corresponding TM radiation is that of an electric quadrupole. The current distribution describes charge moving alternately from the equator to the poles and then from the poles to the equator. Thus, free charge of one sign is accumulated at the poles and of the opposite sign at the equator, the signs changing at the next half cycle.

Among others, one important question that remains is the following: What is physically happening as m changes for a given n ? Plotting $\cos m\phi P_n^m(\cos \theta)$ on the unit sphere for various values of m seems to indicate that m determines two properties of the multipoles: the distribution of multipoles within the array and the orientation of the multipole array. As an example, consider the unit sphere plots for $n = 2$ and $m = 0, 1, 2$:



The shaded portions indicate negative value; nodal lines are solid. For $m = 0$, there is an overlap of 2 like (in sign) "monopoles". For $m = 1$ and 2, all "monopoles" (2 of each sign) are distinct but the quadrupole orientations are different.

4. Determining the SWE coefficients.

As mentioned previously, the summation over m in the spherical wave expansion is eliminated; usually an order of azimuthal variation of $m = 1$ is needed. The $m = 1$ corresponds to the situation in which a pattern may be divided into E and H planes. Since only the incident magnetic field contributes to the induced surface currents, only \bar{H}_i need be expanded. The expansion is:

$$\bar{H}_i(\rho, \theta, \phi) = \frac{k}{j\omega\mu} \sum_{n=1}^N a_n \bar{n}_n + b_n \bar{m}_n$$

where $\bar{n}_n \equiv \bar{n}_{mn}$ and is TE and $\bar{m}_n \equiv \bar{m}_{mn}$ and is TM. This expansion will determine \bar{H}_i , where \bar{H}_i is originally a far field pattern, anywhere in the space outside a sphere of radius ρ_0 which encloses the sources of \bar{H}_i .

Ludwig derives expressions for the expansion coefficients for the case in which the involved data are the tangential components of \bar{E} on a sphere of radius $\rho_1 > \rho_0$. The derivation will be outlined only; the mathematical details may be found in the JPL report.

The vector character of the expansion is eliminated (momentarily) by equating the tangential field components, $E_\theta(\theta, \phi)$ and $E_\phi(\theta, \phi)$, to the summation of the corresponding components of the expansion. The azimuthal terms of the summation are removed by taking the ordinary Fourier expansion of E_θ and E_ϕ , leaving a summation independent of ϕ . $E_\theta(\theta, 90^\circ)$ and $E_\phi(\theta, 0^\circ)$ are the usual far field patterns being expanded. The actual inputs to the program are $A_m(\theta)$ and $B_m(\theta)$, which represent the m^{th} Fourier component of the input pattern. (For $m = 1$, $A_m(\theta)$ and $B_m(\theta)$ are simply the pattern values.)

Using the integral (orthogonality) properties of Legendre functions, an integral (over θ) expression is found for the coefficients (JPL report; page 23,

equation 9). The far field value of the Hankel function and its derivative with respect to ρ is incorporated into the coefficients; that is, the coefficients actually computed are:

$$a_n' = a_n j^{n+1} e^{-jk\rho_1} / k\rho_1$$

$$b_n' = b_n j^n e^{-jk\rho_1} / k\rho_1$$

a_n' and b_n' are written on the computer disk where the SCAT program can gain access to them.

V. ACKNOWLEDGEMENTS

I would like to thank Rick Fisher for his endless helpful suggestions and comments and for his ability to keep a straight face in the midst of sometimes ridiculous questions. I would also like to thank Pat Crane and Marc Damashek for their many valuable programming hints, and Carolyn Dunkle for typing this (initially) illegible report.

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- (All but Ludwig's report are available in the Green Bank Library.)

VII.

PROGRAM CODE FOR BOTH PROGRAMS


```

CALL ADJUST(ETPHI)
CALL ADJUST(EPHI)
ETAMP=ETAMP*SCALE
EPAMP=EPAMP*SCALE
WRITE(6,2012)TC,ETAMP,ETPHI,EPAMP,EPHI
CONTINUE
C
STOP
FORMAT(16A4)
1001 FORMAT(10F10.4)      NRAD SCATTERING PROGRAM //5X,18A4)
2001 FORMAT(1H1,*18A4)
2002 FORMAT(1H0,*5X,18A4)
2003 FORMAT(*0,*     PROPAGATION CONSTANT = *E14.8/)
2006 FORMAT(1H0,*     SCATTERED FIELDS FROM GRID *12)
2007 FORMAT(1H0,*     BEGIN INTEGRATION OVER GRID *12)
2009 FORMAT(1H1,*     DIRECT RADIATION FROM INCIDENT FIELDS)
2C10 FORMAT(1H0,*     SUPERPOSITION OF ALL GRID SCATTERED FIELDS AND DIRECT*
2C10 FORMAT(1H0,*     PHASECENTER TRANSLATED BY X=*F10.4*, Y=*F10.4*, Z=*F10.4*/)
2C10 AMPLITUDE VALUES SCALED BY FACTOR OF *E15.*R)
2011 FORMAT(7HO, PHI=F7.2/, /17X,THE THETA,15X,5HE PHI,/, /9H, THETA,2*(20H, VOLTS, PHASE))
2012 FORMAT(9*2,F11.6,F8.2,F12.6,F8.2)
2013 FORMAT(F10.2,F10.0,F10.2,F10.6,F10.2)
END

02040000          SUBROUTINE SURF(I,MMAX,NMAX,F,FT,FP,IEDGE,TEDGE)
02050000          C
02060000          C
02070000          DIMENSION F(36*91),FT(36*91),FP(36*91),TEDGE(91)
02080000          COMMON/GRID/SIT(36),COT(36),SIT( 91),COT( 91),T(36)*PI* 91
02090000          C
02100000          THIS SUB. PROVIDES MAIN WITH RHO,DRHO/DTHETA AND DRHO/DPHI
02160000          C
02170000          INPUT RHO0,# INTEGRATION GRIDS.# THETA INTEG. GRIDS
02180000          C
02190000          C
02200000          C
02210000          C
02230000          C
02240000          C
02250000          C
02260000          C
02270000          C
READ5*1001ISURF
DT=(T(12)-T(11))/01745329
DP=4(P(12)-P(11))/01745329
02280000          C
02300000          WRITE(6,2002)I,ISURF
02310000          C
02320000          FORMAT(1H0, *FOR INTEGRATION GRID #*13*, ISURF = *I4,*)
02330000          C
02340000          WRITE(6,2004)
02350000          C
FORMAT1H * M N   F   FT   FP *
02360000          C
02370000          C
ISURF CONSIDERS FOLLOWING CASES:
00180000          C
00180000          C
1SURF<0:NOT=-2: F,FT,FP FOUND FROM EQUATIONS INSERTED BELOW.
00180000          C
1SURF=F FROM EQUATION;FT,FP NUMERICALLY DETERMINED.
00180000          C
1SURF>0: F IS FUNCTION OF PHI ONLY AND IS IN TABULAR FORM.
00180000          C
1SURF<0: F IS FUNCTION OF PHI AND IS IN TABULAR
00180000          C
FORM*. FT,FP NUMERICALLY DETERMINED.
00180000          C
00190000          C
IF(1SURF)>20*GO TO 30
20 CONTINUE
C
CALCULATE F,FT,FP FROM EQUATION TO BE INSERTED HERE
C
DO 22 N=1,NMAX
DO 22 M=1,MMAX
DEN=COT(M)
F(M,N)=0.5 /DEN
IF(1SURF .EQ. -2)GO TO 22
FT(M,N)=F(M,N)/COT(M)*SIT(M)
FP(M,N)=0.0
CONTINUE
IF(1SURF .EQ. -2)GO TO 100
GO TO 110
CONTINUE
C
CASE OF ZERO PHI VARIATION
C
C
READ(5,1002)I,E(M+1),M=1,MMAX

```

```

MSTOP=0          00410000 32  CONTINUE
DO 49 N=1,NMAX  00420000 31  CONTINUE
IF(MSTOP .EQ. 1)GO TO 45
C   IF RHO=999.9999, EDGE HAS JUST BEEN PASSED; VALUES BEYOND EDGE
C   WERE ASSIGNED ARBITRARILY.
C   IF(F(M,1) .NE. 999.9999)GO TO 49
C   EXTRAPOLATE FOR RHO VALUES PAST REFLECTOR EDGE
C
C   MSTOP=1          00430000  C   USE BACK & FORWARD DIFFERENCES TO COMPUTE FT & FP
IF(M .GE. 4)F(M,1)=2.*F(M-1,1)-2.*F(M-2,1)+.5*F(M-3,1)
IF(M .EQ. 3)F(M,1)=2.*F(M-1,1)-F(M-2,1)
IF(M .EQ. 2)F(M,1)=F(M-1,1)
CONTINUE          00432000  C   CALL DIFF(I-1, MMAX,NMAX,C,
                           FT,DT,DP)
                           CALL DIFF(I,  MMAX,NMAX,F,
                           FP,DT,DP)
CONTINUE          00433000  C
                           CALL DIFF(I+1, MMAX,NMAX,F,
                           FP,DT,DP)
CONTINUE          00434000  C
                           CALL DIFF(I+1, MMAX,NMAX,F,
                           FP,DT,DP)
CONTINUE          00435000  C
                           CALL DIFF(I+1, MMAX,NMAX,F,
                           FP,DT,DP)
CONTINUE          00436000  C
                           CALL DIFF(I+1, MMAX,NMAX,F,
                           FP,DT,DP)
CONTINUE          00437000  C   WRITE SOME SURFACE PARAMETERS
IF(M .GE. 4)F(M,1)=2.*F(M-1,1)-2.*F(M-2,1)+.5*F(M-3,1)
IF(M .EQ. 3)F(M,1)=2.*F(M-1,1)-F(M-2,1)
IF(M .EQ. 2)F(M,1)=F(M-1,1)
CONTINUE          00440000  C   DO 120 N=1,NMAX+15
                           DO 120 M=1,MMAX
                           WRITE(6,2003) M*N,F(M,N),FP(M,N)
CONTINUE          00441000  C   2003  FORMAT(6,IIX,F10.6)
                           RETURN
                           END
C
C   CALCULATE FT USING FORWARD & BACKWARD DIFFERENCES
C
C   CALL DIFF(I+1, MMAX,1,F*          00450000
                           FT,DT,DP)
DO 43 N=1,MMAX  00451000
C   DRHO/DPHI IDENTICALLY ZERO
C   F(M,1)=0.          00452000
C
C   ASSIGNMENTS FOR N=2 TO NMAX
C
DO 44 N=2,NMAX  00453000
DO 44 N=1,NMAX  00454000
F(M,N)=F(M,1)    00455000
FT(M,N)=FT(M,1)  00456000
FP(M,N)=FP(M,1)  00457000
CONTINUE          00458000
C
C   VARIATION IN THETA AND PHI
C
CONTINUE          00459000
DO 31 N=1,NMAX  00460000
MSTOP=C          00461000
CONTINUE          00462000
DO 32 M=1,MMAX  00463000
IF(MSTOP .EQ. 1)GO TO 35
C   CHECK FOR RHO=999.9999
C   IF(F(M,N) .NE. 999.9999)GO TO 32
C
C   EXTRAPOLATE RHO VALUES PAST REFLECTOR EDGE TO END OF INTEG.GRIDDDOB12000
C
MSTOP=1          00464000
IF(M .GE. 4)F(M,N)=2.5*F(M-1,N)-2.*F(M-2,N)+.5*F(M-3,N)
IF(M .EQ. 3)F(M,N)=2.*F(M-1,N)-F(M-2,N)
IF(M .EQ. 2)F(M,N)=F(M-1,N)
CONTINUE          00465000
35

```

```

CCCCCCCCCCCC
CCCCCCCCCCCC
C
C SUBROUTINE EDGEED(NMAX,EDGE)
C   THIS SU2. NUMERICALLY INTEGRATES THE RADIATION INTEGRAL
C
C   DIMENSION X(36),Y( 91), R(36,91),TDEGE( 91),
C   COMPLEX F(36,9,31)*STO( 3),T1,T2,T3,A*B*C,F23*F14*SUM*TOT
C   COMPLEX FVTEMP(2*3),AT,DELTA
C
C   SOME VARIABLE DEFINITIONS:
C
C   X*Y = THETA*PHI
C   R = PATHLENGTH TERM
C   F = PATHVECTOR RELATED TO SURFACE CURRENTS
C
C   STOT = 3 COMPONENTS OF EVALUATED INTEGRAL (RETURNED TO MAIN)
C   TDEGE = 3 THETA VALUES THAT SPECIFY REFLECTOR EDGE AS A FUNCTION
C   OF PHI
C
C   IEDGE = INTEGER PARAMETER THAT SIGNIFIES PRESENCE OF EDGE AND
C   SPECIFIES HOW EDGE IS OBTAINED.
C
C   RTTEMP*FVTEMP = TEMPORARY STORAGE FOR R*F
C   NCE = PARAMETER THAT INDICATES FOR EACH PHI, WHEN AN EDGE THETA
C   HAS BEEN REACHED.
C
C   TEANG = AVERAGE OF EDGE THETA VALUES CORRESPONDING TO PHI AND
C   PHI+DELTA(PHI). THIS COMPROMISE RESULTS IN A SLIGHTLY
C   DISCONTINUOUS REFLECTOR EDGE.
C
C   NXPL=NMAX+1
C   MXPL=MMAX+1
C
C DO 111 N=1,NMAX
C   IF((P(N)/DTR .GE. 315.) .OR. (P(N)/DTR .LE. 45.))GO TO 10
C   IF((P(N)/DTR .GE. 45.) .AND. (P(N)/DTR .LE. 135.))GO TO 20
C   IF((P(N)/DTR .GE. 135.) .AND. (P(N)/DTR .LE. 225.))GO TO 30
C   IF((P(N)/DTR .GE. 225.) .AND. (P(N)/DTR .LE. 315.))GO TO 40
C
C 10  TDEGE(N)=ATAN(0.069927/ABS(COP(N)))
C   GO TO 111
C 20  TDEGE(N)=ATAN(0.069927/ABS(SIN(N)))
C   GO TO 111
C 30  TDEGE(N)=ATAN(0.069927/ABS(COP(N)))
C   GO TO 111
C
C 40  TDEGE(N)=ATAN(0.069927/ABS(SIN(N)))
C   TDEGE(N)=TDEGE(N)/DTR
C   RETURN
C   END
C
C 111
C
C   IF(TDEGE(N)=TDEGE(N)) GO TO 79
C
C   DO 201 N=1,NMAX
C     STOT(L)=10.0,0,0
C
C     ENTER PHI LOOP
C
C     DO 200 N=1,NMAX
C       NCE=0
C       DY=0.5*(Y(N+1)-Y(N))
C
C     11  ENTER THETA LOOP (WITHIN PHI LOOP)
C
C     DO 201 N=1,MMAX
C       CHECK IF EDGE TO BE ENCOUNTERED ON THIS INTEG GRID
C
C       IF(IEDGE .EQ. 1)GO TO 79
C
C       IF EDGE PRESENT,COMPARE EACH THETA (X(M)) WITH THE EDGE VALUE
C       OF THE TA FOR THAT PHI. AT FIRST HINT OF EDGE, INTEGRATION
C       PROCEDURE IS ALTERED.
C
C       X(M) MUST NOT BE > TDEGE(N) FOR ANY N
C       IF((TDEGE(N) .GE. X(M+1)) .AND. (TDEGE(N+1) .GE. X(M+1)))
C       EGO TO 79
C       IF((X(M) .LE. TDEGE(N)) .AND. (TDEGE(N) .LT. X(M+1)))
C       EGO TO 30
C       IF((X(M) .LE. TDEGE(N+1)) .AND. (TDEGE(N+1) .LT.
C       E(X(M+1))) GO TO 30
C
C 200
C 201
C 79
C
C 01030000  SUBROUTINE FINIT(X,Y,F,R,MMAX,NMAX,STOT,I,J,KK,PC,IEDGE*EDGE)
C 01040000
C 01050000
C 01060000
C 01070000
C 01080000
C 01090000
C 01092000
C 01100000
C 01101000
C 01102000
C 01125000
C 01130000
C 01131000
C 01140000
C 01150000
C 01150100
C 01150200
C 01150300
C 01150400
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150      READ(5,1002)(P(N,I),N=1,NNN)
150      CONTINUE
C      PRINT OUT GRID DATA
C      WRITE(6,2001)
C      WRITE(6,2006)(J,TT(J),J=1,JMAX)
C      WRITE(6,2002)
C      WRITE(6,2006)PP(K),K=1,KMAX)
C      WRITE(6,2004)NG1
DO 155 I=1,NG1
C      WRITE(6,2005)I
      MM=M(M(I))
      NNN=N(N(I))
C      WRITE(6,2006)(M,T(M,I),M=1,MMM)
C      WRITE(6,2002)
C      WRITE(6,2006)(N,P(N,I),N=1,NNN)
C      COMPUTE GRID2 TABLE
C      COMPUTE GRID1 TABLES
C      DTR=0.017453293
DO 200 J=1,JMAX
      TT(J)=TT(J)*DTR
      SITT(J)=SIN(TT(J))
      COTT(J)=COS(SITT(J))
CONTINUE
DO 210 K=1,KMAX
      PP(K)=PP(K)*DTR
      COP(K)=COS(PP(K))
      SIP(K)=SIN(PP(K))
CONTINUE
C      COMPUTE GRID1 TABLES
C      I=0
C      FOR NG1 >1 RE-ENTER SUB. HERE
C      ENTRY RESET(1,NMAX,NMAX)
I=I+1
      MMAX=M(M(I))
      DO 400 M=1,MMAX
      TG1(M)=DTR*T(M,I)
      SIT(M)=SIN(TG1(M))
      COT(M)=COS(TG1(M))
CONTINUE
      NMAX=N(N(I))
      DO 410 N=1,NMAX
      PG1(N)=DTR*P(N,I)
      SIP(N)=SIN(PG1(N))
      COP(N)=COS(PG1(N))
CONTINUE
410      RETURN
C      1001 FORMAT(1S,2F10.2)
C      1002 FORMAT(8F10.2)
2001      FORMAT(1H,* THE FOLLOWING OUTPUT GRID HAS BEEN ESTABLISHED *//,
      * ALL ANGLES IN DEGREES *//)
2002      FORMAT(1//)

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SUBROUTINE FIELDS
C THIS SUB. FINDS THE MAGNETIC FIELD ON A SURFACE S BY
C EVALUATING A SPHERICAL WAVE EXPANSION. THE WAVE COEFFICIENTS
C AS CALCULATED BY ANOTHER PROGRAM ARE OBTAINED VIA A DISK
C DATA SET
C COMMON/GRID1/SIT(36),COT(181),SIP(91),COP(91),T(36),P(91)
C COMMON/GRID2/SIT(181),COT(181),SIP(5),COP(5),T(181),PP(5)
C DIMENSION FT(81),G(81)
C DIMENSION NAME(18)
C DIMENSION A(60,2),B(80,2)

C SOME VARIABLE DEFINITIONS:
C LMAX = MAXIMUM MODE ORDER
C F,G = VECTORS OF LEGENDRE POLYNOMIAL WEIGHTS EVALUATED FOR EACH
C MODE COMPONENT. THE ARGUMENT OF THE POLY IS THE COS OF
C A POLAR ANGLE.
C A,B = TE & TM WAVE COEFFICIENTS

C INPUT WAVE COEFFICIENTS FROM COMPUTER DISK
C
C READ(9,1001)NAME
C WRITE(6,2001)NAME
C READ(9,1002)LMAX,MCOMP
C READ(9,1001)NAME
C READ(9,1003)J,A(J,1),A(J,2),R(J,1),B(J,2),J=1,LMAX)
C FMC=MCOMP
C WRITE(6,2003)NAME
C WRITE(6,2004)J,A(J,1),A(J,2),B(J,1),B(J,2),J=1,LMAX)
C RETURN

CCCCC
CCCCC ENTRY POINT FOR MAGNETIC FIELDS AT FINITE R
C ENTRY FIELD1,NMAX,HR,HT,HP,R
C DIMENSION R(36,91)
C COMPLEX HR(36,91),HT(36,91),HP(36,91)

C MORE VARIABLES:
C R = RHO
C HR,HT,HP = INCIDENT MAGNETIC FIELD COMPONENTS RETURNED
C TO FIELD1
C MMAX,NMAX = # OF THETA,PHI POINTS ON INTEG GRID
C
C IENT=1
C M=0
C ENTER THETA LOOP
C
C N=M+
C SN=SIT(M)
C Z=COT(M)
C TOUT=T(M)
C
C GO TO 99
C ENTRY POINT FOR ELECTRIC FIELDS AT INFINITE R
C ENTRY FIELD2(JD,KC,EITD,EPPD)
C COMPLEX EITD,EPPD
C
C STILL MORE VARIABLES:
C JD,KC = POINT ON OUTPUT GRID
C EITD,EPPD = FIELD VALUES AT JD,KD (RETURNED TO MAIN)
C
C IENT=2
C SN=SIT(JD)
C Z=COT(JD)
C TOUT=T(JD)
C
C FORM LEGENDRE FUNCTION TABLES FOR M-TH THETA VALUE
C
C IF(ABS(SN)-.00001)>200*100*100
C DO 105 N=1,MCOMP
C   F(N)=0
C 105
C
C OBTAIN F & G VECTORS FOR A GIVEN ORDER OF AZIMUTHAL VARIATION
C
C NC=LMAX+1
C CALL LEGENDRE(NC,MCOMP,Z,F)
C DO 110 N=1,LMAX
C   T1=N-MCOMP+1
C   T2=N+1
C   G(N)=T1*F(N+1)-T2*Z*F(N)
C   G(N)=G(N)/SN
C   DO 115 N=1,LMAX
C     F(N)=F(N)/SN
C 115
C
C FOR NEAR FIELD CASE, MUST CONSIDER RADIAL DEPENDENCE OF FIELDS
C AS DEFINED BY THE SPHERICAL HANKEL FUNCTION.
C
C GO TO 1300,5001,IENT
C
C SPECIAL EQUATIONS FOR THETA=0,180 DEG
C
C 01610000 C
C 01620000 C
C 01630000 C
C 01640000 C
C 01650000 C
C 01660000 C
C 01670000 C
C 01680000 C
C 01690000 C
C 01700000 C
C 01710000 C
C 01720000 C
C 01730000 C
C 01740000 C
C 01750000 C
C 01760000 C
C 01770000 C
C 01780000 C
C 01790000 C
C
C GO TO (300,5001,IENT
C
C FOR EACH PHI THE HANKEL FUNCTIONS ARE EVALUATED AT THF
C
C 01220000 C
C 01230000 C
C 01240000 C
C 01250000 C
C 01260000 C
C 01270000 C
C 01280000 C
C 01290000 C
C 01300000 C
C 01310000 C
C 01320000 C
C 01330000 C
C 01340000 C
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C 02300000 C
C 02310000 C
C 02320000 C
C 02330000 C
C 02340000 C
C 02350000 C
C 02360000 C
C 02370000 C

```

```

C CORRESPONDING VALUE OF R FOR EACH MODE COMPONENT.
C DO 450 N=1,NMAX
30U HRR=0
HRI=0
HTR=0
HTI=0
HPR=0
HP1=0
CALL SPHANK(1,R(M,N),SOI,SII)
DO 400 L=1,LMAX
FL=L
FN=(FL+1.0)/R(M,N)
NC=L+1
CALL SPHANK(NC,R(M,N),SIR,SII)

C TIME FOR SOME GLOP
C DETERMINE THE FIELD AT EACH THETA,PHI ON THE INTEG GRID BY
C FIRST OBTAINING THE PRODUCT OF COEFFICIENTS AND THE SPHANK
C VALUES FOR EACH MODE COMPONENT.
C
FL=R=A(L,1)*SIR+A(L,2)*SII
F1I=A(L,1)*SII-A(L,2)*SIR
F2R=N*(A(L,1)*SOI+A(L,2)*SOR)
F3R=B(L,1)*SOR+B(L,2)*SOI
F3I=-B(L,1)*SOI-B(L,2)*SOR

C THEN MULTIPLY BY THE CORRESPONDING MODE LEGENDRE POLY VALUES.
C FOR EACH THETA,PHI THE ABOVE VALUES ARE SUMMED OVER ALL MODE
C COMPONENTS.
C
HRR=HRR+F2R*F(L)*FL
HRI=F1I*(L)*FL
FLI=FL*FL
HTR=HTR*G(L)*FL
HTI=HTI*G(L)*FL
HPR=F1R*F(L)-F2R*F(L)*F3I*FL
HP1=HP1-F1I*FL-F2I*(L-F3I*FL)
SOR=SIR
SOI=SII
CONTINUE
400
C FINALLY, THE PHI EXPANSION TERM IS INTRODUCED, COMPLETING THE
C EXPANSION.
C
HRR=HRR*SIN(COS(FMC*P(N)))
HRI=HRI*SIN(COS(FMC*P(N)))
HR(M,N)=CMPLX(HRR,HRI)
HTR=HTR*COS(FMC*P(N))
HTI=HTI*COS(FMC*P(N))
HT(M,N)=CMPLX(HTR,HTI)
HPR=HPR*SIN(FMC*P(N))
HPI=HPI*SIN(FMC*P(N))
HP(M,N)=CMPLX(HPR,HPI)
WRITE(16,2005)M,N,HR(M,N),HT(M,N),HP(M,N)
2005 FORMAT(1H 'H-FIELDS',215,6F12.5)
CONTINUE
450

```

```

SUBROUTINE PATHL(RHO,J,K,NMAX,NMAX,GAM)
C
C THIS SUB. COMPUTES PATH LENGTH FUNCTION GAMMA
C FOR EACH POINT ON THE INTEC GRID AND IN THE DIRECTION OF AN
C OUTPUT GRID POINT.
C I.E. • THE PATHLENGTH IS A FUNCTION OF 4 VARIABLES.
C
C DIMENSION RHO(36,91),GAM(36,91)
COMMON/GRID1/SIT(36),COT(36),SIP(91),CP(91),T(36),P(91)
COMMON/GRID2/SITT(181),COTT(181),SIPP(51),COPP(51),TP(181),PP(51)
C
C VARIABLES:
C J,K = OUTPUT GRID POINT
C GAM = PATHLENGTH TERM (RETURNED TO MAIN)
C
DO 10 M=1,NMAX
  T1=SIT(M)*SIT(J)*COPP(K)
  T2=SIT(M)*SIT(J)*SIPP(K)
  T3=COT(M)*COTT(J)-1.0
  DO 10 N=1,NMAX
    GAM(N)=RHO(M,N)*(T1+COP(N)+T2*SIP(N)+T3)
    CONTINUE
    RETURN
  END
10

```

10

```

03380000  SUBROUTINE VECTOR(X,Y,AMP,PHI)
03390000  C
03400000  C THIS SUB. CONVERTS COMPLEX VALUES TO POLAR FORM
03410000  C AMP & PHI RETURNED
03420000  C
03430000  C=C=0
03440000  IF(X<100.*200.*300.
03450000  IF(Y<10.,120.,120.
03460000  100
03470000  110
03480000  120
C=360.0
PHI=ATAN(Y/X)*57.295777951*180.-C
03490000  200
03500000  210
03510000  220
03520000  230
03530000  240
03540000  250
03550000  260
03560000  270
03570000  280
03580000  290
03590000  300
03600000  310
03610000  320
03620000  330
03630000  340
03640000  350
03650000  360
03660000  370
03670000  380
03680000  390
03690000  400
03700000  410
03710000  420
03720000  430
03730000  440
03740000  450
03750000  460
03760000  470
03770000  480
03780000  490
03790000  500
03800000  510
03810000  520
03820000  530
03830000  540
03840000  550

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04370000
04380000
04390000
04400000
04410000
04420000
04430000
04440000
04450000
04460000
04470000
04480000

03850000
03860000
03870000
03880000
03890000
03900000
03910000
03920000
03930000
03940000
03950000
03960000
03970000
03980000
03990000
04000000
04010000
04020000
04030000
04040000
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04100000
04110000
04120000
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04200000
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04250000
04260000
04270000
04280000
04290000
04300000
04310000
04320000
04330000
04340000
04350000
04360000

C          SUBROUTINE SPHANK(NR,X,Y)
C
C          THIS SUB. COMPUTES VALUE OF SPHERICAL HANKEL FUNCTION TIMES
C          THE FACTOR (1-J)*(N1+RHO*EXP(J*PHI)). THIS FACTOR IS THE
C          RECIPROCAL OF THE FAR FIELD SPHERICAL HANKEL FUNCTION. THE
C          FAR FIELD DEPENDENCE IS BUILT IN TO THE SPHERICAL WAVE
C          COEFFICIENTS AND MUST BE REMOVED FOR NEAR FIELD CASES.
C          THE HANKEL FUNCTION PROVIDES THE NEAR FIELD RADIAL VARIATION
C          OF THE PATTERN.
C
C          DOUBLE PRECISION TERM,AI,AR
C
C          VARIABLES:
C
C          N = MODE ORDER + 1
C          R = ARGUMENT OF FUNCTION = PHASE CONSTANT * RHO
C          X,Y = REAL & IMAG COMPUTED VALUES (RE TURNED TO FIELD)
C
C          AR=0
C          AI=0
C          PI=3.1415927
C          K=0
C          TERM=1
C          GO TO 100
C          K=K+1
C          T1=N+K
C          T2=N-K+1
C          T3=2*K
C          TERM=TERM*T1*T2/T3
C
C          IF(DABS(TERM) .LT. .000001)GO TO 1000
C          GO TO(100,1000),IGO
C          AR=AR+TERM
C          IGO=1
C          IF((K-N)20.1000,1000
C          AI=AI-TERM
C          IGG=2
C          TERM=-TERM
C          IF((K-N)20.1000,1000
C          X=AR
C          Y=AI
C
C          RETURN
C
C          THIS SUB. SHIFTS PHI UNTIL IT LIES IN THE RANGE -180+180
C          IF(PHI-180.)20.,20.+10.
C          PHI=PHI-360
C          GO TO 1
C          IF(PHI+180.)130.,40.+40.
C          PHI=PHI+360
C          GO TO 20
C
C          RETURN
C
C          END
C
C          SUBROUTINE ADJUSTPHI
C
C          THIS SUB. SHIFTS PHI UNTIL IT LIES IN THE RANGE -180+180
C          IF(PHI-180.)20.,20.+10.
C          PHI=PHI-360
C          GO TO 1
C          IF(PHI+180.)130.,40.+40.
C          PHI=PHI+360
C          GO TO 20
C
C          RETURN
C
C          END

```

```

SUBROUTINE LEGEND(NMAX,M,Z,VAL)
C THIS SUB. CALCULATES VALUES OF THE ASSOCIATED LEGENDRE
C FUNCTION WITH INDICES N=M TO N=NMAX. DO NOT USE NMAX LESS
C THAN M+1. VALUES CHECKED WITH TABULATED VALUES TO 5 PLACES
C THRU N=56 FOR M=1 AND N=10 FOR M=5.
C
C DIMENSION VAL( 81)
C DOUBLE PRECISION TERM1,TERM2,TERM3,T0
C
C VARIABLES:
C
C NMAX = MODE ORDER + 1
C MCOMP = ORDER OF AZIMUTHAL VARIATION
C Z = ARGUMENT OF FUNCTION = COS(THETA) OR COS(BIGTHETA)
C VAL = VALUES OF LEGENDRE POLY FOR EACH MODE (RETURNED TO
C FIELD1,FIELD2)
C
C ZD=2
C FM=M
C TERM1=0.0
C IFUM) 8.8*9
8 TERM2=1.0
GO TO 11
9 KMAX=2*M-1
TERM2=(1.0-ZD*ZD)**(FM/2.0)
DO 10 K=1,KMAX,2
FK=K
TERM2=TERM2*(2.0*FM-FK)
VAL(M)=TERM3
TERM1=TERM2
TERM2=TERM3
CONTINUE
RETURN
END
10 TERM3=((2.0*FN+1.0)*ZD*TERM2-(FN+FM)*(TERM1/(FN-FM+1.0))
VAL(N+1)=TERM3
TERM1=TERM2
TERM2=TERM3
CONTINUE
20 RETURN
END
C
SUBROUTINE DIFF(I,J, NMAX,NMAX,F, Fx,T0,DP)
C THIS SUB. RETURNS TO SUPER THE VALUE OF FT OR FP. THESE CERVIS
C ARE COMPUTED USING FORWARD AND BACKWARD DIFFERENCES.
C
C DIMENSION F(36,91),Fx(36,91),
C ITERM(21,20)
C
C SOME VARIABLES USED:
C
C ITERM = COEFFICIENTS WITHIN A TERM OF THE DIFFERENCE EQUATION
C NT = # TERMS IN EQUATION
C TI = INTEG GRID #
C J = INDICATOR OF WHETHER FT OR FP IS BEING SMOOTHED
C FX = EITHER FT OR FP
C DTR=0.01745329
C
C FOR II=1 CALCULATE TERM COEFFICIENT FOR NUMERICAL DERIVATIVE
C FORMULA. THESE COEFFICIENTS ARE IDENTICAL TO THOSE OF A
C BINOMIAL EXPANSION.
C
C IF(IJ .NE. 1)GO TO 10
C
C TERM(1,1)=1
C TERM(2,1)=-1
C DO 11 J=2,20
C JJP1=JJ+1
C DO 12 I=1,JJP1
C IF(IJ .NE. 1)GO TO 13
C TERM(I,JJ)=1
C GO TO 12
C IF(IJ .NE. JJP1)GO TO 14
C TERM(I,JJ)=ITERM(I-1,JJ-1)-ITERM(I-1,JJ-1)
C GO TO 12
C
C TERM(I,JJ)=ITERM(I,JJ-1)-ITERM(I-1,JJ-1)
C CONTINUE
C DO 71 JJ=1,1C
C JJP1=JJ+1
C WRITE(6,*)(0.001*(TERM(I,JJ)-ITERM(I,JJ)),I=1,JJP1)
C FORMAT(1H ,21f5)
C CONTINUE
C
C IF J=-1, COMPUTE FP; OTHERWISE COMPUTE FT.
C
C IF(IJ .EQ. -1)GO TO 50
C CONTINUE
C
C CALCULATE FP
C
C DEN=DP*DTR*0.0254
C
C ENTER THETA & PHI LONPS (PHI WITHIN THETA). PERFORM CALCULATION
C FOR EACH POINT ON INTEGR. GRID.
C
C DO 54 M=1,MMAX
C DO 54 N=1,NMAX
C Fx(M,N)=0.
C
C
04490000
04500000
04510000
04520000
04530000
04540000
04550000
04560000
04570000
04580000
04590000
04600000
04610000
04620000
04630000
04640000
04650000
04660000
04670000
04680000
04690000
04700000
04710000
04720000
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05460000

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NATIONAL RADIO ASTRONOMY OBSERVATORY
GREEN BANK, WEST VIRGINIA

ADDENDUM
TO
ELECTRONICS DIVISION INTERNAL REPORT No. 221

UPDATED DESCRIPTION OF USING THE JPL
PHYSICAL-OPTICS SCATTERING PROGRAM

JAMES R. LYONS

AUGUST 1982

NUMBER OF COPIES: 150

UPDATED DESCRIPTION OF USING THE
JPL PHYSICAL-OPTICS SCATTERING PROGRAM

James R. Lyons

I. Introduction

A field plotting routine, multiple-scattering capability, and several other features, have been added to the JPL program as described in EDIR No. 221 (September 1981). This memo is designed to briefly describe these additions and to provide the user with a convenient list of program inputs. Specifically, it is a revision of section III, Details of Using the Program, of report 221. It is suggested that the previous report be read before attempting to read and apply this memo to the program. Neither the basic structure nor the underlying theory of the program have been altered, but there are additional input parameters.

II. Added Program Features

A. Field plotting routine.

A subroutine, EPLOT, has been written which will plot the amplitude (dB) and phase of a particular component of an electric or magnetic field. The spherical wave expansion (SWE) program employs this subroutine to plot the E and H plane of the input field and of the far-field form of the spherical wave representation. The scattering (SCAT) program uses EPLOT to plot the polar, azimuthal, or radial component of the near-field of the incident pattern, and to plot the polar, azimuthal or Cartesian-Y component of the far-field form of the incident and scattered patterns. Only the first 30° of the far-field incident pattern can be plotted; for $\theta(1) > 30^\circ$, the plot and power of this pattern are skipped. Input parameters are used to choose the field component to be plotted.

In the output grid, the Y component of the electric field has the form:

$$E_Y(\theta, \phi) = E_\theta(\theta, \phi) \cos \theta \sin \phi + E_\phi(\theta, \phi) \cos \phi.$$

The $\cos \theta$ term induces an azimuthal variation on E_Y which would destroy any beam circularity. Since our currently relevant problems involve circularly symmetric beams, the amplitude of this term has been neglected. E_Y then has the form

$$E_Y(\theta, \phi) = E_\theta(\theta, \phi) \operatorname{sign}(\cos \theta) \sin \phi + E_\phi(\theta, \phi) \cos \phi.$$

Note also that $|\cos \theta| \approx 1$ for most scattering situations.

B. Multiple scattering.

In some scattering problems, there exists several reflectors, so it becomes necessary to perform a spherical wave expansion on a scattered field and reflect this off yet another surface. SCAT can do this by writing the Y component of the far-field scattered field values it computes into a disk data set. The SWE program then reads in the values and determines the expansion. SCAT stores 31 values (for $\Delta\theta = 1.0^\circ$, this corresponds to 30°), either those on the "right side" of the beam or the mirror-image of those on the "left side" of the beam. In either case, the polar angle of the field values is translated such that the beam peak occurs at $\theta = 0^\circ$. Since only half the beam is stored, the entire beam must be assumed circular. If the scattered field is asymmetric with respect to its beam peak, individually expanding both beam halves will define the extremes within which the true beam must lie.

The following table shows the data flow for the two programs when scattering from multiple reflectors. The numbers are the data set reference numbers used in the READ and WRITE statements.

		<u>READ</u>	<u>WRITE</u>
First reflection	{SWE SCAT	Cards 9 9	9 12
Subsequent reflections	{SWE SCAT	12 9	9 12

Note: SCAT only stores the pattern values corresponding to the final Φ of the output grid.

C. Power calculations.

Relative power calculations are made for the same patterns that can be plotted by EPLOT. For the near-field of the incident pattern in SCAT, the power integral has the form

$$P = \frac{1}{2\pi} \int_0^{2\pi} \int_0^\theta |\mathbf{H}_\theta(\theta, \phi)|^2 + |\mathbf{H}_\phi(\theta, \phi)|^2 \sin \theta d\theta d\phi.$$

For the far-field patterns in SCAT and the patterns in the SWE program, the integral is:

$$P = \int_0^{\theta_0} [|\mathbf{E}_\theta(\theta, \phi_0)|^2 + |\mathbf{E}_\phi(\theta, \phi_0)|^2] \sin \theta d\theta.$$

In the near-field integral, the power is calculated out to the reflector edge associated with each azimuthal angle. The cumulative power thru each integration grid is output along with the average value of the edge-angle of the integration grid. Power is also calculated for any section of the integration grid which lies beyond the reflector edge, thus allowing the spillover efficiency to be determined, if desired.

The far-field power integral can only be evaluated out to a constant polar angle. For the far-field incident pattern of SCAT, this angle is the average of the average edge-angles computed in the near-field case. For the scattered

pattern in SCAT and the patterns in the SWE program, the angle thru which the power is calculated is READ in by the programs.

The power contribution between two polar angles on the grid is evaluated at the midpoint of the angles:

$$\Delta P = \frac{1}{2} [|E(\theta_1)|^2 + |E(\theta_2)|^2] \cdot \sin\left(\frac{\theta_1 + \theta_2}{2}\right) \cdot \Delta\theta.$$

When the integration limit lies between grid values, ΔP is linearly scaled to the limit value. For the incident and SWE patterns, ΔP is doubled for each annulus because the pattern is actually that of a half-beam.

D. Field interpolation.

To accurately represent an input pattern, the SWE program requires values to be input at least every 0.25° or 0.5° in polar angle. Since most patterns are reasonably "smooth" thru the first 20 dB or so, points are input to the program every 1.0° or 2.0° and a linear interpolation is used to include points every 0.25° or 0.5° , respectively. To better represent the rounded beam peak, the first three interpolated values (i.e., .25, .5, .75, or .5, 1.0, 1.5) have a perturbation added to the interpolation.

The maximum number of points MAINOR can handle in computations is 121. Therefore, the maximum number of input points is 31 (e.g., 0° thru 30° , every 1.0°).

III. Input Variables and Their Formats

For both programs this section lists the input variables, gives a brief description of each variable, provides a concise table of inputs and their formats, and lists the output produced by the program.

A. SWE program.

1. Input variables descriptions:

Angles and phase are in degrees. All inputs read-in by MAIN.

NAMEJ	- Alphanumeric: Use to identify program; ≤ 72 characters.
MCOMP	- Order of azimuthal variation; usually = 1.
LMAX	- Maximum mode order, ≤ 80 ; usually ≈ 70 .
NAME	- Alphanumeric: Use to identify reflector, pattern, etc.
JIN	- Number of pattern points used in program, ≤ 121 ; $= 4^* (\text{actual number of card inputs} - 1) + 1$.
IC1	- If ≤ 0 , convert input field from dB to volts.
IC2	- If > 0 , set input field phase to zero.
IC3	- If > 0 , compute pattern from equations inserted
IC4	- If > 0 , E and H plane patterns are identical; input only E plane.
NDISK	- Pattern parameter: = 1 if values are to be input from cards; = 2 if values are read from disk data set (#12).
IPLANE	- Specifies plane: = 1 if E plane power is calculated and E plane pattern is plotted; = 2 for H-plane.
IPLOT1	- Input pattern print-out parameter: = 1 for a plot; = -1 for numerical values; = -11 for both; = 0 for neither.
IPLOT2	- Output (SWE) pattern print-out parameter: Same as IPLOT1.
TEDGE	- Polar angle thru which power flow is calculated.
T(J)	- Polar angle of input pattern values.
E(J), H(J)	- E, H-plane input pattern amplitude.
EP(J), HP(J)	- E, H-plane input pattern phase.
JMAX0	- $= 180/\Delta\theta + 1$, where $\Delta\theta$ is desired output increment of the output (SWE) pattern; typically = 181.
JOUT	- Number of output points starting with $\theta = 0^\circ$; typically = 91.

2. Table of inputs and their formats.

Card No.	Variables	Format
1	NAMEJ	18A4
2	MCOMP, LMAX	2I5
3	NAME	18A4
4	JIN, IC1, IC2, IC3, IC4, NDISK	6I5
5	IPLANE, IPLOT1, IPLOT2	3I5
6	TEDGE	F10.5
7	IC4≤0: T(1), E(1), EP(1), H(1), HP(1)	5F10.5
⋮	⋮	⋮
If NDISK = 1 JW + 6	IC4>0: T(1), E(1), EP(1) ⋮ IC4≤0: T(JIN), etc. IC4>0: T(JIN), etc.	3F10.5 5F10.5 3F10.5
JIN + 7	JMAXO, JOUT	2I5
No cards: Data read from computer disk (#12)	NAMEJ T(1), E(1), EP(1) ⋮ T(JIN), etc.	(18A4) (3F10.5) (3F10.5)
7	JMAXO, JOUT	2I5

3. SWE program print-out.

- Two alphanumeric statements.
- Input pattern: plot or numerical values.
- Power in input pattern.
- Real and imaginary values of SWE coefficients.
- Fraction of total mode power in the coefficients for each mode order.
- Total coefficient mode power (not related to computed power in pattern).
- Output pattern (far-field of SWE): plot or numerical values.
- Power in output pattern.

B. SCAT Program.1. Input variable descriptions.

All linear measures are in meters; all angles are in degrees.

Parentheses indicate which subroutine is reading in data.

(MAIN)

- TITLE - Alphanumeric: describe reflector, wavelength, etc.,
≤ 72 characters.
 - NDISK - disk writing parameter: = 1 if 1st half of scattered
values are to be stored in a disk data set; = 2 for
2nd half of pattern; = 0 if no values are to be stored.
 - PC - phase constant = $2\pi/\lambda$
 - XT, YT, ZT - translation to expected phase center of scattered
pattern.
 - ALPHA - used as a reflector rotation parameter when surface
is specified analytically; set to 0.0 if not needed.
 - RHOØ - distance to the reflector at $\theta = 0^\circ$.
 - THETAØ - used as an edge parameter when theta-edge is specified
analytically; set to 0.0 if not needed.
- - - - -

(SETUP)

- JMAX - number of θ values in output grid; ≤ 361.
- TT1 - initial θ value.
- DTT - θ increment, usually = 1.0° .
- KMAX - number of ϕ values in output grid; ≤ 46.
- PP1 - initial ϕ value.
- DPP - ϕ increment.
- NG1 - number of integration grids, ≤ 21.
- MM(I) - number of θ values on Ith integration grid; ≤ 36.
- T1 - initial θ value.
- DT - θ increment, typically between 0.2° and 1.0° .

NN(I) - number of ϕ values on Ith integration grid; ≤ 91 .

P1 - initial ϕ value.

DP - ϕ increment, typically between 2° and 10° .

(Note: Subroutine EPLOT and the power calculations assume that DT and DP are constant within an integration grid, and that DTT is constant for the output grid.)

- - - - -

(FIELDS)

NAME - alphanumeric: identify SWE program.

LMAX - maximum mode order.

MCOMP - order of azimuthal variation.

NAME - alphanumeric: identify reflector, pattern, etc.

J - 1 to LMAX.

A(J,1), A(J,2) - real and imaginary components of SWE "A" coefficients.

B(J,1), B(J,2) - real and imaginary components of SWE "B" coefficients.

- - - - -

(MAIN)

IEDGE - reflector parameter used to determine if and how program obtains the reflector edge-values of θ . Cases:

< 0 - calculate values using equations inserted in EDGEEQ . READ in 1 value, TEDGE1, to be used in equations.

= 0 - θ at edge is a constant. READ 1 value, TEDGE(1).

= 11 - edge will not be encountered on present integration grid. READ no values.

> 0 - θ at edge is in tabular form and is read in for each ϕ on the integration grid. READ in NMAX values.

TEDGE(N) - edge value of θ . Read in either NMAX, 1, or no values.

- - - - -

(SURF)

- ISURF**
- reflector parameter used to determine how $\rho(\theta, \phi)$, $\partial\rho/\partial\theta(\theta, \phi)$, and $\partial\rho/\partial\phi(\theta, \phi)$ are obtained. Cases:
 - < 0 - ρ , $\partial\rho/\partial\theta$, $\partial\rho/\partial\phi$ all determined analytically ($\neq -2$)
from equations in SURF. READ no values.
 - = -2 - ρ determined from equations; derivatives numerically computed. READ no values.
 - = 0 - ρ is in tabular form and is a function of θ ; derivatives numerically computed. READ MMAX values.
 - 0 - ρ is in tabular form and is a function of θ and ϕ ; derivatives numerically computed.
READ MMAX x NMAX values.
- F(M, N)**
- $\rho(\theta, \phi)$. READ either MMAX x NMAX, MMAX, or no values.

(MAIN)

- NIPLOT**
- near-field incident pattern parameter:
 - = 1, 2, 3 - plot $H_r(\theta, \phi_0)$, $H_\theta(\theta, \phi_0)$, $H_\phi(\theta, \phi_0)$, respectively.
 - = 0 - no plot.
- PHIPLT**
- ϕ_0 used with NIPLOT: ϕ_0 must be on present integration grid segment.
- FIPILOT**
- far-field incident pattern parameter:
 - = 1 - plot $E_Y(\theta, \phi_0)$ = Y component.
 - = 2, 3 - plot $E_\theta(\theta, \phi_0)$, $E_\phi(\theta, \phi_0)$, respectively.
 - = -1 - numerical values listed.
 - = 0 - neither plot nor numerical values.
- FSPILOT**
- far-field scattered pattern parameter.
(Same as FIPILOT.)
- TEDGEA**
- half-angle from beam center, thru which power in scattered pattern is computed.

2. Table of inputs and their formats.

A dash (-) in the Card No. column indicates that the card number is not known or is difficult to determine.

Card No.	Variables	Format
1	TITLE	18A4
2	NDISK	I5
3	PC, XT, YT, ZT, ALPHA, RHO \emptyset , THETA \emptyset	7F10.4
4	JMAX, TT1, DTT	I5, 2F10.2
5	KMAX, PP1, DPP	I5, 2F10.2
6	NG1	I5
7	MM(1), T1, DT	I5, 2F10.2
8	NN(1), P1, DP	I5, 2F10.2
:	:	:
2NG1 + 5	MM(NG1), T1, DT	I5, 2F10.2
2NG1 + 6	NN(NG1), P1, DP	I5, 2F10.2
NO CARDS:	NAME	(18A4)
Data	LMAX, MCOMP	(I5)
read	NAME	(18A4)
from	1, A(1,1), A(1,2), B(1,1), B(1,2)	(I5, 2E17.8, 2X, 2E17.8)
computer	:	:
disk (#9).	LMAX, A(LMAX,1), A(LMAX,2)	(I5, 2E17.8, 2X, 2E17.8)
	B(LMAX,1), B(LMAX,2)	
Read for each integration grid. Repeat NG1 times.	IEDGE	I5
	TEDGE(N) (NMAX, 1, or no cards)	7(1X,F9.5)
	ISURF	I5
	F(M,N) (MMAX x NMAX, MMAX or no cards)	7(1X,F9.5)
	NIPILOT, PHIPLT, FIPILOT, FSPILOT	I5,F7.2,2I5
	TEDGEAE	F10.4

3. SCAT program print-out.

```

-- 2 alphanumeric statements followed by propa-
gation constant.

-- θ and ϕ for each integration grid segment.

For each integration grid. { -- integration grid number and value of IEDGE.
                           -- if edge is to be encountered, list of θ-edge
                           and ϕ-edge as functions of ϕ.
                           -- integration grid number and value of ISURF.
                           -- plot of near-field of incident pattern.
                           -- near-field incident pattern power.

For each ϕ. { -- far-field of incident pattern: plot or numbers.
              -- far-field incident pattern power.

              -- phase center translations, scale factor (set
                 to 1.0).

For each ϕ. { -- far-field of scattered pattern: plot or numbers.
              -- far-field scattered pattern power.

```

C. Submitting a job and organization of JCL.

The SCAT and SWE programs have been run on the IBM 4341 computer located in Charlottesville. The operating system used is called Pandora; blocks of program code (e.g., subroutines) are stored under various Pandora member names (PMN^S).

To submit a job with the Pandora system, the SUBMIT command is used followed by the Pandora members making up the program and data. All the Pandora members (except EPLOT, which is just the plotting routine) have already been described in EDIR #221.

The submit command for the SWE program is SUBMIT_JCLSWE_MAINOR_EPLOT_LVM_DATASWE.

The JCL and the actual subroutines submitted are shown below. The PMN for a subroutine or block of subroutines is shown in the left-hand column.

PMN

JCL and Subroutines

MAINOR

```
//SWAVE_ _JOB_(userI.D.),user name,MSGLEVEL=(2,0),
  CLASS=L,TIME=1
  /*ROUTE_ _PRINT_REMOTE1
  // _EXEC _FORTGCLG,ERROR=E,PARM.FORT=ID
  //FORT.SYSPRINT_DD_DUMMY
  //FORT.SYSIN_DD_*
```

[MAIN

E PLOT

```
[ Subroutines LEGEND, VECTOR, MULT
  /*
  //LKED.SYSPRINT_DD_DUMMY
  //GO.FT09F001_DD_DSN=user name.DATA,DISP=SHR
  //GO.FT12F001_DD_DSN=user name.DSCAT,DISP=SHR
  //GO.SYSIN_DD_*
```

LVM

DATA SWE

```
[ DATASWE
  /*
```

A CLASS = L job uses a 216 K byte memory partition, double the size of a standard partition. The SWE program could be reduced to a standard partition by (carefully) reducing array sizes. If this is done, the corresponding dimensions in the CALL MULT statements in MAINOR must be changed along with the array dimensions in several subroutines.

For both SWE and SCAT jobs the TIME parameter is in (CPU) minutes.

The submit command for the SCAT program is

```
SUBMIT_JCLSCAT_MAIN1SC_EPLOT_SSURF_FINT_SFPVSAJD_DATASCAT
```

The JCL and subroutine structure is as follows:

PMN	<u>JCL and Subroutines</u>	
JCLSCAT	//SCAT_ _ _ _ _ JOB_(userI.D.), user name, MSGLEVEL=(2,0), CLASS=0, TIME=user set ----- ----- Same as last 4 lines of JCLSWE -----	
	MAIN1SC [MAIN	
	EPLOT [Subroutine EPLOT	
	SSURF [Subroutine SURF	
	FINT [Subroutine FINT	
SFPVSAJD	Subroutines SETUP, FIELDS, PATHL, VECTOR, SPHANK, ADJUST, LEGEND, DIFF ----- ----- Same 5 lines of JCL as at end of member LVM -----	
	DATASCAT	DATASCAT
		/*

A CLASS = 0 job is an extra large partition of 880 K bytes. SCAT (between 300 and 400 K bytes) could be reduced to an L job by reducing the array dimensions of the output grid and the integration grid, if necessary, and by decreasing the number of allowable integration grids. Note that the array dimensions in several subroutines must be changed.

D. Creating and listing disk data sets.

The following JCL is used to create a sequential data set on the computer disk.

```
//CREATEDS__JOB_(userI.D.),user name,CLASS=Q,  
MSGLEVEL=1  
  
/*ROUTE__PRINT_REMOTE1  
  
//NEWFILE__EXEC_PGM=(,CATLG),DSN=user name.DATA  
or DSCAT,UNIT=3300,SPACE=(CYL,(1,1)),  
DCB=(RECFM=FB,LRECL=80,BLKSIZE=1600).
```

To list the contents of a data set:

```
//LIST_JOB_(userI.D.),user name,CLASS=Q,MSGLEVEL=1  
/*ROUTE__PRINT_REMOTE1  
  
//_EXEC_LIST  
  
//SYSIN_DD_DSN=user name.DATA or DSCAT,DISP=SHR
```

Only the above JCL need be submitted to perform the desired task.

Acknowledgements

I thank Rick Fisher for his many helpful suggestions and constructive criticisms and Carolyn Dunkle for typing this report.

IV. Changed and Additional Program Code.

```

G LEVEL 21          MAIN          DATE = 82211      11/15/21      G LEVEL 21          MAIN          DATE = 82211      11/15/21
C          SWE PROGRAM
C          THIS IS A SPHERICAL WAVE EXPANSION PROGRAM-ORTHOGONALITY      00010000 110    CONTINUE
C          TECHNIQUE. THE SPHERICAL WAVE COEFFICIENTS FOR THE M-TH      00020000 C      READ-IN PATTERN FROM CARDS
C          FOURIER COMPONENT OF A RADIATION PATTERN ARE COMPUTED.      00030000
C          THE EXPANSION MATCHES THE FAR-FIELD INPUT PATTERN WITH      00040000
C          THE FAR-FIELD FORM OF SPHERICAL WAVES.      00050000
C          DIMENSION NAME(16),NAMEJ(18)
C          DIMENSION PA(80),PB(80)
C          DIMENSION T(121),E(121),H(121),HP(121)
C          DIMENSION F(80,121)*G(80,121),M(81)
C          DIMENSION A(121,2),B(121,2)
C          DIMENSION ACCE(80,2),BCDE(80,2)
C          DIMENSION ACUT(1,21)*BCUT(1,21)
C          CTR=C*017453293
C          DT2=DR/2.0
C          SCHE VARIABLE DEFINITIONS:
C          T,E,EP,R,HP=INPUT PATTERN PARAMETERS
C          JIN(<121)=# THETA INPUT VALUES
C          NMAX(<=80)= MAXIMUM MODE ORDER
C          ACOE,BCOE = REAL & IMAG COMPONENTS OF TE AND TM WAVE COEFFICIENTS
C          WRITTEN ON DISK
C          A*B = AMP & PHASE OF INPUT FIELDS VALUES :: DELTA(THETA)
C          F*G = MATRICES INMAX BY JIN OR NMAX BY JOUT OF VALUES RELATED TO
C          ASSOCIATE LEGENDRE POLYNOMIALS
C          G PROP. TO CERTRATIVE WRT THETA OF N-TH POLY.
C          F PRCP. TO N-TH POLY.
C          INPUT ALPHANUMERIC STATEMENTS,PATTERN & PLOTTING PARAMETERS.
C          READ(5,2223)NAMEJ
C          WRITE(6,2001)NAMEJ
C          READ(5,1002)IMCCMP,NMAX
C          READ(5,2223)NAME
C          WRITE(6,2001)NAME
C          READ(5,1002)JIN,IC1,IC2,IC3,IC4,NDISK
C          READ(5,1002)IPLANE,IPLC1,IPLC2
C          READ(5,2013)TEGE
C          IF IC3>0,CALCULATE INPUT PATTERN PARAMETERS FROM EQUATION
C          IF IC3<0 AND IC4>0,READ IN ONLY T,E,EP. I.E.,PATTERN IS CIRC.
C          IF IC4=0,CONTINUE
C          IF IC3=113,113,133
C          IF IC3=153,153,163
C          CONTINUE
C          READ-IN PATTERN FROM DISK DATA SET (#12)
C          IF INCISK .EQ. 1,CC TC 110
C          READ(12,223)NAMEJ
C          WRITE(6,2014)NAMEJ
C          READ(12,2013)(T(J),E(J),EP(J),J=1,JIN,4)
C          GO TC 143
C          CONTINUE
C          DELT=0.25
C          DC 123 N=1,JIN
C          READ(5,2013)(T(M),E(M),EP(M),M=1,JIN,4)
C          CONTINUE
C          READ(5,2015)(T(M),E(M),EP(M),H(M),HP(M),M=1,JIN,4)
C          CONTINUE
C          IF IC1 < CR = 0 CONVERT FROM DB TO VOLTS
C          IF IC1 10,10,20
C          DO 15 J=1,JIN,4
C          E(J)=10.0*(E(J)/20.0)
C          H(J)=10.0*(H(J)/20.0)
C          CONTINUE
C          INTERPOLATE PATTERN VALUES
C          JINM=JIN-4
C          DO 167 J=1,JIN,4
C          JP2=J+2
C          DO 166 K=J,JP2
C          FD4=FLCAT(K+1-J)/4.
C          T(K+1)=(T(J)+FD4*(T(J+4)-T(J))
C          E(K+1)=E(J)+FD4*(E(J+4)-E(J))
C          EP(K+1)=EP(J)+FD4*(EP(J+4)-EP(J))
C          H(K+1)=H(J)+FD4*(H(J+4)-H(J))
C          HP(K+1)=HP(J)+FD4*(HP(J+4)-HP(J))
C          ADD PERTURBATION
C          IF(J .GT. 1)GO TO 166
C          PERT=25/FLOAT(K+J)
C          E(K+1)=E(K+1)+PERT*(E(J)-E(J+4))
C          H(K+1)=H(K+1)+PERT*(H(J)-H(J+4))
C          CONTINUE
C          00873000
C          00874000
C          00875000
C          00876000
C          00877000
C          00878000
C          00880000
C          00931000
C          00940000
C          00950000
C          00970000
C          00980000
C          00990000
C          01000000
C          01010000
C          01020000
C          01025000
C          01030000
C          01040000

```

```

G LEVEL 21          MAIN          DATE = 82211    11/15/21    G LEVEL 21          MAIN          DATE = 82211    11/15/21
123  CONTINUE          01070000          C          01070000          C          WRITE(6,2011)EDGE*PINPS
143  CONTINUE          01080000          C          01080000          C          CCNVERT TC RADIAN AND REAL AND IMAG
C          FCR IC2 > 0 NEGLECT PHASE          01090000          C          01090000          C          01590000
C          IF(IC2)4 C=40*30          01100000          C          01100000          C          01560000
C          DC 35 J=1,JIN          01110000          C          01110000          C          01570000
C          EP(JJ)=0.          01120000          45          01120000          45          01600000
C          HP(JJ)=0.          01130000          50          01130000          50          01610000
C          H(JJ)=0.          01140000          50          01140000          50          01620000
C          CONTINUE          01150000          50          01150000          50          01630000
C          DC 45 J=1,JIN          01160000          C          01160000          C          01640000
C          T(JJ)=DR*T(JJ)          01170000          C          01170000          C          01650000
C          IF(IC2)50*60          01180000          C          01180000          C          01660000
C          DC 55 J=1,JIN          01190000          C          01190000          C          01670000
C          TH=CTR*EP(JJ)          01200000          C          01200000          C          01680000
C          EP(JJ)=E(JJ)*SINTH          01210000          C          01210000          C          01690000
C          E(JJ)=E(JJ)*COSTH          01220000          C          01220000          C          01700000
C          TH=CTR*T(JJ)          01230000          C          01230000          C          01710000
C          HP(JJ)=H(JJ)*SINTH          01240000          C          01240000          C          01720000
C          H(JJ)=H(JJ)*COSTH          01250000          C          01250000          C          01730000
C          CONTINUE          01260000          C          01260000          C          01740000
C          ESTABLISH DIFFERENTIAL PATTERN VALUES AT EACH INPUT POINT OF
C          PATTERN. VALUES ARE STORED IN VECTORS A & B AND ARE USEC IN
C          EVALUATING THE COEFFICIENT INTEGRALS.          01270000          C          01750000
C          01280000          C          01280000          C          01760000
C          01290000          C          01290000          C          01770000
C          01300000          C          01300000          C          01780000
C          01310000          C          01310000          C          01790000
C          01320000          C          01320000          C          01800000
C          01330000          65          01330000          65          01810000
C          PRINT CUT INPUT PATTERN          01340000          C          01340000          C          01820000
C          WRITE(6,2006)MCOMP          01350000          C          01350000          C          01830000
C          WRITE(6,2002)          01360000          C          01360000          C          01840000
C          WRITE(6,2003)(T(M)*EP(M)*H(M)*HP(M)*M=1,JIN)          01380000          C          01850000
C          CONTINUE          01390000          C          01390000          C          01860000
C          CCNTINUE          01400000          C          01400000          C          01870000
C          FMC=MCCMP          01410000          C          01410000          C          01880000
C          DC 80 J=1,JIN          01420000          C          01420000          C          01890000
C          Z=CCS(T(JJ))          01430000          C          01430000          C          01900000
C          00 85 N=1,MCCMP          01431000          85          01431000          85          01910000
C          F(N)=0.          01442000          C          01442000          C          01920000
C          NC=NMAX+1          01450000          C          01450000          C          01930000
C          CALL LEGENDINC,MCCMP,Z,PM)          01460000          C          01460000          C          01940000
C          DC 80 T=1,NMAX          01470000          C          01470000          C          01950000
C          F(I,J)=FMC*PM(I)          01480000          C          01480000          C          01960000
C          T1=1-PCCMP+1          01490000          C          01490000          C          01970000
C          T2=I+1          01482000          8C          01482000          8C          02020000
C          G(I,J)=T1*PM(I+1)-T2*Z*PM(I)          01483000          C          01483000          C          02030000
C          CCNTINUE          01490000          C          01490000          C          02040000
C          EVALUATE COEFFICIENT INTEGRALS BY MULTIPLICATION OF F & G
C          MATRICES AND A & B MATRICES. ACOE & BCCE RETURNED FROM MULT.
C          CALL MULT(JIN,NMAX,2,F,A,ACOE,0,80,121,121,2,80,2)
C          CALL MULT(JIN,NMAX,2,G,B,BCCE,1,80,121,121,2,80,2)
C          CALL MULT(JIN,NMAX,2,F,B,ACOE,0,80,121,121,2,80,2)
C          CALL MULT(JIN,NMAX,2,G,A,BCOE,1,80,121,121,2,80,2)
C          CALL MULT(JIN,NMAX,2,G,A,BCOE,1,80,121,121,2,80,2)
C          PRINT-CUT PCWER RESULTS
C
11
C

```

18

```

G LEVEL 21          MAIN          DATE = 82211    11/15/21   MAIN          DATE = 82211    11/15/21
270 CONTINUE
C   IF(J .EQ. 1) GO TO 274
C   IF(J .GT. 1) GO TO 276
C   CALCULATE POWER IN OUTPUT PATTERN
C
C   IF(IPLANE .EQ. 1) INPUT=(EAMP+EAMP+EAMP+HAMP+HAMP+HAMP)*(TH+TH1)+C1)
C   IF(IPLANE .EQ. -1) INPUT=(HAMP+HAMP+HAMP+HAMP+HAMP+HAMP)*(TH+TH1)+C2)
C   POUTS=POUTS+PCUT*DT
C   EAMP=HAMP
C   HAMP=HAMP
C   TH=TH
C   CONTINUE
C   E1(J)=EAMP
C   EP1(J)=EPH1
C   H(J)=HAMP
C   HP(J)=HPH1
C   T(J)=TH
C   J=J+1
C   IF(J-JIN)200,-280,300
C   IF(J-JMAX)200,+285,300
C   ISIGN=1
C   TH=180.
C   IF(MCMMP-1)370,380,370
C   DC 375 N=1,NMAX
C   F(1+N)=C
C   G(1+N)=0
C   GC TC 215
C   DC 296 N=1,NMAX
C   ISIGN=-ISIGN
C   FN=N*(N+1)*ISIGN
C   F1+N=-FN/2.C
C   G1+N=FN/2.0
C   GC TC 215
C   CCNTINUE
C   IF(IPLCT2 .EQ. C1 .OR. (IPLCT2 .EQ. 1))GC TO 310
C   WRITE(6,2002)
C   PRINT-CLT PATTERN VALUES
C
C   WRITE(6,2003)(T(J),E(J)*EP(J)*H(J)*HP(J),J=1,JPLT)
C   310 CONTINUE
C   WRITE(6,2012)EDGE,PCLT$ STOP
C
IC02 FORMAT(6I5)
1C11 FORMAT(1HO,1X,2I4)
2C01 FORMAT(1HO,1B4)
2C02 FORMAT(IHO,1X,POLAR,1X,E-PLANE,1X,H-PLANE,1X,DEG,1)
2C03 FORMAT(F1C,2,F12.6,F8.2,F13.6,F8.2)
2C04 FORMAT(IHC,SUPERFLC,WAVE COEFFICIENTS FOR AZIMUTHAL ORDER • IZ//,
          E20X,4H(N)•32X,4HB(N)•27X,4HIMAG,15X,4HREAL,13X,4HIMAG,
          E5H,4HREAL,12X,4HIMAG,15X,4HREAL,13X,4HIMAG,
          E13X,7H,MCODE,6X,7H,MCODE,5X,16HIMAG,15X,4HREAL,13X,4HIMAG,
          E20X,15•2E17.8,2X,2E17.8,2X,2E14.5,F14.8)
2C05 FORMAT(1HO,1X,INPUT PATTERN FOR AZIMUTHAL CUMPCNENT OF CRCR,• IZ)
2C06 FORMAT(IHO,1X,INPUT PATTERN FOR AZIMUTHAL CCMPCNENT OF CRCR,• IZ)

03240000 2007 FORMAT(IHO,1X,TOTAL MODE POWER • E15.8,• WATTS•)
03250000 2008 FORMAT(IHC,1X,FAR FIELD SUMMATION CF SPHERICAL MODES•)
03260000 2009 FORMAT(F1C,2,F10.6,F10.2,F10.6,F10.2)
03270000 2011 E• SQUARE VOLTS• THE INPUT POWER THRU • F6•2,• DEGREES IS • F6•4,
03271000 2012 FORMAT(IHO,1X,C• SQUARE VOLTS•) THE OUTPUT POWER THRU • F6•2,• DEGREES IS • F6•4,
03272000 2013 FORMAT(3F10.5)
03280000 2014 FORMAT(C• RADIATION SOURCE: • 1RA4)
03285000 2015 FORMAT(5F10.5)
03290000 2016 FORMAT('O',53X,*E-PLANE*) INPUT PATTERN• /14X,*E-PLANE*)
03300000 2017 FORMAT('C',53X,*H-PLANE*) INPUT PATTERN• /14X,*H-PLANE*)
03310000 2018 FORMAT('•',53X,*E-PLANE*) OUTPUT PATTERN• /14X,*E-PLANE*)
03315000 2019 FORMAT('1',53X,*H-PLANE*) OUTPUT PATTERN• /14X,*H-PLANE*)
03320000 2022 FORMAT(15.2E17.8*2X,*E17.8)
03330000 2223 FORMAT(18A4), END
03340000
03350000
03360000
03370000
03380000
03390000
03400000
03410000
03420000
03430000
03440000
03450000
03460000
03470000
03480000
03490000
03500000
03510000
03520000
03530000
03540000
03550000
03560000
03561000
03562000
03563000
03570000
03580000
03590000
03600000
03610000
03620000
03630000
03640000
03650000
03660000
03670000
03680000
03690000
03700000
03710000
03720000

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G LEVEL 21 E PLOT DATE = 82211 11/15/21 G LEVEL 21 E PLOT DATE = 82211 11/15/21
 00010000 C CO 861 M=1•121•10
 00020000 IARRAY(1)=EYE
 00030000 CONTINUE
 00040000 DB INCREMENT; PHASE RANGE IS
 00050000 00630000 CONVERT AMP FROM VOLTS TO DB
 00060000 00640000
 00070000 00650000 IF(EAMP •GE• 0.001)AMPDB=20•ALOGIO(EAMP)
 00080000 IF(EAMP •LT• 0.001)AMPDB=-50.0
 00090000 00670000
 00100000 00680000
 00110000 00590000
 00120000 00700000
 00130000 00710000
 00140000 00720000
 00150000 00730000 DETERMINE ARRAY LOCATIONS
 00160000 00740000
 00170000 00750000 SAMP=2.5*(EAMPDB + 46.0)
 00180000 SPHASE=.3333333*(EPHI+180.)
 00190000 INT(SAMP+.5)
 00200000 NAP=
 00210000 NPP= INT(SPHASE+.5)
 00220000 IF(NPP •NE• NAPIGO TC 830
 00230000 IF AMP & PHASE LOCATIONS COINCIDE,PRINT A •••
 00240000 C 00800000
 00250000 C 00820000
 00260000 C 00830000
 00270000 C 00840000
 00280000 C 00850000
 00290000 C ASSIGN AMP & PHASE SYMBOLS TO IARRAY
 00300000 C 00860000
 00310000 C 00870000
 00320000 C 00880000
 00330000 C 00890000
 00340000 C 00900000
 00350000 C 00910000
 00360000 C SUPPRESS PHASE SYMBOLS FOR AMP < -46 DB
 00370000 C IF((NAP •LT• 0) IARRAY(NPP-1)=BLANK
 00380000 DO 850 M=1•IDIT
 00390000 IF(INPRINT •NE• 1)GC TO 840
 00400000 C 00920000
 00410000 C PRINT ANGLE•DASH•REFERENCE LINE WITH AMP & PHASE,AND DASH
 00420000 C 00930000
 00430000 C 00940000
 00440000 C 00950000
 00450000 C 00960000
 00460000 C PRINT DASH•AMP & PHASE•AND DASH
 00470000 C 01040000
 00480000 C 01050000
 00490000 C 01060000
 00500000 C 01070000
 00510000 C 01080000
 00520000 C 01090000
 00530000 C 01100000
 00540000 C 01110000
 00550000 C PRINT AMP & PHASE SCALE VALUES IN CENTER CF PLOT
 00560000 C 01120000
 00570000 C 01130000
 00580000 C 01140000
 00590000 C 01150000
 00600000 C 01160000

815 C WHEN ABOVE CONDITION IS MET,PRINT A REFERENCE LINE OF DASHES
 816 C NCDAASH=NCDAASH+1
 817 C PRINT PRELIMINARY PLOT LABELS
 818 C WRITE(6•21071)
 819 C WRITE(6•21021)
 820 C WHEN ABOVE CONDITION IS MET,PRINT A REFERENCE LINE OF DASHES
 821 C NCDAASH=NCDAASH+1
 822 C PRINT DASH•AMP & PHASE•AND DASH
 823 C 01030000
 824 C 01040000
 825 C 01050000
 826 C 01060000
 827 C 01070000
 828 C 01080000
 829 C 01090000
 830 C 01100000
 831 C 01110000
 832 C 01120000
 833 C 01130000
 834 C 01140000
 835 C PRINT AMP & PHASE GRID MARKS IN CENTER CF PLT. SKIP FOR
 836 C J < 61
 837 C PRINT AMP & PHASE SCALE VALUES IN CENTER CF PLOT
 838 C IF(J •EQ• (JMAX-1)/2)WRITE(6•2113)
 839 C IF(J •EQ• (JMAX+3)/2)WRITE(6•2114)

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G LEVEL 21          EPLOT          DATE = 82211      11/15/21      G LEVEL 21          MULT          DATE = 82211      11/15/21
870  CONTINUE
C   FINISH PLCT
C   IF(IJ .NE. JMAX) RETURN
C   PRINT COUTPLT PLOT LABELS AND PEAK AMP
C   WRITE(6,2104)
C   WRITE(6,2110)
C   WRITE(6,1091)
C   WRITE(6,2112) PKAMP
C   RETURN
2101  FORMAT(8X,1A1,121A1,1A1)
2102  FORMAT(8X,1A1,121A1,1A1)
E,,   26.0   22.0   18.0   42.0   38.0   34.0   30.0*
E,,   2.0    0     +2.0    0     10.0   6.0*
2103  FORMAT(4X,F4.0,122A1)
2104  FORMAT(7X,1A1,121A1,1A1)
E,,   12.0   10.0   8.0    6.0    4.0    2.0    1.0*
E,,   1.0    0.8   0.6    0.4   0.2    0.1    0.0*
2107  FORMAT(10X,1A1,121A1,1A1)
E,,   X: AMPLITUDE*(1-DB*)
2109  FORMAT(8X,1A1,121A1,1A1)
E,,   0: PHASE*DEGREES*
2110  FORMAT(7X,-180     -150    -120    -90    -60    -30    0
E,,   -30     0      30     60     90     120    180
E,,   180     180     180     180     180     180     180
2112  FORMAT(10X,1A1,121A1,1A1)
2113  FORMAT(10X,1A1,121A1,1A1)
E,,   26     18     14     10     6     34     38
E,,   4.2    1.8    1.4    1.0    0.6    0.34    0.38
2114  FORMAT(10X,1A1,121A1,1A1)
E,,   -30    -150   -120   -90   -60   -30    0
E,,   150    180    180    180    180    180    180
END
01170000  SUBROUTINE MULT(M,N,K,A*B,C,IC,NA,MA,NB,MB,NC,MC)
01180000
01190000
01200000  C
01210000
01220000  C
01230000  10
01240000
01250000
01260000  15
01270000
01280000  20
01290000
01300000
01310000
01320000  25
01330000
01340000  30
01350000
01360000
01370000  35
01380000
01390000
01400000
01410000
01420000
01430000
01440000
01450000
01460000
01470000
01480000
01490000
01500000
01510000
01520000
00540000
00550000
00560000
00570000
00580000
00590000
00600000
00610000
00620000
00630000
00640000
00650000
00660000
00670000
00680000
00690000
00700000
00710000
00720000
00730000
00740000
00750000
00530000
THIS SUB. MULTIPLIES 2 FULL MATRICES
DIMENSION A(NA,MA),B(NB,MB),C(NC,MC)
IF(IC).LT.20,30
CC 15 I=1,N
DO 15 J=1,K
DO 15 L=1,M
C(I,J)=C(I,J)-A(I,L)*B(L,J)
RETURN
DC 25 I=1,N
DO 25 J=1,K
C(I,J)=0.0
DO 25 L=1,M
C(I,J)=C(I,J)+A(I,L)*B(L,J)
RETURN
DC 35 I=1,N
DO 35 J=1,K
DO 35 L=1,M
C(I,J)=C(I,J)-A(I,L)*B(L,J)
RETURN
END
THIS SUBROUTINE WAS INADVERTENTLY LEFT OUT OF THE FIRST REPORT.

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G LEVEL 21      MAIN          DATE = 82211    11/19/25      G LEVEL 21      MAIN          DATE = 82211    11/19/25
C   MAIN PROGRAM FOR SCATTERING FROM ASYMMETRICAL SUBREFLECTORS      CALL SETUP(NGL,I,JMAX,KMAX,MMAX,NMAX)
C   ASSUME SMOOTH PERFECTLY CONDUCTING SURFACE OF ARBITRARY      WRITE(6,2CQ2)TITLE
C   SHAPE      INPUT INCIDENT FIELD DATA
C   PROGRAM ORIGINALLY CUE TO A. LUDWIG • FEB. 1970      CALL FIELDS
C   CCCCCCCC      BEGIN LCOP FOR INTEGRATION GRID SEGMENTS
C   CCCCCCCC      DC 90 J=1,JMAX
C   COMMON/GRID1/SIT(36),COT(36),SIP(91),TCP(91),PP(46)      00080000 00090000 00600000 00530000
C   COMMON/GRID2/SIT(36),CCT(36),SIP(46),CPP(46),TT(36),PP(46)      00090000 00100000 00610000 00540000
C   DIMENSION F(36,91)*RHOFG(91)*TTCUT(31)*PHI(31)*TAMP(31)      00100000 00110000 00620000 00550000
C   DIMENSION F(36,91)*FT(36,91)*FP(36,91)*GAM(36,91)*TEDGE(91)      00110000 00120000 00630000 00560000
C   INTEGER TITLE,FPILOT,FSPLLOT      90 00130000 00140000 00640000 00570000
C   COMPLEX HRI36,*91*HT(36,91)*HP(36,91)      00150000 00160000 00650000 00580000
C   COMPLEX ETT(361,*46)*EPP(361,*46)      00170000 00180000 00660000 00590000
C   COMPLEX ETC(EPPC,EPLANE      00190000 00200000 00670000 00600000
C   COMPLEX A(36,91,*3)      00210000 00220000 00680000 00610000
C   COMPLEX STOT(3)      00230000 00240000 00690000 00620000
C   COMPLEX T1,T2,T3      00250000 00260000 00700000 00630000
C   COMPLEX TET,TEP      00270000 00280000 00710000 00640000
C   EQUIVALENCE (A1,1,1,1)*R1,(A1,1,2)*R2,(A1,1,3)*R3      00290000 00300000 00720000 00650000
C   EQUIVALENCE (GAM,FI)      00310000 00320000 00730000 00660000
C   T•P = THETA & PHI OF INTEGRATION GRID      00330000 00340000 00740000 00670000
C   TT•PP = BIGTHETA & BIGPHI OF OUTPUT GRID      00350000 00360000 00750000 00680000
C   COMMON PARAMETERS = SIN,COS, CF ABOVE ANGLES      00370000 00380000 00760000 00690000
C   F•FP = R•O AND ITS DERIVATIVES WRT TO THETA & PHI      00390000 00400000 00770000 00700000
C   HR•HT•HP = INCIDENT MAGNETIC FIELDS COMPONENTS      00410000 00420000 00780000 00710000
C   A = VECTOR RELATED TO SURFACE CURRENTS      00430000 00440000 00790000 00720000
C   STOT = VALUE OF RADIATION INTEGRAL FROM FIN      00450000 00460000 00800000 00730000
C   GAM = PATHLENGTH TERM (A FUNCTION OF BOTH GRIDS)      00470000 00480000 00810000 00740000
C   TEDGE = POLAR ANGLE VALUES WHICH SPECIFY REFLECTOR EDGE AS A      00490000 00500000 00820000 00750000
C   ETTC•EPP = SCATTERED FIELDS VALUES ON THE CUTPUT GRID      00510000 00520000 00830000 00760000
C   JMAX•NMAX = # BIGTHETA•#BIGPHI POINTS ON CUTPUT GRID      00530000 00540000 00840000 00780000
C   MMAX•IMAX = # THETA•#PHI POINTS ON PRESENT INTEGRATION GRID      00550000 00560000 00850000 00790000
C   NGL = # INTEG. GRIDS      00570000 00580000 00860000 00800000
C   ETC•EPC = FIELDS VALUES AT A SPECIFIC CUTPUT GRID POINT      00590000 00600000 00870000 00810000
C   READ(5,1C01)TITLE      00610000 00620000 00880000 00820000
C   WRITE(6,203)TITLE      00630000 00640000 00890000 00850000
C   READ(5,1C03)INCISK      00650000 00660000 00900000 00850000
C   WRITE INT DISK DATA SET      00670000 00680000 00910000 00860000
C   IF(INDISK .NE. 0)WRITE(12,1001)TITLE      00690000 00700000 00920000 00870000
C   INPUT PHASE TRANSLATIONS AND REFLECTOR PARAMETERS      00710000 00720000 00930000 00880000
C   READ(5,1C02)PC,X1,Y1,Z1,ALPHA,RHCO,THETAG      00730000 00740000 00940000 00890000
C   SCALE1•C      00750000 00760000 00950000 00900000
C   CONTINUE      00770000 00780000 00960000 00910000
C   CTR=0.017453293      00790000 00800000 00970000 00920000
C   WRITE(6,2003)PC      00810000 00820000 00980000 00930000
C   INPUT GRD DATA AND ESTABLISH CUTPUT GRD( #2) AND      00830000 00840000 00990000 00940000
C   FIRST INTEGRATION GRID( #1)      00850000 00860000 00950000 00953000
C   PRINT-CUT CCNSTANT THETA-EDGE      00870000 00880000 00960000 00951000
C   IF(EDGE .EQ. 0)WRITE(6,2019) (JJ,TEDGE(JJ),JJ=1,1)

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G LEVEL 21          MAIN          DATE = 82211    11/19/25    G LEVEL 21          MAIN          DATE = 82211    11/19/25
DO 30 N=2,NMAX
  TEDGE(N)=TEDGE(1)
CONTINUE
C
PRINT-CUT THETA-EDGE AS A FUNCTION OF PHI
C
IF(IEDGE .NE. 0) WRITE(6,2019)(JJ,TEDGE(JJ),JJ=1,NMAX)
29  WRITE(6,2018)
C
CONVERT THETA-EDGE TO RADIANS; CALCULATE RHO-EDGE
C
DO 31 N=1,NMAX
  TEDGE(N)=TEDGE(N)*0.01745329
  RHEDGE(N)=RHCC/(SIN(TEDGE(N))*
31
  OUTPUT RHO-EDGE
C
PRINT-CUT RHO-EDGE AS A FUNCTION OF PHI
C
WRITE(6,2019)(JJ,RHEDGE(JJ),JJ=1,NMAX)
28  CONTINUE
C
FIND AVERAGE-IN-PHI OF THETA-EDGE
C
TEDGEA=C.
DC 405 N=1,NMAX
  TEDGEA=TEDGE + TEDGE(N)
405  IF(IEDGE .EQ. 1)TEDGE(N)=THMAX
  TEDGE=TEDGE*FLOAT(NMAX)/0.017453293
C
  IF EDGE NOT TO BE ENCOUNTERED THIS GRID, SET TEDGEA TO THETAMAX
C
  IF(IEDGE .EQ. C)TEDGEA=THMAX
C
ESTABLISH SURFACE PARAMETERS!!.. RHO & ITS DERIVS.! ON
INTEGRATION GRID
C
CALL SURF1,MMAX,NMAX,F,FT,FP,IEDEG,TEGGE,RHO0,ALPHA)
410  CC 410 N=1,NMAX
  CC 411 N=1,NMAX
  F(M,N)=C*(F(M,N),
CONTINUE
C
ESTABLISH INCIDENT FIELDS ON INTEGRATION GRID BY EVALUATING
  NEAR-FIELD OF SWE
C
CALL FIELD1,MMAX,NMAX,HR,HT,HP,F)
C
READ IN PLOTTING PARAMETERS
C
REAC(5,1)CSINIPLOT,PHIPLT,FIPLT,FSPLT
C
PLCT NEAR-FIELD CF INCIDENT PATTERN
C
TT1=T(1)/CTR
  CC 430 N=1,NMAX
  TC=T(M)/DTR
  IF(N .NE. 1)GC TC 300
C
  SUM POWER CONTRIBUTIONS FRM EACH DELTA-THETA (TO THETA-EDGE)
  CC 340 N=2,NMAX
  CC 350 N=3,NMAX
  CC 360 N=4,NMAX
  CC 370 N=5,NMAX
  CC 380 N=6,NMAX
  CC 390 N=7,NMAX
  CC 400 N=8,NMAX
  CC 410 N=9,NMAX
  CC 420 N=10,NMAX
  CC 430 N=11,NMAX
  CC 440 N=12,NMAX
  CC 450 N=13,NMAX
  CC 460 N=14,NMAX
  CC 470 N=15,NMAX
  CC 480 N=16,NMAX
  CC 490 N=17,NMAX
  CC 500 N=18,NMAX
  CC 510 N=19,NMAX
  CC 520 N=20,NMAX
  CC 530 N=21,NMAX
  CC 540 N=22,NMAX
  CC 550 N=23,NMAX
  CC 560 N=24,NMAX
  CC 570 N=25,NMAX
  CC 580 N=26,NMAX
  CC 590 N=27,NMAX
  CC 600 N=28,NMAX
  CC 610 N=29,NMAX
  CC 620 N=30,NMAX
  CC 630 N=31,NMAX
  CC 640 N=32,NMAX
  CC 650 N=33,NMAX
  CC 660 N=34,NMAX
  CC 670 N=35,NMAX
  CC 680 N=36,NMAX
  CC 690 N=37,NMAX
  CC 700 N=38,NMAX
  CC 710 N=39,NMAX
  CC 720 N=40,NMAX
  CC 730 N=41,NMAX
  CC 740 N=42,NMAX
  CC 750 N=43,NMAX
  CC 760 N=44,NMAX
  CC 770 N=45,NMAX
  CC 780 N=46,NMAX
  CC 790 N=47,NMAX
  CC 800 N=48,NMAX
  CC 810 N=49,NMAX
  CC 820 N=50,NMAX
  CC 830 N=51,NMAX
  CC 840 N=52,NMAX
  CC 850 N=53,NMAX
  CC 860 N=54,NMAX
  CC 870 N=55,NMAX
  CC 880 N=56,NMAX
  CC 890 N=57,NMAX
  CC 900 N=58,NMAX
  CC 910 N=59,NMAX
  CC 920 N=60,NMAX
  CC 930 N=61,NMAX
  CC 940 N=62,NMAX
  CC 950 N=63,NMAX
  CC 960 N=64,NMAX
  CC 970 N=65,NMAX
  CC 980 N=66,NMAX
  CC 990 N=67,NMAX
  CC 1000 N=68,NMAX
  CC 1010 N=69,NMAX
  CC 1020 N=70,NMAX
  CC 1030 N=71,NMAX
  CC 1040 N=72,NMAX
  CC 1050 N=73,NMAX
  CC 1060 N=74,NMAX
  CC 1070 N=75,NMAX
  CC 1080 N=76,NMAX
  CC 1090 N=77,NMAX
  CC 1100 N=78,NMAX
  CC 1110 N=79,NMAX
  CC 1120 N=80,NMAX
  CC 1130 N=81,NMAX
  CC 1140 N=82,NMAX
  CC 1150 N=83,NMAX
  CC 1160 N=84,NMAX
  CC 1170 N=85,NMAX
  CC 1180 N=86,NMAX
  CC 1190 N=87,NMAX
  CC 1200 N=88,NMAX
  CC 1210 N=89,NMAX
  CC 1220 N=90,NMAX
  CC 1230 N=91,NMAX
  CC 1240 N=92,NMAX
  CC 1250 N=93,NMAX
  CC 1260 N=94,NMAX
  CC 1270 N=95,NMAX
  CC 1280 N=96,NMAX
  CC 1290 N=97,NMAX
  CC 1300 N=98,NMAX
  CC 1310 N=99,NMAX
  CC 1320 N=100,NMAX
  CC 1330 N=101,NMAX
  CC 1340 N=102,NMAX
  CC 1350 N=103,NMAX
  CC 1360 N=104,NMAX
  CC 1370 N=105,NMAX
  CC 1380 N=106,NMAX
  CC 1390 N=107,NMAX
  CC 1400 N=108,NMAX
  CC 1410 N=109,NMAX
  CC 1420 N=110,NMAX
  CC 1430 N=111,NMAX
  CC 1440 N=112,NMAX
  CC 1450 N=113,NMAX
  CC 1460 N=114,NMAX
  CC 1470 N=115,NMAX
  CC 1480 N=116,NMAX
  CC 1490 N=117,NMAX
  CC 1500 N=118,NMAX
  CC 1510 N=119,NMAX
  CC 1520 N=120,NMAX
  CC 1530 N=121,NMAX
  CC 1540 N=122,NMAX
  CC 1550 N=123,NMAX
  CC 1560 N=124,NMAX
  CC 1570 N=125,NMAX
  CC 1580 N=126,NMAX
  CC 1590 N=127,NMAX
  CC 1600 N=128,NMAX
  CC 1610 N=129,NMAX
  CC 1620 N=130,NMAX
  CC 1630 N=131,NMAX
  CC 1640 N=132,NMAX
  CC 1650 N=133,NMAX
  CC 1660 N=134,NMAX
  CC 1670 N=135,NMAX
  CC 1680 N=136,NMAX
  CC 1690 N=137,NMAX
  CC 1700 N=138,NMAX
  CC 1710 N=139,NMAX
  CC 1720 N=140,NMAX
  CC 1730 N=141,NMAX
  CC 1740 N=142,NMAX
  CC 1750 N=143,NMAX
  CC 1760 N=144,NMAX
  CC 1770 N=145,NMAX
  CC 1780 N=146,NMAX
  CC 1790 N=147,NMAX
  CC 1800 N=148,NMAX
  CC 1810 N=149,NMAX
  CC 1820 N=150,NMAX
  CC 1830 N=151,NMAX
  CC 1840 N=152,NMAX
  CC 1850 N=153,NMAX
  CC 1860 N=154,NMAX
  CC 1870 N=155,NMAX
  CC 1880 N=156,NMAX
  CC 1890 N=157,NMAX
  CC 1900 N=158,NMAX
  CC 1910 N=159,NMAX
  CC 1920 N=160,NMAX
  CC 1930 N=161,NMAX
  CC 1940 N=162,NMAX
  CC 1950 N=163,NMAX
  CC 1960 N=164,NMAX
  CC 1970 N=165,NMAX
  CC 1980 N=166,NMAX
  CC 1990 N=167,NMAX
  CC 2000 N=168,NMAX
  CC 2010 N=169,NMAX
  CC 2020 N=170,NMAX
  CC 2030 N=171,NMAX
  CC 2040 N=172,NMAX
  CC 2050 N=173,NMAX
  CC 2060 N=174,NMAX
  CC 2070 N=175,NMAX
  CC 2080 N=176,NMAX
  CC 2090 N=177,NMAX
  CC 2100 N=178,NMAX
  CC 2110 N=179,NMAX
  CC 2120 N=180,NMAX
  CC 2130 N=181,NMAX
  CC 2140 N=182,NMAX
  CC 2150 N=183,NMAX
  CC 2160 N=184,NMAX
  CC 2170 N=185,NMAX
  CC 2180 N=186,NMAX
  CC 2190 N=187,NMAX
  CC 2200 N=188,NMAX
  CC 2210 N=189,NMAX
  CC 2220 N=190,NMAX
  CC 2230 N=191,NMAX
  CC 2240 N=192,NMAX
  CC 2250 N=193,NMAX
  CC 2260 N=194,NMAX
  CC 2270 N=195,NMAX
  CC 2280 N=196,NMAX
  CC 2290 N=197,NMAX
  CC 2300 N=198,NMAX
  CC 2310 N=199,NMAX
  CC 2320 N=200,NMAX
  CC 2330 N=201,NMAX
  CC 2340 N=202,NMAX
  CC 2350 N=203,NMAX
  CC 2360 N=204,NMAX
  CC 2370 N=205,NMAX
  CC 2380 N=206,NMAX
  CC 2390 N=207,NMAX
  CC 2400 N=208,NMAX
  CC 2410 N=209,NMAX
  CC 2420 N=210,NMAX
  CC 2430 N=211,NMAX
  CC 2440 N=212,NMAX
  CC 2450 N=213,NMAX
  CC 2460 N=214,NMAX
  CC 2470 N=215,NMAX
  CC 2480 N=216,NMAX
  CC 2490 N=217,NMAX
  CC 2500 N=218,NMAX
  CC 2510 N=219,NMAX
  CC 2520 N=220,NMAX
  CC 2530 N=221,NMAX
  CC 2540 N=222,NMAX
  CC 2550 N=223,NMAX
  CC 2560 N=224,NMAX
  CC 2570 N=225,NMAX
  CC 2580 N=226,NMAX
  CC 2590 N=227,NMAX
  CC 2600 N=228,NMAX
  CC 2610 N=229,NMAX
  CC 2620 N=230,NMAX
  CC 2630 N=231,NMAX
  CC 2640 N=232,NMAX
  CC 2650 N=233,NMAX
  CC 2660 N=234,NMAX
  CC 2670 N=235,NMAX
  CC 2680 N=236,NMAX
  CC 2690 N=237,NMAX
  CC 2700 N=238,NMAX
  CC 2710 N=239,NMAX
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  CC 7550 N=723,NMAX
  CC 7560 N=724,NMAX
  CC 7570 N=725
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G LEVEL 21          MAIN          DATE = 82211    11/19/25    G LEVEL 21          MAIN          DATE = 82211    11/19/25
DC 33C M=1,NMAX          02090000
A3=REAL(H(M,N))          02100000
A4=IAVAG1(H(M,N))        02110000
A5=REAL(HP(M,N))         02111000
A6=IAVAG1(HP(M,N))       02112000
CALL VECTOR(A3,A4,HTEMP,HTPHI)
CALL VECTOR(A5,A6,HPAMP,HPPHI)
C
SUM SQUARE OF THETA & PHI FIELD COMPONENTS
C
EEM=HTEMP+HTAMP+HPAMP+HPHAMP
IF(M .EQ. 1)GO TO 325
C
POWER AT MIDPOINT OF ANULUS; MULTIPLY BY WIDTH OF ANULUS
C
POWM=(EEMM1+EEM)*SIN(T(M)-TCIFF)
TPCWX=TPCWX+POWM*DT
C
TAKE AVERAGE OF PAST & PRESENT THETA-EDGE VALUES
C
TEAVG=.5*(TEGE(N)+TEGE(N-1))
IF(T(M) .LT. TEAVG) GO TO 320
C
SCALE POWER CONTRIBUTIONS AT EDGE IF TEAVG LIES BETWEEN THETAS
SCALH=(TEAVG-T(M-1))/(T(M)-T(M-1))
IF(SCALH .LE. .01)GO TO 325
POWX=POWX+SCALH*PCWN*DT
GO TO 325
CONTINUE
POWX=POWX+POWN*DT
CONTINUE
EEMM1=EEM
320
CONTINUE
FNMXM1=FLOAT(NMAX-1)
C
SCALE AZIPUTHAL CONTRIBUTION
TPCW1=TPCWX/FNMXM1*DP+TPOW1
POW1=POWX/FNMXM1*DP+PCWI
C
PRINT-CUT AVERAGE THETA-EDGE AND POWER THRU GRID SEGMENT
C
WRITE(6,2304) TEDGEA,PCWI
C
PRINT-CUT MAX THETA AND PCWER THRU MAX THETA (I.E.,TPOW1)
C
WRITE(6,2307) THMAX,TPCW1
RAT=POW1/TPCW1*100.0
WRITE(6,2308) RAT
C
COMBINE SURFACE AND FIELD DATA TO DETERMINE VECTOR A
A CONTAINS THREE COORDINATE COMPONENTS.
C
DO 400 M=1,NMAX
DC 400 N=1,NMAX
C
        DATE = 82211    11/19/25
T1=FT(M,N)*SIT(M)*HP(M,N)*HT(M,N)
T2=FP(M,N)*HR(M,N)*F(M,N)*SIT(M)*HP(M,N)/PC
T3=-FT(M,N)*SIT(M)*HR(M,N)-F(M,N)*SIT(M)*HP(M,N)/PC
C
        CONVERT FROM SPHERICAL TO CARTESIAN COORDINATES
C
A(M,N,1)=T1*SIT(M)*CGP(N)+T2*COT(M)*COP(N)-T3*SIP(N)
A(M,N,2)=T1*SIT(M)*SIP(N)+T2*COT(M)*SIP(N)+T3*COP(N)
A(M,N,3)=T1*CCT(M)-T2*SIT(M)
CONTINUE
C
        BEGIN OUTPUT GRID LOOP. FOR EACH POINT ON OUTPUT GRID DETERMINE
SCATTERED FIELD CONTRIBUTION FROM ENTIRE INTEG GRID.
C
        WRITE(6,2C07) I
C
        400
        01792000
        01793000
        01800000
        01810000
        01811000
        01812000
        01813000
        01820000
        01830000
        01831000
        01832000
        01833000
        01840000
        01850000
        01851000
        01852000
        01853000
        01860000
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        01940000
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        01960000
        01961000
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        01970000
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        01982000
        01983000
        01990000
        01991000
        01992000
        01993000
        02000000
        02010000
        02020000
        02030000
        02040000
        02050000
        02060000
        02070000
        02080000
        01730000
        01740000
        01750000
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        01790000
        01791000
        01792000
        01793000
        01800000
        01810000
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        01991000
        01992000
        01993000
        02000000
        02010000
        02020000
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        02060000
        02070000
        02080000
        0171=SIT(M)*CGP(N)+T2*COT(M)*COP(N)-T3*SIP(N)
        0172=SIT(M)*SIP(N)+T2*COT(M)*SIP(N)+T3*COP(N)
        0173=CCT(M)-T2*SIT(M)
        0174=CONTINUE
        0175=WRITE(6,2011)TO
        0176=DO 500 J=1,JMAX
        0177=00 500
        0178=00 500
        0179=00 500
        0180=00 500
        0181=00 500
        0182=00 500
        0183=00 500
        0184=00 500
        0185=00 500
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        0197=00 500
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        0199=00 500
        0200=00 500
        0201=00 500
        0202=00 500
        0203=00 500
        0204=00 500
        0205=00 500
        0206=00 500
        0207=00 500
        0208=00 500
        0170=CONTINUE
        0171=CONVERT FROM SPHERICAL TO CARTESIAN COORDINATES
        0172=A(M,N,1)=T1*SIT(M)*CGP(N)+T2*COT(M)*COP(N)-T3*SIP(N)
        0173=A(M,N,2)=T1*SIT(M)*SIP(N)+T2*COT(M)*SIP(N)+T3*COP(N)
        0174=A(M,N,3)=T1*CCT(M)-T2*SIT(M)
        0175=CONTINUE
        0176=WRITE(6,2011)TO
        0177=DO 500 J=1,JMAX
        0178=00 500
        0179=00 500
        0180=00 500
        0181=00 500
        0182=00 500
        0183=00 500
        0184=00 500
        0185=00 500
        0186=00 500
        0187=00 500
        0188=00 500
        0189=00 500
        0190=00 500
        0191=00 500
        0192=00 500
        0193=00 500
        0194=00 500
        0195=00 500
        0196=00 500
        0197=00 500
        0198=00 500
        0199=00 500
        0200=00 500
        0201=00 500
        0202=00 500
        0203=00 500
        0204=00 500
        0205=00 500
        0206=00 500
        0207=00 500
        0208=00 500
        0170=CONTINUE
        0171=CALL PATHL(F,J,K,MMAX,NMAX,GAM)
        0172=PERFORM INTEGRATION
        0173=CALL FINIT(T,P,A,GAM,MMAX-1,NMAX-1,STOT,I,J,K,PC,IEDEG,IEEDGE)
        0174=ASSIGN SCATTERED FIELDS AT OUTPUT POINT
        0175=ETTC=COIT(J)*(STCT(1)*COPP(K)+STOT(2)*SIPP(K))-STGT(3)
        0176=J*SIT(M)
        0177=EPP0=STOT(2)*COPP(K)-STOT(1)*SIPP(K)
        0178=ETTO=-10.0*1.0)*PC/6.*2831854*ETTO
        0179=EPPC=-10.0*1.0)*PC/6.*2831854*EPPC
        0180=TC=TT1)/C.017453293
        0181=AL=REAL(ETTC)
        0182=A2=A1WG(ETTO)
        0183=A3=REAL(EPPC)
        0184=A4=A1WG(EPPC)
        0185=CONVERT TC POLAR COORDINATES
        0186=CALL VECTOR1,A2*ETAMP,EPHFI
        0187=CALL VECTOR1,A3*4*EPAMP,EPHFI
        0188=WRITE(6,2012)TO
        0189=CONTINUE
        0190=SUPERIMPOSE SCATTERED FIELD VALUE WITH VALUES CORRESPONDING
        0191=TC OTHER INTEG GRIDS.
        0192=ETT(J,K)=ETT1(J,K)*ETTO
        0193=EPP(J,K)=EPP1(J,K)*EPPC
        0194=CONTINUE
        0195=COMPUTE AVERAGE THETA-EDGE FOR ALL INTEGRATION GRIDS. SKIP
        0196=COMPUTATION IF THETA-EDGE NOT ON PRESENT GRID SEGMENT
        0197=02160000
        0198=02170000
        0199=02180000
        0200=02190000
        0201=02200000
        0202=02210000
        0203=02220000
        0204=02230000
        0205=02240000
        0206=02250000
        0207=02260000
        0208=02270000
        0209=02280000
        0210=02290000
        0211=02300000
        0212=02310000
        0213=02320000
        0214=02330000
        0215=02340000
        0216=02350000
        0217=02360000
        0218=02370000
        0219=02380000
        0220=02390000
        0221=02400000
        0222=02410000
        0223=02420000
        0224=02430000
        0225=02440000
        0226=02450000
        0227=02460000
        0228=02470000
        0229=02480000
        0230=02490000
        0231=02500000
        0232=02510000
        0233=02520000
        0234=02530000
        0235=02540000
        0236=02550000
        0237=02560000
        0238=02570000
        0239=02580000
        0240=02590000
        0241=02600000
        0242=02610000
        0243=02620000
        0244=02630000
        0245=02640000
        0246=02650000
        0247=02660000
        0248=02670000
        0249=02680000
        0250=02690000
        0251=02700000
        0252=02710000
        0253=02720000
        0254=02730000
        0255=02740000
        0256=02750000
        0257=02760000
        0258=02770000
        0259=02780000
        0260=02790000
        0261=02800000
        0262=02810000
        0263=02820000
        0264=02830000
        0265=02840000
        0266=02850000
        0267=02860000
        0268=02870000
        0269=02880000
        0270=02890000
        0271=02900000
        0272=02910000
        0273=02920000
        0274=02930000
        0275=02940000
        0276=02950000
        0277=02960000
        0278=02970000
        0279=02980000
        0280=02990000
        0281=02991000
        0282=02992000
        0283=02993000
        0284=02994000
        0285=02995000
        0286=02996000
        0287=02997000
        0288=02998000
        0289=02999000
        0290=02999100
        0291=02999200
        0292=02999300
        0293=02999400
        0294=02999500
        0295=02999600
        0296=02999700
        0297=02999800
        0298=02999900
        0299=02999910
        0300=02999920
        0301=02999930
        0302=02999940
        0303=02999950
        0304=02999960
        0305=02999970
        0306=02999980
        0307=02999990
        0308=02999991
        0309=02999992
        0310=02999993
        0311=02999994
        0312=02999995
        0313=02999996
        0314=02999997
        0315=02999998
        0316=02999999
        0317=029999991
        0318=029999992
        0319=029999993
        0320=029999994
        0321=029999995
        0322=029999996
        0323=029999997
        0324=029999998
        0325=029999999
        0326=0299999991
        0327=0299999992
        0328=0299999993
        0329=0299999994
        0330=0299999995
        0331=0299999996
        0332=0299999997
        0333=0299999998
        0334=0299999999
        0335=02999999991
        0336=02999999992
        0337=02999999993
        0338=02999999994
        0339=02999999995
        0340=02999999996
        0341=02999999997
        0342=02999999998
        0343=02999999999
        0344=029999999991
        0345=029999999992
        0346=029999999993
        0347=029999999994
        0348=029999999995
        0349=029999999996
        0350=029999999997
        0351=029999999998
        0352=029999999999
        0353=0299999999991
        0354=0299999999992
        0355=0299999999993
        0356=0299999999994
        0357=0299999999995
        0358=0299999999996
        0359=0299999999997
        0360=0299999999998
        0361=0299999999999
        0362=02999999999991
        0363=02999999999992
        0364=02999999999993
        0365=02999999999994
        0366=02999999999995
        0367=02999999999996
        0368=02999999999997
        0369=02999999999998
        0370=02999999999999
        0371=029999999999991
        0372=029999999999992
        0373=029999999999993
        0374=029999999999994
        0375=029999999999995
        0376=029999999999996
        0377=029999999999997
        0378=029999999999998
        0379=029999999999999
        0380=0299999999999991
        0381=0299999999999992
        0382=0299999999999993
        0383=0299999999999994
        0384=0299999999999995
        0385=0299999999999996
        0386=0299999999999997
        0387=0299999999999998
        0388=0299999999999999
        0389=02999999999999991
        0390=02999999999999992
        0391=02999999999999993
        0392=02999999999999994
        0393=02999999999999995
        0394=02999999999999996
        0395=02999999999999997
        0396=02999999999999998
        0397=02999999999999999
        0398=029999999999999991
        0399=029999999999999992
        0400=029999999999999993
        0401=029999999999999994
        0402=029999999999999995
        0403=029999999999999996
        0404=029999999999999997
        0405=029999999999999998
        0406=029999999999999999
        0407=0299999999999999991
        0408=0299999999999999992
        0409=0299999999999999993
        0410=0299999999999999994
        0411=0299999999999999995
        0412=0299999999999999996
        0413=0299999999999999997
        0414=0299999999999999998
        0415=0299999999999999999
        0416=02999999999999999991
        0417=02999999999999999992
        0418=02999999999999999993
        0419=02999999999999999994
        0420=02999999999999999995
        0421=02999999999999999996
        0422=02999999999999999997
        0423=02999999999999999998
        0424=02999999999999999999
        0425=029999999999999999991
        0426=029999999999999999992
        0427=029999999999999999993
        0428=029999999999999999994
        0429=029999999999999999995
        0430=029999999999999999996
        0431=029999999999999999997
        0432=029999999999999999998
        0433=029999999999999999999
        0434=0299999999999999999991
        0435=0299999999999999999992
        0436=0299999999999999999993
        0437=0299999999999999999994
        0438=0299999999999999999995
        0439=0299999999999999999996
        0440=0299999999999999999997
        0441=0299999999999999999998
        0442=0299999999999999999999
        0443=02999999999999999999991
        0444=02999999999999999999992
        0445=02999999999999999999993
        0446=02999999999999999999994
        0447=02999999999999999999995
        0448=02999999999999999999996
        0449=02999999999999999999997
        0450=02999999999999999999998
        0451=02999999999999999999999
        0452=029999999999999999999991
        0453=029999999999999999999992
        0454=029999999999999999999993
        0455=029999999999999999999994
        0456=029999999999999999999995
        0457=029999999999999999999996
        0458=029999999999999999999997
        0459=029999999999999999999998
        0460=029999999999999999999999
        0461=0299999999999999999999991
        0462=0299999999999999999999992
        0463=0299999999999999999999993
        0464=0299999999999999999999994
        0465=0299999999999999999999995
        0466=0299999999999999999999996
        0467=0299999999999999999999997
        0468=0299999999999999999999998
        0469=0299999999999999999999999
        0470=02999999999999999999999991
        0471=02999999999999999999999992
        0472=02999999999999999999999993
        0473=02999999999999999999999994
        0474=02999999999999999999999995
        0475=02999999999999999999999996
        0476=02999999999999999999999997
        0477=02999999999999999999999998
        0478=02999999999999999999999999
        0479=029999999999999999999999991
        0480=029999999999999999999999992
        0481=029999999999999999999999993
        0482=029999999999999999999999994
        0483=029999999999999999999999995
        0484=029999999999999999999999996
        0485=029999999999999999999999997
        0486=029999999999999999999999998
        0487=029999999999999999999999999
        0488=0299999999999999999999999991
        0489=0299999999999999999999999992
        0490=0299999999999999999999999993
        0491=0299999999999999999999999994
        0492=0299999999999999999999999995
        0493=0299999999999999999999999996
        0494=0299999999999999999999999997
        0495=0299999999999999999999999998
        0496=0299999999999999999999999999
        0497=02999999999999999999999999991
        0498=02999999999999999999999999992
        0499=02999999999999999999999999993
        0500=02999999999999999999999999994
        0501=02999999999999999999999999995
        0502=02999999999999999999999999996
        0503=02999999999999999999999999997
        0504=02999999999999999999999999998
        0505=02999999999999999999999999999
        0506=029999999999999999999999999991
        0507=029999999999999999999999999992
        0508=029999999999999999999999999993
        0509=029999999999999999999999999994
        0510=029999999999999999999999999995
        0511=029999999999999999999999999996
        0512=029999999999999999999999999997
        0513=029999999999999999999999999998
        0514=029999999999999999999999999999
        0515=0299999999999999999999999999991
        0516=0299999999999999999999999999992
        0517=0299999999999999999999999999993
        0518=0299999999999999999999999999994
        0519=02
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25

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G LEVEL 21      MAIN          DATE = 82211    11/19/25      G LEVEL 21      MAIN          DATE = 82211    11/19/25
C   IF(ITT1.*GE. 29.91GO TO 755
C   PERC=100.*EPOW/TPOW
C   WRITE(6,2301)TDEGAE,EPOW,PERC
C   RAT=POMI/EPOW
C
C   PRINT-CUT RATIO OF NEAR TC FAR-FIELD INCIDENT POWER
C
C   WRITE(6,2306)TDEGAE,RAT
C   CONTINUE
C
C   TRANSLATE PHASE CENTER,SCALE FIELD AMPLITUDES,AND OUTPUT
C   TOTAL FIELDS
C   WRITE(6,2002)TITLE
C   WRITE(6,2010)XT,Y1,ZT,SCALE
C
C   PARAMETERS USED IN WRITING SCATTERED FIELDS ON DISK
C
C   TTPEAK=180.0-2.*ALPHA
C   JP_EAK=INT((TTPEAK-ITT1)/DTT + .1) + 1
C
C   READ-IN HALF-ANGLE FOR POWER CALCULATION
C
C   READ(5,1002)TDEGAE
C
C   DO 795 K=1,KMAX
C   PC=PP1K/0.017453293
C   IF(FSPLOT .EQ. -1) WRITE(6,2011)PO
C   IF(FSPLOT .EQ. 0) WRITE(6,2004)PO
C   TPOW=0.
C   EPOW=0.
C
C   TRANSLATE ANGLES TOWARDS 0.0 DEG AT BEAM PEAK
C
C   TDIFF=1180.-2.*ALPHA+5*DFTT1*DTR
C   DO 760 J=1,JMAX
C   T0=TT1(J)/C.017453293
C   SIGNTT=0.0
C   IF(COTT1(JJ) *EQ. 0.00000)GO TO 765
C   SIGNTT=COTT1(JJ)/ABS(COTT1(JJ))
C   CONTINUE
C
C   EVALUATE Y COMPONENT OF SCATTERED FIELD
C
C   EPLANE=EFTT(J,K)*SIPP(K)*SIGNTT + EPP(J,K)*COSP(K)
C
C   A1=REAL(EFTT(J,K))
C   A2=AIMAG(EFTT(J,K))
C   A3=REAL(EPP(J,K))
C   A4=AIMAG(EPP(J,K))
C   A5=REAL(EPLANE)
C   A6=AIMAG(EPLANE)
C
C   CALL VECTOR(A1,A2*ETAMP*EPHI)
C   CALL_VECTOR(A3,A4*EPAMP*EPHI)
C   CALL_VECTOR(A5,A6*EPAMP*EPHI)
C
C   TRANSLATE PHASE CENTER OF SCATTERED PATTERN
C
C   DATE = 82211      11/19/25
C
C   033300000
C   033350000
C   034000000
C   033600000
C   033610000
C
C   ADJUST PHASES TO -180,180 RANGE
C
C   033630000
C   033700000
C   033710000
C
C   CALL ADJUST(EPHI)
C   CALL_ADJUST(EPHI)
C   CALL ADJUST(EPHI)
C
C   033800000
C   033900000
C   034000000
C   034100000
C   034200000
C   034300000
C   034400000
C   034500000
C   034600000
C   034700000
C   034800000
C   034900000
C   035000000
C
C   SCALE FIELD AMPLITUDES
C
C   ETAMP=ETAMP*SCALE
C   EPAMP=EPAMP*SCALE
C   EYAMP=EYAMP*SCALE
C
C   IF(FSPLOT .EQ. -1) GO TO 780
C   IF(J .NE. 1) GO TO 770
C   IF(FSPLOT .EQ. 1) WRITE(6,2106)PO
C   IF(FSPLOT .EQ. 2) WRITE(6,2107)PO
C   IF(FSPLOT .EQ. 3) WRITE(6,2108)PO
C
C   CONTINUE
C   IF(FSPLOT .EQ. 1) CALL EPILOT(J,EYAMP,EYPHI,JMAX*DTT,TT1,T0)
C   IF(FSPLOT .EQ. 2) CALL EPILOT(J,ETAMP,EPHI,JMAX*DTT,TT1,T0)
C   IF(FSPLOT .EQ. 3) CALL EPILOT(J,EPAMP,EPHI,JMAX*DTT,TT1,T0)
C
C   034010000
C   034100000
C   034200000
C   034300000
C   034400000
C   034500000
C   034630000
C   034760000
C   034890000
C   034920000
C   035050000
C   035180000
C   035200000
C   035330000
C
C   STORE 31 SCATTERED PATTERN VALUES IN THETA*AMP & PHASE ARRAYS.
C   CHOOSE EITHER LEFT OR RIGHT SIDE OF BEAM. TRANSLATE ANGLES
C
C   0353200
C   0353300
C   0354000
C   0355000
C   0356000
C   0357000
C   0358000
C   0359000
C   0360000
C   0360200
C   0360300
C   0361000
C   0363000
C   0364000
C   0365000
C
C   IF(INDISK .EQ. 0) GO TO 790
C   IF(INDISK .EQ. 2) GO TO 791
C   IF((J .LT. (JPEAK-30)) .OR. (J .GT. JPEAK)) GO TO 790
C   L=JPEAK-J+1
C   TTOUT(L)=TPEAK-T0
C   TAMP(L)=EYAMP
C   TPHIL=EPHI
C   GC TO 790
C   CONTINUE
C   IF((J .LT. JPEAK) .OR. (J .GT. (JPEAK+30))) GO TO 790
C   L=J-JPEAK+
C   TTOUT(L)=T-TPEAK
C   TAMP(L)=EPAMP
C   TPHIL=EPHI
C   CCNTINUE
C
C   PRINT-CUT SCATTERED PATTERN VALUES
C
C   IF(FSPLOT .EQ. -1) WRITE(6,2102)TO,ETAMP
C   ,EPHI,EPAMP,EPHI
C
C   CALCULATE POWER IN SCATTERED PATTERN
C
C   EEJ=ETAMP*ETAMP + EPAMP*EPAMP
C   IF(J .EQ. 1) GC TO 785
C
C   0367000
C   0368000
C   0369000
C   0370000
C   0371000
C   0371200
C   0371300
C
C   DATE = 82211      11/19/25
C
C   037200000
C   037300000
C   037400000
C   037500000
C   037600000
C   037700000
C   037800000
C   037900000
C   038000000
C   038100000
C   038200000
C   038230000
C   038300000
C   038400000
C   038500000
C   038600000
C   038700000
C   038800000
C   038900000
C   039000000
C   039100000
C   039200000
C   039300000
C   039400000
C   039500000
C   039600000
C   039800000
C   039910000
C   039900000
C   039950000
C   040000000
C   040100000
C   040213000
C   040214000
C   040215000
C   040216000
C   040217000
C   040218000
C   040219000
C   040250000
C   040300000
C   040310000
C   040320000
C   040400000
C   040500000
C   040700000
C   040800000
C   041000000
C   041010000

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G LEVEL	21	MAIN	DATE = 82211	11/19/25	G LEVEL	21	MAIN	DATE = 82211	11/19/25
C	POWER AT CENTER OF ANULUS. ADD TO TOTAL POWER		04102000	2004	FORMAT(0*,*10X,*PHI= *F6*2* DEGREES*)		04380000		
C	$PCW_j = 5*ABS((EEJ+EEJM1)*SIN(TT(j)-TDIFF))$		04103000	2006	FORMAT(1H0,* SCATTERED FIELDS FROM GRID*12)		04390000		
	$TPOW = TPCW*PCW_j;DTT$		04110000	2007	FORMAT(1H0,* BEGIN INTEGRATION OVER GRID*12)		04400000		
C	TRANSLATE THETA TO 0 DEG AT BEAM PEAK;< 0 DEG BEFORE PEAK;		04120000	2009	FORMAT(1H1,* DIRECT RADIATION FROM INCIDENT FIELDS*)		04400000		
C	>0 DEG AFTER PEAK.		04121000	2010	FORMAT(1H0,*SUPERPOSITION OF ALL GRID SCATTERED FIELDS AND DIRECT RADIATION*)		04420000		
C	$\epsilon_{SH} Z= *F8*4*/*$		04122000	2011	*. FIELDS /*/*, PHASECENTER TRANSLATED BY X= *F8*4*, Y= *F8*4*, Z= *F8*4*/		04430000		
C	ϵ_{SH} AMPLITUDE VALUES SCALED BY FACTOR OF *F4*21		04123000	2012	*. FORMAT(1H0, PHI=*F7*2*/		04440000		
C	FORMAT(1H0, PHI=*F7*2*/		04124000	2013	E17X*7E THE THETA,15X*5HE PHI/*		04450000		
C	FORMAT(1H0, PHI=*F7*2*/		04125000	2014	E9H THE TA,2*120H VOLTS PHASE*)		04470000		
C	FORMAT(1H0, PHI=*F7*2*/		04130000	2015	FORMAT(F9*2*F11*6*F8*2*F12*6*F8*2)		04490000		
C	FORMAT(F10*2*F10*6*F10*2*F10*6*F10*2)		04134000	2016	FORMAT(F10*2*F10*6*F10*2*F10*6*F10*2)		04500000		
C	RHO-EDGE */		04135000	2017	FORMAT(1H0,*N RHO-EDGE *)		04510000		
C	FORMAT(1H0,*N THE TA EDGE (DEG.1),		04136000	2018	FORMAT(1H1,*15*1X*F8*41)		04520000		
C	FORMAT(1H0,*N THE TA EDGE (DEG.1),		04136200	2019	FORMAT(1H0,*N THE TA EDGE (DEG.1),		04530000		
C	FORMAT(1H0,*N THE TA EDGE (DEG.1),		04136300	2020	FORMAT(1H0,*N THE TA EDGE (DEG.1),		04540000		
C	FORMAT(1H0,*N THE TA EDGE (DEG.1),		04136400	2021	FORMAT(1H0,*N THE TA EDGE (DEG.1),		04560000		
C	FORMAT(1H0,*N THE TA EDGE (DEG.1),		04136500	2022	FORMAT(1H0,*N THE TA EDGE (DEG.1),		04562000		
C	FORMAT(1H0,*N THE TA EDGE (DEG.1),		04136600	2023	FORMAT(1H0,*N THE FAR-FIELD INCIDENT POWER THRU *F6*2* DEGREES*,		04570000		
C	FORMAT(1H0,*N THE FAR-FIELD INCIDENT POWER THRU *F6*2* DEGREES*,		04136700	2024	C IS *F6*4*, VOLTS SQUARED = *F6*2*, % OF TOTAL POWER *,		04580000		
C	FORMAT(1H0,*N THE FAR-FIELD INCIDENT POWER THRU *F6*2* DEGREES*,		04136800	2025	FORMAT(1H0,*N THE RATIO OF FAR-FIELD SCATTERED POWER TO NEAR-FIELD*,		04590000		
C	FORMAT(1H0,*N THE RATIO OF FAR-FIELD SCATTERED POWER TO NEAR-FIELD*,		04136900	2026	C, INCIDENT POWER THRU *F6*2*, DEGREES IS *F6*4*)		04600000		
C	FORMAT(1H0,*N THE NEAR-FIELD INCIDENT POWER THRU *F6*2*,		04137000	2027	FORMAT(1H0,*N THE NEAR-FIELD INCIDENT POWER THRU *F6*2*,		04610000		
C	FORMAT(1H0,*N THE NEAR-FIELD INCIDENT POWER THRU *F6*2*,		04137100	2028	C, DEGREES IS *F6*4*, SQUARE VOLTS*)		04620000		
C	FORMAT(1H0,*N THE NEAR-FIELD INCIDENT POWER THRU *F6*2*,		04137200	2029	FORMAT(1H0,*N THE FAR-FIELD SCATTERED POWER THRU *F6*2*,		04630000		
C	FORMAT(1H0,*N THE FAR-FIELD SCATTERED POWER THRU *F6*2*,		04137300	2030	C IS *F6*4*, VOLTS SQUARED = *F6*2*, % OF TOTAL POWER *,		04640000		
C	FORMAT(1H0,*N THE FAR-FIELD SCATTERED POWER THRU *F6*2*,		04137400	2031	C, INCIDENT POWER THRU *F6*2*, DEGREES IS *F6*4*)		04650000		
C	FORMAT(1H0,*N THE FAR-FIELD SCATTERED POWER THRU *F6*2*,		04137500	2032	FORMAT(1H0,*N THE RATIO OF NEAR-FIELD TO FAR-FIELD INCIDENT POWER *,		04660000		
C	FORMAT(1H0,*N THE RATIO OF NEAR-FIELD TO FAR-FIELD INCIDENT POWER *,		04137600	2033	C THRU *F6*2*, DEGREES IS *F6*4*)		04662000		
C	FORMAT(1H0,*N THE NEAR-FIELD INCIDENT POWER THRU *F6*2*,		04137700	2034	FORMAT(1H0,*N THE NEAR-FIELD INCIDENT POWER THRU *F6*2*,		04670000		
C	FORMAT(1H0,*N THE NEAR-FIELD INCIDENT POWER THRU *F6*2*,		04137800	2035	C, DEGREES IS *F6*4*, TOTAL POWER IS *F6*4*, SQUARE VOLTS*)		04680000		
C	SQUARE VOLTS*)		04137900	2036	FORMAT(1H0,*N THE POWER STRIKING THE REFLECTOR IS *F6*2*, % OF *		04690000		
C	THE POWER STRIKING THE REFLECTOR IS *F6*2*, % OF *		04138000	2037	E, THE TOTAL POWER*)		04700000		
C	E, THE TOTAL POWER*)		04200000	2100	FORMAT(1H1,*26X*, RADIAL COMPONENT OF (NEAR) INCIDENT MAGNETIC*)		04710000		
C	FORMAT(1H1,*26X*, RADIAL COMPONENT OF (NEAR) INCIDENT MAGNETIC*)		04210000	2101	E, FIELD VS. POLAR ANGLE *14X* PHI= *F6*1* DEG*)		04720000		
C	E, FIELD VS. POLAR ANGLE *14X* PHI= *F6*1* DEG*)		04212000	2102	FORMAT(1H1,*26X*, THETA COMPONENT OF (NEAR) INCIDENT MAGNETIC*)		04730000		
C	FORMAT(1H1,*26X*, THETA COMPONENT OF (NEAR) INCIDENT MAGNETIC*)		04220000	2103	E, FIELD VS. POLAR ANGLE *14X* PHI= *F6*1* DEG*)		04740000		
C	E, FIELD VS. POLAR ANGLE *14X* PHI= *F6*1* DEG*)		04230000	2104	FORMAT(1H1,*28X*, THETA COMPONENT OF (FAR) INCIDENT MAGNETIC*)		04750000		
C	FORMAT(1H1,*28X*, THETA COMPONENT OF (FAR) INCIDENT MAGNETIC*)		04240000	2105	E, FIELD VS. POLAR ANGLE *14X* PHI= *F6*1* DEG*)		04760000		
C	E, FIELD VS. POLAR ANGLE *14X* PHI= *F6*1* DEG*)		04250000	2106	FORMAT(1H1,*35X*, Y-COMP. OF (FAR) INCIDENT ELECTRIC FIELD VS. *)		04770000		
C	FORMAT(1H1,*35X*, Y-COMP. OF (FAR) INCIDENT ELECTRIC FIELD VS. *)		04260000	2107	E, POLAR ANGLE *14X* PHI= *F6*1* DEG*)		04780000		
C	E, POLAR ANGLE *14X* PHI= *F6*1* DEG*)		04270000	2108	FORMAT(1H1,*26X*, THETA COMPONENT OF (FAR) SCATTERED ELECTRIC *)		04790000		
C	FORMAT(1H1,*26X*, THETA COMPONENT OF (FAR) SCATTERED ELECTRIC *)		04280000	2109	E, FIELD VS. POLAR ANGLE *14X* PHI= *F6*1* DEG*)		04800000		
C	E, FIELD VS. POLAR ANGLE *14X* PHI= *F6*1* DEG*)		04290000	2110	FORMAT(1H1,*29X*, PHI COMPONENT OF (FAR) INCIDENT ELECTRIC FIELD*)		04810000		
C	FORMAT(1H1,*29X*, PHI COMPONENT OF (FAR) INCIDENT ELECTRIC FIELD*)		04300000	2111	E, VS. POLAR ANGLE /*14X* PHI= *F6*1* DEG*)		04820000		
C	E, VS. POLAR ANGLE /*14X* PHI= *F6*1* DEG*)		04310000	2112	FORMAT(1H1,*32X*, Y-COMP. OF (FAR) SCATTERED ELECTRIC FIELD VS. *)		04830000		
C	FORMAT(1H1,*32X*, Y-COMP. OF (FAR) SCATTERED ELECTRIC FIELD VS. *)		04320000	2113	E, POLAR ANGLE /*14X* PHI= *F6*1* DEG*)		04840000		
C	E, POLAR ANGLE /*14X* PHI= *F6*1* DEG*)		04330000	2114	FORMAT(1H1,*28X*, THETA COMPONENT OF (FAR) SCATTERED ELECTRIC *)		04850000		
C	FORMAT(1H1,*28X*, THETA COMPONENT OF (FAR) SCATTERED ELECTRIC *)		04340000	2115	E, FIELD VS. POLAR ANGLE /*14X* PHI= *F6*1* DEG*)		04860000		
C	E, FIELD VS. POLAR ANGLE /*14X* PHI= *F6*1* DEG*)		04350000	2116	FORMAT(1H1,*27X*, PHI COMPONENT OF (FAR) SCATTERED ELECTRIC *)		04870000		
C	FORMAT(1H1,*27X*, PHI COMPONENT OF (FAR) SCATTERED ELECTRIC *)		04360000	2117	E, FIELD VS. POLAR ANGLE /*14X* PHI= *F6*1* DEG*)		04880000		
C	E, FIELD VS. POLAR ANGLE /*14X* PHI= *F6*1* DEG*)		04370000	2118	FORMAT(1H1,*27X*, PHI COMPONENT OF (FAR) SCATTERED ELECTRIC *)		04890000		
C	FORMAT(1H1,*27X*, PHI COMPONENT OF (FAR) SCATTERED ELECTRIC *)				MAIN IS FOLLOWED BY EPLT.				
1001	FORMAT(18A4)				END				
1002	FORMAT(1ICF10*4)								
1003	FORMAT(15IS1)								
1005	FORMAT(15*F7*2*215)								
1004	FORMAT(1LX*F9*5)								
2001	FORMAT(1L1*, NRAC								
2002	FORMAT(1LH*5X*, L8A4)								
2003	FORMAT(*C*, PROPAGATION CONSTANT= *F10*5)								