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JPL PHYSICAL OPTICS SCATTERING PROGRAM

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I. INTRODUCTION

This report describes a FORTRAN program developed at JPL for the calculation of the scattered pattern from a reflector of arbitrary shape. The program written by A. Ludwig is contained in a report entitled "Calculation of Scattered Patterns from Asymmetrical Reflectors" (Technical Report 32-1430; 2/15/70).

The JPL report stresses the following: the use of a Fourier expansion representation of the reflector surface, the possible inaccuracy of a far field assumption for the fields incident on the reflector, and the use of a fast integration technique developed by Ludwig. Some good examples of the accuracy of the program are included. A copy of the JPL report may be obtained from Rick Fisher in Green Bank or from Sarah Martin (microfiche) in Charlottesville.

The program is actually two programs: one which calculates a scattered pattern and one which computes, using a spherical wave expansion, a completely general representation of the incident field pattern. The programs will be referred to as SCAT and SWE, respectively.

The method of surface specification has been altered to that of tabular form and the integration routine has been slightly adjusted to better handle reflector edge contributions. No major changes have been made to the SWE program.

SCAT employs the technique of physical optics to calculate the scattered field. The program takes the Fourier transform, using a fast integration technique, of an array of induced current dipoles.

The incident field representation provided by the SWE program allows accurate calculations to be made even with the reflector well into the near field of the source. Input into the SWE program is the far field source pattern, such as those typically measured on an antenna range. The far field form of

spherical waves is matched to the (far field) incident pattern. Backing up into the near field, the SWE accurately approximates the transverse and radial components of the electric field.

(An attempt has been made towards the end of this report to describe in physical terms some of the theory behind the SCAT and SWE programs.)

One application of this program is in checking the taper at the edge of the main dish of the 140-ft Cassegrain telescope. For this problem the scattering surface is the subreflector, which has an asymmetrical edge, and the incident field is the antenna range (or computed) pattern of the appropriate Cassegrain feed.

A slight variation of this situation involves the re-design of the subreflector to include a conical vertex plate and a flange about the reflector edge.

Another application is to aid in the design of beam polarization splitters which may be added to the 140-ft. To the SCAT program, the beam splitter would be a tilted plane reflector placed very near the feed aperture.

II. GENERAL PROGRAM CHARACTERISTICS

A. Summary of What Program Does

1. SCAT program.

Given a scattering surface specified by $\rho(\theta, \phi)$ (i.e., spherical coordinates) and the magnetic fields incident on this surface, the scattered pattern is computed over a grid of observation points. The scattering surface is assumed to be perfectly conducting; therefore, the electric field is zero at the surface.

The values of ρ are either input in tabular form or calculated from an equation in θ and/or ϕ .

The E and H plane values (dB or volts, and degrees) of the incident magnetic field are read in by the SWE program. The SWE program computes the coefficients of the expansion and stores them. These stored values are read in by the SCAT programs.

Ludwig has shown that assuming far field conditions for the source field for the trivial case of scattering from an infinite plane reflector located within or near the traditional far field boundary line ($2D^2/\lambda$, D = source diameter) can result in rather strong backlobes in the scattered pattern. The strength of the backlobe and distortion of the main (reflected) lobe increases as the reflector further penetrates the near field; in other words, the magnitude of the error depends on frequency. (See JPL report, figures 3-5 on pages 6-8.) Ludwig further shows (Figure 7) that use of the spherical wave representation almost entirely eliminates the problem.

The scattered fields are, as mentioned above, computed using the method of physical optics. It is assumed that the induced currents are zero on shadowed portions of the reflector and that on directly illuminated areas the induced surface current value is twice the value of the tangential (to the surface) incident H field. These two assumptions constitute the physical optics approximations.

As a means of reducing computation time the far field form of the scattered pattern is calculated; the near field form can be found by submitting the scattered field values to the SWE program.

Summing the product of the induced surface currents and phase delay (path-length) for each Δs of the surface, the scattered field as seen from a particular point on the output grid is determined. This integration is performed for each point on the output grid. Finally, the far field source pattern is added to the scattered pattern yielding what is termed the "total fields". For the

situation of an infinite plane reflector, the total fields would be the incident fields pointed in the opposite direction on the output grid and with a 180° phase change.

From comparisons of computed and measured subreflector patterns, Ludwig has shown the computed pattern to be accurate down to -35 dB (relative power) and even through the first side lobe. Fairly reliable spillover efficiency calculations should be possible, since the major spillover contribution is generally from fields within 25 dB of the pattern maximum. Proper handling of the side lobes requires the use of the Geometrical Theory of Diffraction (GTD) which does not make the assumption of zero induced currents on shadowed reflector surfaces.

2. SWE program.

The SWE program finds the coefficients of expansion of the incident (far zone) magnetic field pattern in terms of transverse electric (TE) and transverse magnetic (TM) spherical waves. The TE and TM spherical waves are the general vector solutions to Maxwell's equations for an electromagnetic wave travelling in a source-free region, V . For the case of representing a magnetic field, the TM vector solution has no radial components, whereas the TE solution has components in all three coordinate directions.

A spherical wave expansion operates under the same mathematical principles as the Fourier expansion of a function. In a Fourier expansion, a function is represented by a summation of sines and cosines each multiplied by a coefficient particular to the order of sine and cosine variation; this order of variation could be called the mode order. To represent the function, the two sets of coefficients (one set multiplying the sines, the other the cosines) must be determined. This is done by evaluating integrals involving sines and cosines and specific function values. The integrals are simplified through the use of the orthogonality relations between sines and cosines.

In a spherical wave expansion the sines and cosines of the Fourier expansion are replaced with the TE and TM spherical wave solutions of Maxwell's equations. The integrals for the mode coefficients contain the corresponding TE and TM mode functions and the far field incident pattern values. Thus, the procedure for determining the coefficients of the spherical wave expansion is basically the same as for that of the Fourier expansion.

The expansion is in three variables, in this case the spherical coordinates, and contains two mode orders. Thus, the expansion is a double summation; one summation for each mode variable. One mode order, generally called the mode order and designated by a "n", specifies the degree of ρ variation. The other mode order, called the order of azimuthal variation and designated by an "m", specifies the degree of ϕ variation. The degree of θ variation depends on both mode orders.

In principle, the expansion scheme outlined above can represent any input pattern at any distance from its source. However, because the case of $m = 1$ is of particular importance, the double summation has been reduced in the program to a single summation over n. Also, one set of spherical wave solutions has been neglected, which means that the incident radiation is assumed to be linearly polarized.

B. Coordinate System Used.

The origin of the system will be at the phase center of the source fields. The reflecting surface is specified by the vector $\bar{\rho}$ which is a function of the angles θ and ϕ ; ρ , θ , ϕ represent a point on the surface in spherical coordinates. Since $\bar{\rho}$ is defined by θ and ϕ , these two angles will specify the point of integration.

The output grid over which the observer views the scattered pattern is also designated by spherical coordinates (R, Θ , Φ). Since the far field form of the scattered pattern is computed, the output grid is defined entirely by Θ and Φ .

The Z-axis is the reflector axis. The three coordinate sets — (ρ, θ, ϕ) , (R, Θ, Φ) and (X, Y, Z) — all have the same origin. The coordinate system is shown in Figure 1.

C. Program Structure and Order of Calculation.

1. SCAT program.

For storage reasons, the reflector surface over which the integration is performed is divided into integration grids. The scattered fields from each grid are superimposed and added to the incident fields to yield the total (scattered) pattern.

The reflector is divided along θ and/or ϕ (Figure 2), depending on the size of the reflector, the choice of $\Delta\theta$ and $\Delta\phi$ and storage requirements. This partitioning introduces no appreciable error and is quite useful since the total integration may cover several thousand points.

The output grid is not segmented because it generally includes only several hundred points at most.

For each integration grid, $\rho(\theta, \phi)$ is either read in or calculated for every (θ, ϕ) point on the grid. Since the normal to the surface is required at each integration point, $\frac{\partial\rho}{\partial\theta}(\theta, \phi)$ and $\frac{\partial\rho}{\partial\phi}(\theta, \phi)$ are computed over the grid. The derivatives are computed using the numerical method of backward and forward differences or by an analytic expression in θ and ϕ which is inserted in subroutine SURF.

In general the outermost (in θ) integration grids will contain points that are beyond the edge of the reflector. This is so because in the program, θ and ϕ are independent of each other; that is, the integration grid is represented by two vectors rather than by a matrix of points. However, at the reflector edge, the boundary is defined by θ_{edge} , which will generally be a function of ϕ . To correctly represent the reflector, the edge values of θ are read in or calculated for each ϕ of the integration grid.

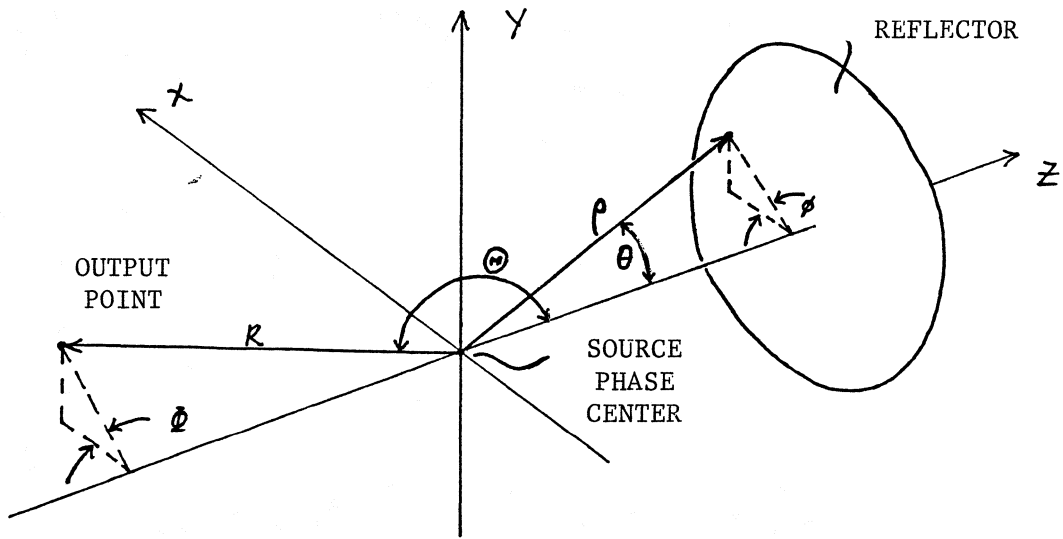


FIGURE 1

Coordinate system (from JPL report, page 3).

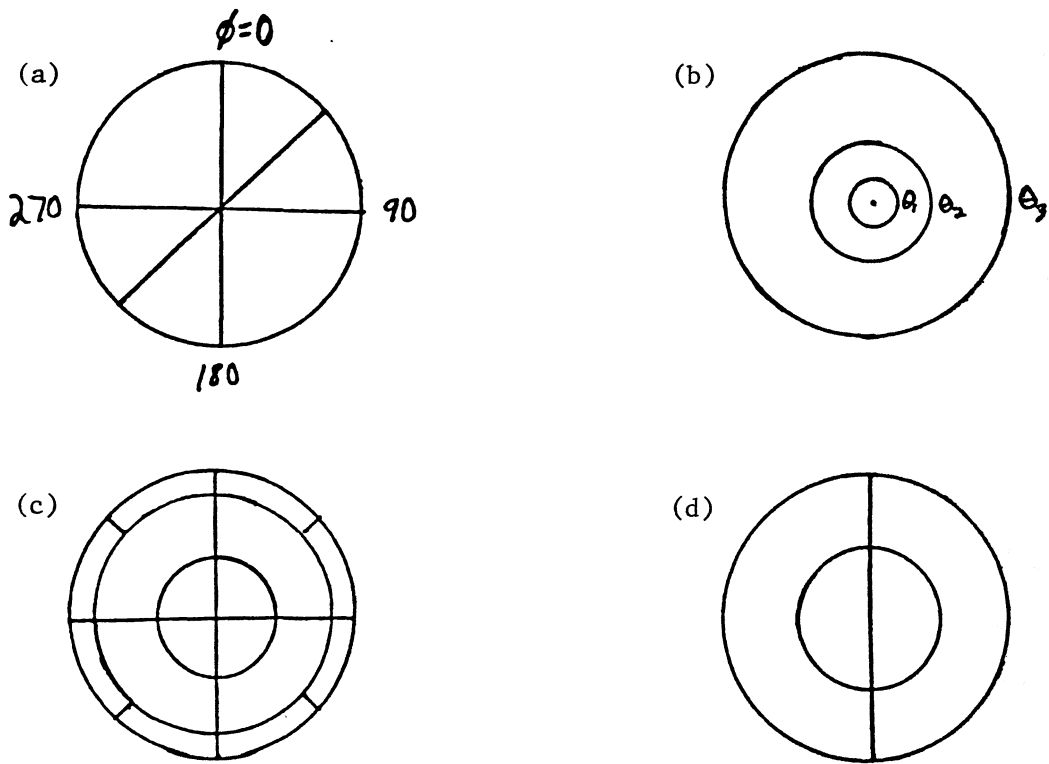


FIGURE 2

Sample integration grids.

An integer parameter, IEDGE, is used to determine if the edge will be encountered for the present integration grid and, if it is encountered, how the edge values will be obtained. IEDGE is read in by the MAIN subroutine of the SCAT program for each integration grid.

As mentioned above, the edge integration grids will contain points beyond the actual reflector. Since such values may not be known, they are not required inputs. When the values of $\rho(\theta, \phi)$ are in tabular form, the user can set to 999.99999 the first ρ value that corresponds to a point beyond the reflector edge; the ρ values remaining for that given value of ϕ can (and must) be set to some arbitrary value (e.g., zero).

If all the values of ρ for an edge integration grid are known, as is the case when ρ is calculated analytically, then including them will result in slightly better accuracy.

Note that the SCAT program is at its maximum user-convenience when ρ (in particular) and the edge values of θ can be computed from equations inserted in the program.

An informal flowchart of the SCAT program is shown in Figure 3. Subroutine names are written next to most flowchart components.

2. SWE program.

The organization of the SWE program is considerably simpler; therefore, no flowchart of it is shown. A mode order term and the far field E and H plane pattern values from the appropriate feed comprise the inputs. Details of input considerations will be given below.

After reading in the input pattern, the program establishes a matrix of θ mode weights (forms of Legendre polynomials) and a matrix of input pattern differentials. The evaluation of the coefficient integrals is then done by matrix multiplication. The real and imaginary values of the TE and TM coefficients are normalized and stored on the computer disk. As a check on the validity of the coefficients, the far field form is computed allowing a comparison to be made with the input pattern.

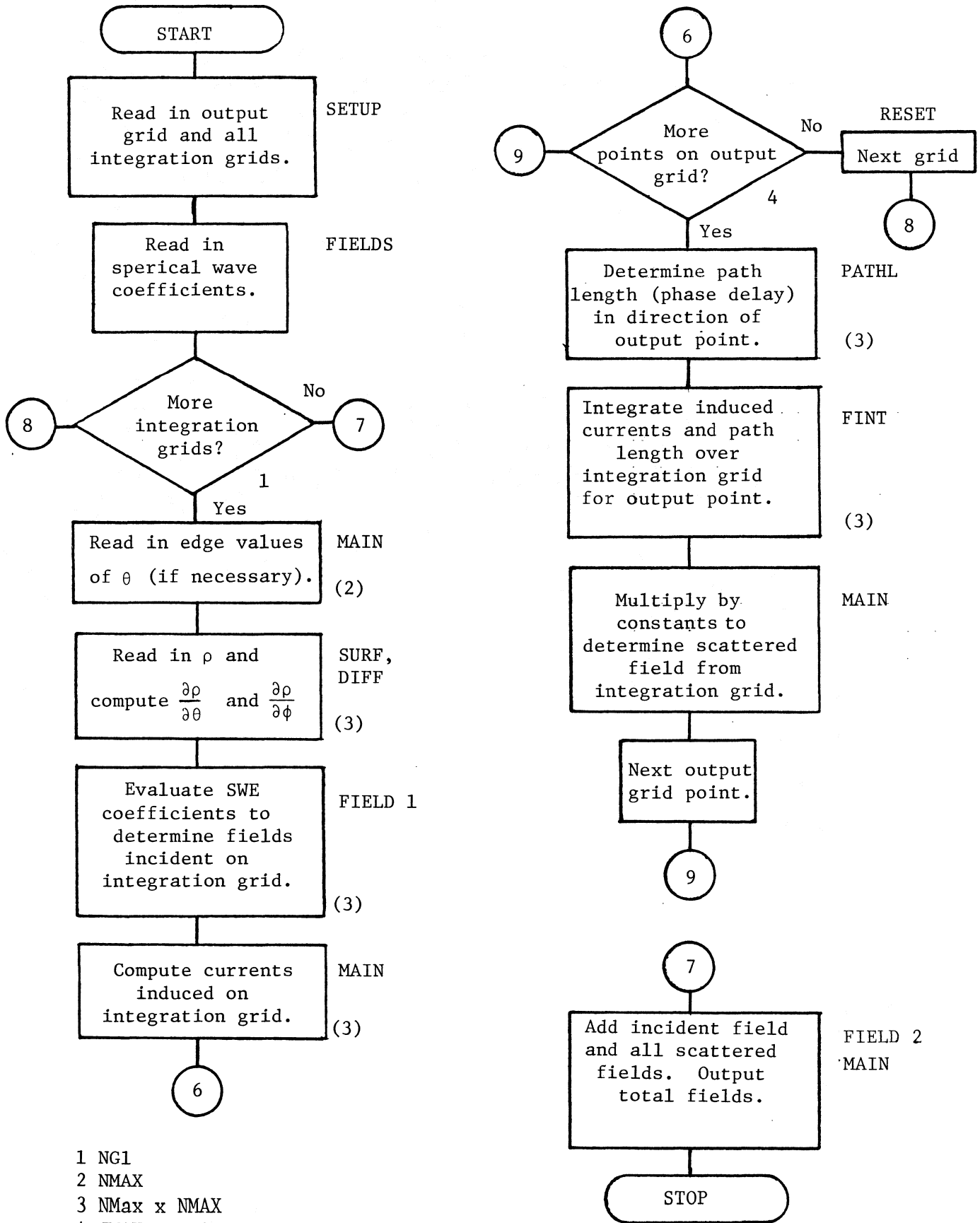


FIGURE 3

SCAT program flowchart.

III. DETAILS OF USING THE PROGRAM

Both programs are stored on the disk of the IBM 360 computer located in Charlottesville. The user command system is called Pandora. Each subroutine is stored as a Pandora member and has a member name either identical to or similar to the subroutine name. The Pandora member name (PMN) will be given for each subroutine.

For both the SCAT program and the SWE program this section provides the following information:

For each subroutine of the program:

Subroutine description, input variables (that READ data), input considerations, and a table of inputs and their formats.

Discussion of program output.

A. Subroutines of Scattering Program

1. MAIN (PMN MAIN1SC)

Subroutine description:

This routine coordinates the efforts of the program. The procedure followed by MAIN has been outlined by the SCAT program flowchart (Figure 3). Note the SUBROUTINE names written by the flowchart blocks.

Input variables and their tasks:

TITLE - any alphanumeric statement ≤ 72 characters.

PC - phase constant, $k = 2\pi/\lambda$ (λ in meters).

XT, YT, ZT - translations to expected phase center of scattered pattern (affects only output phase data).

SCALE - scale factor for output fields. If input is zero, program sets SCALE = 1.0.

(The following variables are read in for each integration grid.)

IEDGE - an integer; determines if and how program obtains θ at the reflector edge. There are four cases:

IEDGE (continued):

IEDGE < 0 - Program calculates values for that integration grid using equations inserted into subroutine EDGEEQ.

IEDGE = 0 - θ at edge is constant. Program reads in just one value.

IEDGE = 11 - Edge will not be encountered on present integration grid. No values read in.

IEDGE > 0 - θ at edge is a function of ϕ ($\neq 11$) and is read in for each value of ϕ in the integration grid (NMAX values).

TEDGE (NMAX) - edge values of θ in degrees. Either NMAX values, one value or no values are read in depending on IEDGE.

Input considerations:

TITLE usually describes reflector and frequency. The scattered field phase center translations may be set to zero if the output phase is of no interest.

The choice of $\Delta\theta$ and $\Delta\phi$ (which determines the number of integration grids) is covered in SETUP.

Table of inputs and their formats:

Card	Input Variable(s)	Format
1	TITLE	18A4
2	PC, XT, YT, ZT, SCALE	5F10.4
	[SETUP DATA]	
	[FIELDS DATA]	
	IEDGE	I5
	TEDGE (Either NMAX, 1 or no values.)	7(1X, F9.5)
	[SURF DATA]	

Read in
for each
integra-
tion
grid.

2. Setup (PMN SETUP)

Subroutine description:

This routine sets up the output grid and all integration grids. The integration grid parameters are reset for each new grid by ENTRY RESET within SETUP.

The input parameters of the output grid consist of the number of values (JMAX and KMAX), initial value and increment which define the θ and ϕ points over which the fields are viewed. A typical grid could be $\phi = 0^\circ$ and 90° and $\theta = -90^\circ$ to $+90^\circ$ with a step of 2° . Thus, $KMAX = 2$, $JMAX = 91$.

The input parameters of each input grid are identical to those of the output grid but define θ and ϕ instead. As an example, consider Figure 2(d). The first grid might be $\theta = 0^\circ$ to 6° with $\Delta\theta = 0.2^\circ$ and $\phi = 0^\circ$ to 180° with $\Delta\phi = 4^\circ$. Thus, $MMAX = 31$ and $NMAX = 46$. The next grid might be $\theta = 0^\circ$ to 6° and $\phi = 180^\circ$ to 360° . Two more grids would follow. The order of the four grids is of no consequence.

Input variables and their tasks:

JMAX - number of θ values on output grid, ≤ 181 .
 TT1 - initial θ value.
 DDT - θ increment.
 KMAX - number of ϕ values on output grid, ≤ 5 .
 PP1 - initial ϕ value.
 DPP - ϕ increment.
 NG1 - number of integration grids, ≤ 21 .
 MM(I) - number of θ values on Ith integration grid, ≤ 36 .
 TI - initial θ value.
 DT - θ increment ($\Delta\theta$).
 NN(I) - number of ϕ values on Ith integration grid, ≤ 91 .
 PI - initial ϕ value.
 DP - ϕ increment ($\Delta\phi$).

(All initial values and increments in degrees.)

Input considerations:

$\Delta\theta$ and $\Delta\phi$ are chosen such that the dimensions of the largest ΔS on the reflector surface are roughly one wavelength by one wavelength. This ensures that the pathlength term does not vary by more than 2π over any ΔS . Thus, $\Delta\theta$ can be found from $\Delta\theta \approx \lambda/\rho_{\max}$, where ρ_{\max} is the greatest distance to the reflector and $\Delta\theta$ is in radians. The corresponding equation for $\Delta\phi$ is $\Delta\phi \approx \Delta\theta/\sin(\theta_{\max})$, where θ_{\max} is the largest θ angle subtended by the reflector edge.

The values obtained for $\Delta\theta$ and $\Delta\phi$ should be rounded to the nearest convenient value (tenths or hundredths of a degree). An upper limit of $\Delta\theta = 1^\circ$ is necessary, since larger values sample the incident field too sparsely. Decreasing $\Delta\theta$ and/or $\Delta\phi$ offers a good check on the validity of the chosen increments. However, halving either increment nearly doubles the computation time. To give some idea of the CPU time needed for a given JMAX, KMAX, TMMAX, TNMAX combination (define TMMAX and TNMAX as total number of θ and ϕ points on reflector surface), the following table is provided.

<u>KMAX</u>	<u>JMAX</u>	<u>TMMAX</u>	<u>TNMAX</u>	<u>CPU TIME</u> <u>(minutes)</u>	
1	91	35	92	18	
1	91	35	184	36	IBM
1	181	35	184	50	360/65
2	181	35	184	89	

To ensure accurate edge contributions, the user should be certain that for any value of ϕ the edge value of θ does not fall within the θ range of the first two values of θ (inclusive) on the integration grid. For most cases, the edge value of θ could be within this range and would produce no noticeable differences in the total scattered pattern. It is, however, absolutely necessary that all edge θ values be greater than or equal to the first θ of integration.

The integration routine and SURF subroutine set the limit of MMAX and NMAX ≥ 4 .

Table of inputs and their formats:

<u>Card</u>	<u>Input Variable(s)</u>	<u>Format</u>
1	JMAX, TT1, DTT	I5, 2F10.2
2	KMAX, PP1, DPP	I5, 2F10.2
3	NG1	I5
4	MM(1), T1, DT	I5, 2F10.2
5	NN(1), P1, DP	I5, 2F10.2
⋮	⋮	⋮
2NG1 + 2	MM(NG1), T1, DT	I5, 2F10.2
2NG1 + 3	NN (NG1), P1, DP	I5, 2F10.2

3. FIELDS (PMN FIELDS)

Subroutine FIELDS reads, off the computer disk, the spherical wave coefficients of the expansion of the incident fields. As mentioned previously, these coefficients are written on the disk by the SWE program.

FIELDS contains two entry points, ENTRY FIELD1 and ENTRY FIELD2. FIELD1 evaluates the SWE coefficients to find the near field form of the pattern incident on the reflector.

FIELD2 determines the electric field at an infinite distance from its source. FIELD2 is used in calculation of the far field form of the incident pattern.

Both FIELD1 and FIELD2 use the SWE coefficients in their computations.

Input variables:

- TITLE - alphanumeric statement, ≤ 70 characters.
- LMAX - maximum mode order, ≤ 70 .
- MCOMP - order of azimuthal variation.
- A(N,1), A(N,2) - real and imaginary components of $TE_{MCOMP,N}$ (N = 1 to LMAX) spherical wave coefficient.
- B(N,1), B(N,2) - real and imaginary components of $TM_{MCOMP,N}$ spherical wave coefficient.

Input considerations:

The user need not include the above information with the other data for the SCAT program. Instead, FIELDS reads this data directly off the computer disk. The first three variables are inputs of the SWE program; the coefficients are calculated by it. The input considerations and table of formats for these variables are given in the SWE program description.

4. SURF (PMN SSURF).

Subroutine description:

For each integration grid, this subroutine reads in all $\rho(\theta, \phi)$ values and calculates $\frac{\partial \rho}{\partial \theta}(\theta, \phi)$ and $\frac{\partial \rho}{\partial \phi}(\theta, \phi)$. There are four input cases that SURF considers; these are determined by the integer input parameter ISURF. The four cases are the following: ρ is a function of θ and ϕ and is read in from a table of values; ρ is a function of θ and is read in from a table of values; ρ is known in analytic forms as a function of θ and/or ϕ ; ρ , $\frac{\partial \rho}{\partial \theta}$ and $\frac{\partial \rho}{\partial \phi}$ are all known in analytic form. In the first situation, MMAX x NMAX values are read in (every point on the integration grid); in the second case, MMAX values are read in (every θ point on integration grid); in the remaining cases, no values are read in. When $\frac{\partial \rho}{\partial \theta}$ and $\frac{\partial \rho}{\partial \phi}$ are not analytically specified, they are numerically determined using the techniques of forward and backward differences. (See Computer Methods for Science and Engineering, LaFara.) This method of computing derivatives at the points of a table of data has been found to often be more accurate than is necessary for the SCAT program. For differentiable surfaces and not overly varying integration grids (as outlined in SETUP), this accuracy is maintained down to tables of four values. Thus, MMAX and NMAX should be greater than or equal to four.

To compute ρ and/or its partial derivatives, equations in θ and/or ϕ are inserted into SURF. The values of $\sin \theta$, $\cos \theta$, $\sin \phi$, $\cos \phi$, θ and ϕ (radians)

are stored for each integration grid in the variables SIT(M), COT(M), SIP(N), COP(N), T(M), P(N).

As previously mentioned, when ρ is input in tabular form, a value of 999.99999 indicates the reflector edge has just been passed. For a given ϕ , the remaining θ points on the integration grid must be assigned some arbitrary ρ value, which does not figure in any computations. If ρ is known for points beyond the edge, it would be slightly more accurate to include the value immediately beyond (on the integration grid) the edge value. Doing this would allow ρ at the edge to be interpolated rather than extrapolated.

Input variables:

ISURF - integer parameter that determines how ρ is obtained.

There are four cases as follows:

ISURF < 0 - ρ , $\frac{\partial \rho}{\partial \theta}$ and $\frac{\partial \rho}{\partial \phi}$ all determined analytically.
($\neq -2$)

Read no values.

ISURF = -2 - ρ determined analytically; $\frac{\partial \rho}{\partial \theta}$ and $\frac{\partial \rho}{\partial \phi}$ numerically computed (backward and forward differences). Read no values.

ISURF = 0 - ρ is a function of θ only and is read (MMAX values) from a table.

ISURF > 0 - ρ is a function of θ and ϕ and is read (MMAX x NMAX values) from a table.

F (M, N) - $\rho(\theta, \phi)$

FT (M,N) - $\frac{\partial \rho}{\partial \theta}(\theta, \phi)$

FP (M,N) - $\frac{\partial \rho}{\partial \phi}(\theta, \phi)$

Input considerations:

If the partial derivatives of ρ are to be determined numerically, MMAX and NMAX can at minimum be two but it is suggested that they be no less than four.

A value of ρ is required for every (θ, ϕ) or θ (depending on ISURF) on the integration grid.

For the case of reading in ρ for each (θ, ϕ) point on the integration grid, the θ DO loop is within the ϕ DO loop. Thus, for each ϕ value a new block of data is started.

The above two restrictions must be kept in mind when setting a value of ρ to 999.99999.

If at all possible, ρ should be represented by an equation. For this situation the input procedure is at its simplest.

Table of inputs and their formats:

	<u>Card</u>	<u>Variable(s)</u>	<u>Format</u>
Read in for each integra- tion grid.	1	ISURF (For ISURF < 0, no values read.)	I5
	2	F(1, 1) to F(MMAX, 1)	7(1X, F9.5)
	.	or to F(MMAX, NMAX)	
	.		

5. DIFF (PMN DIFF)

Subroutine description:

This subroutine is called by SURF to numerically approximate the derivative of a dependent variable with respect to the independent variable at each point in a table of data. The technique of forward and backward differences is employed; forward near the front of the table and backward elsewhere. By reaching back or forth into the table of data and taking the differences of the dependent variables along the way, the method is able to accurately compute derivatives of rapidly changing data within the table. The spacing of the independent variable ($\Delta\theta, \Delta\phi$) must be constant throughout the table (integration grid). The table should contain a minimum of four values.

There are no input variables.

6. FINT (PMN EFINT, FINT)

This subroutine performs the integration of the product of the induced surface currents and the pathlength over each integration grid for each point on the output grid. It is here in the program where the great majority of number crunching is performed. This integration routine was developed by Ludwig specifically for the form of integral encountered in scattering problems. (The fast Fourier transform is not applicable.) Ludwig's method reduces by a factor of 16, relative to a Simpson's rule integration, the number of integration points needed.

In this integration routine the θ loop is embedded within the ϕ loop.

The reflector edge is handled in the following way: If IEDGE indicates an edge is present, then for a given azimuthal angle FINT compares each $\theta + \Delta\theta$ value with the edge value of θ (which is a function of ϕ). When the edge is encountered, the edge θ for ϕ and for $\phi + \Delta\phi$ are averaged. The contribution from this smaller (in some cases, slightly larger) ΔS is then evaluated by extrapolating or interpolating the pathlength and surface currents to the averaged edge θ . (See Figure 4.). The reflector, in effect, has a discontinuous edge.

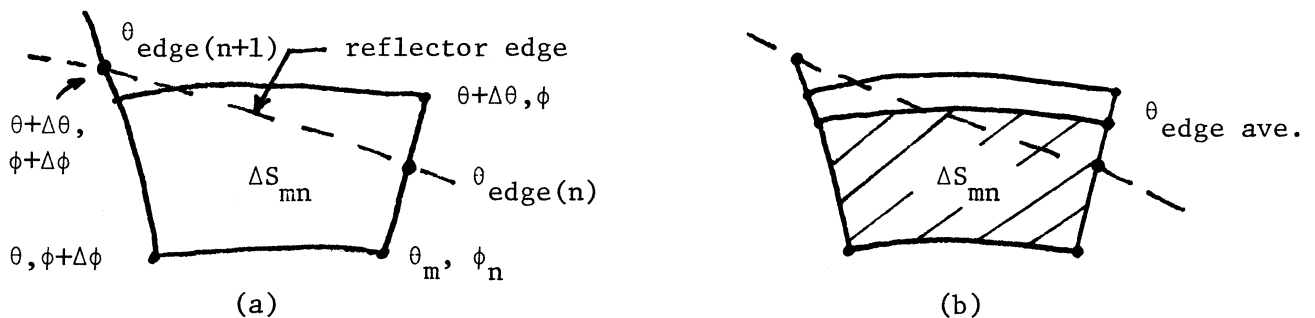


FIGURE 4

Another routine has been developed with the intent of better representing the edge. The ΔS_{mm} in Figure 4(a) is divided along ϕ , the number of divisions being determined by the difference between the two edge values of θ . θ_{edge} is then interpolated to each ϕ division. The same technique of θ averaging is applied over the reduced edge section, yielding smaller edge discontinuities. This integration routine is stored in PMN EFINT.

7. PATHL (PMN PATHL).

Subroutine description:

This member computes the pathlength from every point on the integration surface in the direction of each point on the output grid. The pathlength is the dot product between the vector ρ and the unit vector in the direction of the output variables.

8. SPHANK (PMN SPHANK).

Subroutine description:

SPHANK computes the term $h_n^{(2)}(k\rho) \cdot (-j)^{n+1} \cdot k\rho \cdot e^{jk\rho}$, where $h_n^{(2)}(k\rho)$ is the spherical Hankel function of the second kind, n is the mode order, and $j = \sqrt{-1}$.

This subroutine is called by FIELD1 (magnetic field at finite distance from source). The value SPHANK determines provides the near field ρ dependence. For far field calculations, this factor must be removed. Otherwise, SPHANK would provide only the value of $h_n^{(2)}(k\rho)$.

The spherical Hankel function represents an outward travelling spherical wave. $h_n^{(2)}(\chi)$ is defined as $j_n(\chi) - j y_n(\chi)$ where j_n and y_n are spherical Bessel functions. $j_n(\chi)$ and $y_n(\chi)$ are analogous to $\sin \chi$ and $\cos \chi$, and $h_n^{(2)}(\chi)$ is analogous to $e^{-j\chi}$.

SPHANK provides the only (known) numerical program limitation. For $\chi < 100$ (roughly), $y_n(\chi)$ blows up as n increases. As χ decreases, the blow up

occurs for smaller n . This is not a computing problem but is a characteristic of the spherical Bessel function, y_n . (See Handbook of Mathematical Functions, Abramowitz and Stegun; p. 438, Fig. 10.1 to 10.3 and pp. 465-466, Table 10.5.)

To avoid computer overflow SPHANK stops calculation when $y_n \sim 10^7$. For a wavelength of 5 cm ($k = 125.7 \text{ m}^{-1}$) and a plane reflector at a distance of 1 m, this truncation produced no changes in the scattered field phase or amplitude. It is suggested that the problem be such that $kp > 100$. For $kp < 100$, the above tables should be consulted to determine the maximum "safe" mode order.

There are no input variables.

9. LEGEND (PMN LEGEND)

Subroutine description:

LEGEND computes the value of the associated Legendre polynomial, $P_n^m(\cos \theta)$, where m is the order of azimuthal variation and n is the mode order.

The Legendre polynomial provides the correct polar angle (θ and Θ) variation of the spherical wave field representation. Therefore, it is called by both FIELD1 and FIELD2.

LEGEND has no numerical difficulties analogous to those of SHPANK. (See Tables of Functions, Johnke and Emde; pp. 112-113.)

There are no input variables.

10. VECTOR (PMN VECTOR)

Subroutine description:

This subroutine converts complex numbers in rectangular form (real and imaginary) to polar form (amplitude and phase). Most program computations are done with numbers in rectangular form.

There are no input variables.

11. ADJUST (PMN ADJUST)

Subroutine description:

ADJUST normalizes phase angles to the range of -180° to $+180^\circ$.

There are no input variables.

12. EDGEEQ (within PMN SSURF)

Subroutine description:

Equations for the edge value of theta are inserted into this subroutine. The same values of $\sin \phi$, $\cos \phi$, etc., used in SURF are available. The edge equations will be functions of ϕ only.

If necessary, the equations can be divided up for the different ϕ integration grids.

There are no input variables.

B. Output of Scattering Program.

The printout of the program is fairly self-explanatory. Briefly, the bulk of the output is the following:

-----	..	-- Output grid of all integration grids.
Output	..	-- Spherical wave coefficients.
for	..	-- If edge is to be encountered, all edge values of theta.
each	..	-- A selection of the values of ρ , $\frac{\partial \rho}{\partial \theta}$ and $\frac{\partial \rho}{\partial \phi}$.
inte-	..	-- Scattered fields (E and H planes of far-electric field)
gra-	..	from integration grid.
tion	..	-- Far field incident pattern.
grid.	..	-- Superposition of incident fields and all grid scattered
-----	..	fields (i.e., yields total scattered fields).

C. Subroutines of Spherical Wave Expansion Program.

The SWE program determines the coefficients of the spherical wave expansion of an input far field pattern, such as those typically measured on an antenna range. In principle, this expansion is completely general and provides one with the near field values, including radial components, of the input pattern. (For a mathematical account of spherical wave theory, see Electromagnetic Theory, Stratton; Chapter 7.)

1. MAIN (PMN MAINOR)

Subroutine description:

This routine is essentially the entire program. It reads in the far field pattern, computes the coefficients, writes the coefficients into a disk data set, and computes the far field form of the spherical wave expansion (for comparison with the input pattern).

Input variables:

TITLE	-	Alphanumeric statement, ≤ 72 characters (identify program).
MCOMP	-	Order of azimuthal variation.
LMAX	-	Maximum mode order (θ), ≤ 80 .
TITLE	-	Same as above (identify input pattern).
JIN	-	Number of input field points, ≤ 121 .
IC1	-	If ≤ 0 , convert incident field from dB to volts.
IC2	-	If > 0 , neglect incident field phase.
IC3	-	If > 0 , compute incident field amplitude from equations inserted in MAIN.
IC4	-	If > 0 , $E_\theta = E_\phi$ in phase and amplitude; input each phase and amplitude once.
PSI	-	Polar angle θ .
E	-	$E_\theta(\theta, \phi)$ amplitude as a function of θ of input pattern; volts or dB.
EP	-	$E_\theta(\theta, \phi)$ phase; degrees.
H	-	$E_\phi(\theta, \phi)$ amplitude; volts or dB.
HP	-	$E_\phi(\theta, \phi)$ phase.

Input variables (continued):

- JMAX0 - $180/\Delta\theta + 1$ where $\Delta\theta$ is the desired output increment of the far field SWE pattern.
- JOUT - Number of output values starting with $\theta = 0^\circ$.

Input considerations:

As mentioned previously, only one azimuthal expansion component can be handled by the SWE program. This component is usually $m = 1$, although it can be zero or > 1 . The $m = 1$ mode typically corresponds to a pattern amplitude which varies as a full cycle sinusoid in total azimuth. It is for this case that one refers to the "E and H planes" of a pattern. If more than one azimuth component is needed, the components can be run separately and the resulting coefficients superimposed.

The number of n modes necessary to accurately represent the input pattern depends on the complexity of that pattern. For the E plane (E_θ) pattern of the C-band Cassegrain feed at 5 cm, roughly $n = 70$ modes were needed. Since the SWE program requires only about one minute of CPU time, the user can easily test various maximum mode orders and check which yields the most accurate far field pattern.

Comparing the far field spherical wave expansion pattern to the input pattern, one can see that the former tends to oscillate about the latter. Frequent sampling of the input pattern will help to minimize this. For the 5 cm case above, feed pattern values were input every 0.25° until the input pattern was about 35 dB down ($\theta = 20^\circ$) and every 0.5° out to -45 dB ($\theta \approx 30^\circ$).

An oscillation of several hundredths of a volt in the approximation pattern values is of little concern since the input pattern is usually not reliable to such resolutions. A more important consideration is that the patterns contain roughly the same power through the average angle out to the edge of the reflecting surface.

Table of inputs and their formats:

<u>Card</u>		<u>Format</u>
1	TITLE	18A4
2	MCOMP, LMAX	2I5
3	TITLE	18A4
4	J1N, IC1, IC2, IC3, IC4	5I5
5	⋮ IC4 ≤ 0 : T(1), E(1), EP(1), H(1), HP(1)	5(1X, F9.5)
	⋮ IC4 > 0 : T(1), E(1), EP(1)	3(1X, F9.5)
⋮	⋮	⋮
JIN+4	⋮ IC4 ≤ 0 : T(JIN), E(JIN), EP(JIN), H(JIN), HP(JIN)	5(1X, F9.5)
	⋮ IC4 > 0 : T(JIN), E(JIN), EP(JIN)	3(1X, F9.5)
JIN+5	JMAX0, JOUT	2I5

2. MULT (PMN MULT)

Subroutine description:

MULT multiplies two matrices. The matrix dimensions are specified in the CALL MULT statement. If any variable dimensions are changed in MAIN, the CALL MULT statement must be changed accordingly.

3. LEGEND (PMN LEGEND)

4. VECTOR (PMN VECTOR)

LEGEND and VECTOR described previously.

D. Output of the SWE Program.

- E and H planes of input pattern in volts and degrees.
- Real and imaginary values of TE and TM wave coefficients for each mode.
- Fraction of total mode power of the coefficients for each mode.

- Total mode power of coefficients.
- Far field summation of spherical modes;
E and H planes in volts and degrees.

E. Submitting a Job.

It will be assumed that the user is already familiar with the Pandora command system of the IBM 360. If the user is not familiar with the system, the Pandora Guide (assembled by the Charlottesville Computer Division) nicely explains it. Some essential member-oriented commands are the following: ENTER, FETCH, SAVES, CLEAR, SCRATCH, CWS, CONCAT, SEQUENCE and SUBMIT. Some essential line editing commands are the following: CHANGE, INSERT, DELETE, MOVE, COPY, EDIT and SEEK.

The necessary JCL (Job Control Language) parameters for the programs are contained in the Pandora members JCLSCAT and JCLSWE (some message suppression). Only the TIME parameter is set by the user. This specifies a maximum CPU time in minutes. In JCLSWE, TIME = 1 and need not be adjusted.

To submit the SCAT program, the following statement is entered:

```
SUBMIT_JCLSCAT_MAINISC_SSURF_FINT_SFPVSALD_DATANAME
```

(PMN SFPVSALD contains eight subroutines.)

In this statement, JCLSCAT must be first and SFPVSALD and DATANAME must be last and in the shown order. The three other members can be arbitrarily shuffled. DATANAME is specified by the user.

To submit the SWE program, the following statement is entered:

```
SUBMIT_JCLSWE_MAINOR_LVM_DATANAME
```

(PMN LVM contains three subroutines.)

This statement must be entered with the order as shown.

It must be kept in mind that whenever the incident field is to be changed in the SCAT program, the SWE program must be run so that the proper coefficients are stored on the disk.

The following Pandora members may prove useful:

- DSCATSUB - Contains input data for scattering the 140-ft subreflector; $\Delta\theta = 0.2^\circ$, $\Delta\phi = 2^\circ$.
- DSWE5G - Input data for SWE program. For 5 GHz C-band Cassegrain feed pattern.
- DSWE10G - For 10 GHz X-band Cassegrain feed.

Shown in Figure 5 is the scattered pattern from the 140-ft subreflector at 5 GHz.

IV. SOME OF THE THEORY UNDERLYING THE PROGRAMS

A. Determining the Scattered Fields.

1. Problem definition:

Given a perfectly conducting reflecting surface and the magnetic fields incident on this surface, the far field scattered pattern is to be found. The incident electric field does not contribute to the fields scattered from a perfect conductor.

In the coordinate system of Figure 1, the following variables are defined:

- $\bar{E}_s(R, \theta, \phi)$ = far zone scattered electric field = $[E_\theta(\theta, \phi) \hat{\theta} + E_\phi(\theta, \phi) \hat{\phi}] \frac{e^{-jkR}}{R}$, $R \rightarrow \infty$.
- $\hat{R}, \hat{\theta}, \hat{\phi}$ = unit vectors.
- $\bar{H}_i(\rho, \theta, \phi)$ = incident magnetic field.
- \hat{n} = outward unit normal to scattering surface.
- $\hat{\rho}, \hat{\theta}, \hat{\phi}$ = unit vectors.
- $\bar{K}(\rho, \theta, \phi)$ = induced surface current.
- $\bar{\rho}$ = $\rho \hat{\rho}$.
- $\hat{i}, \hat{j}, \hat{k}$ = Cartesian unit vectors.
- k = $2\pi/\lambda$.
- j = $\sqrt{-1}$.

140-FT SUBREFLECTOR SCATTERED PATTERN 8/21/81

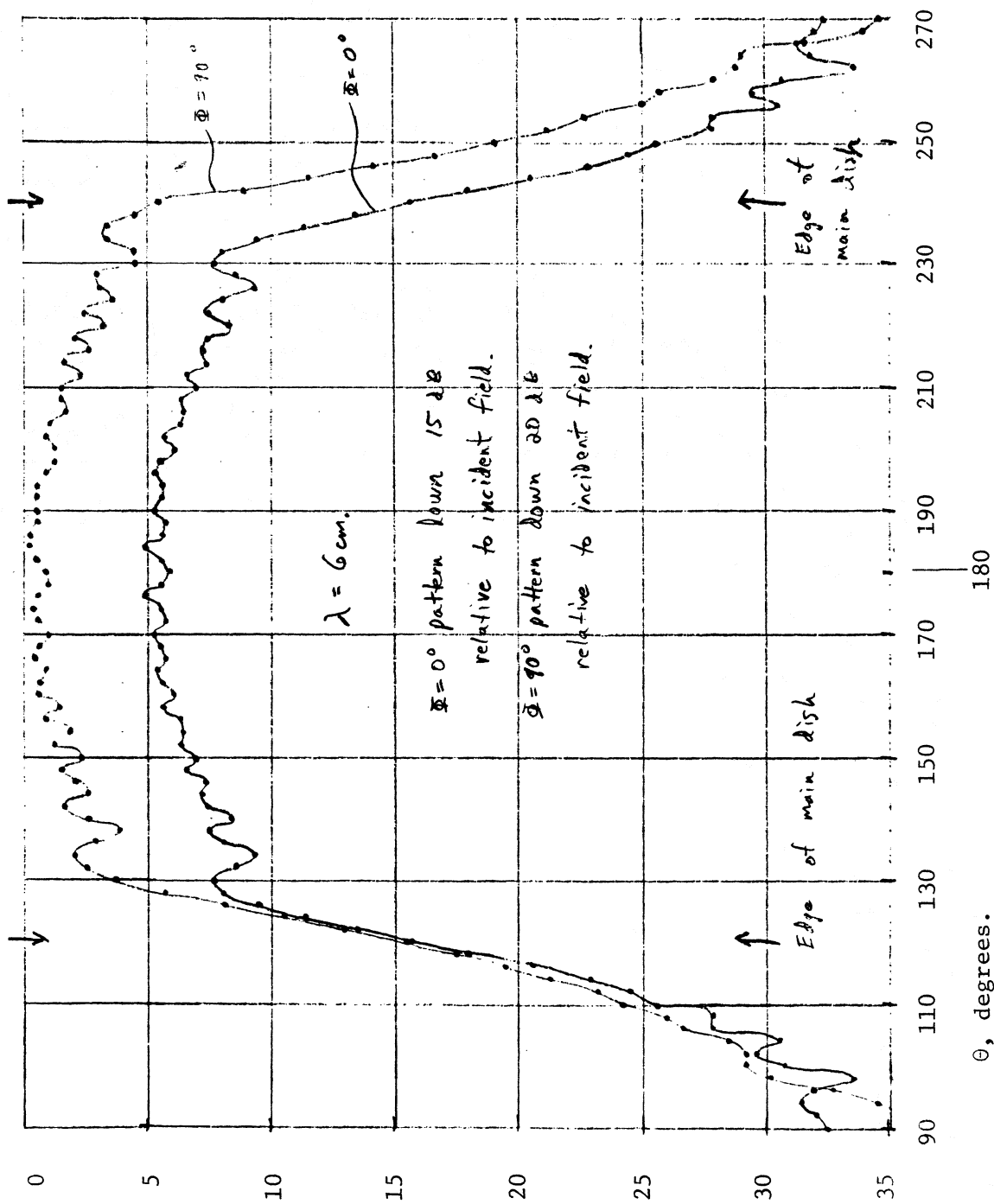


FIGURE 5

The program computes E_θ and E_ϕ in volts. If near field values are needed, then E_θ and E_ϕ can be input to the spherical wave expansion (SWE) program.

To find the total scattered field, the far field incident pattern is added to E_s .

2. Induced currents and the physical optics approximations:

The scattered field, E_s , is determined by taking the 2-dimensional Fourier transform (mapping from $\theta, \phi \rightarrow \Theta, \Phi$) of the array of current dipoles. Diffraction effects will be accounted for since the wavefronts of each current dipole are summed. In principle, the accuracy of the technique is limited by the precision with which the induced currents are known.

To determine \bar{K} , the following assumptions are made: On directly illuminated sections of the reflector, the induced currents are those of an optically reflected \bar{H}_i ; on shadowed portions, there are no induced currents. These two assumptions comprise the physical optics approximations.

An expression for \bar{K} can be found by applying the field boundary conditions to the scattering surface. For two arbitrary mediums,

$$\bar{K} = \hat{n} \times (\bar{H}_1 - \bar{H}_2).$$

Since there are no fields beneath the reflector surface, $\bar{H}_2 = 0$ and since \bar{H}_i is optically reflected, $\bar{H}_1 = 2\bar{H}_i$. Thus

$$\bar{K} = 2\hat{n} \times \bar{H}_i \tag{1}$$

Note that \bar{K} gives rise to all fields (incident and scattered) in free space (medium 1).

Ludwig notes that this expression is, in some cases, a poor approximation to the true induced currents. Generally, the calculated currents oscillate about the true currents, making little net contribution to the scattered

fields. Equation (1) can be accurately used for reflectors as small as several wavelengths in diameter.

3. Fields due to the induced currents.

To determine the scattered field at some far field output point (θ, ϕ) , the amplitude and phase delay (pathlength) of the radiation from each current dipole is vectorially summed. In principle, the current dipoles are of infinitesimal length; for the numerical integration this, of course, is not true.

For each dS , $d\bar{E}_s$ is proportional to the currents induced on dS . Thus, the radiation amplitude is \bar{K} .

Since the far field of \bar{E}_s is desired, it is instructive to view the output grid, over which the scattered pattern is calculated, as lying on a sphere of infinite radius centered on the source phase center. If the source fields travel directly to the output grid without striking a reflector, the pathlength covered is defined to be zero. When intercepted first by a reflector, the pathlength is defined as the additional distance travelled due to the reflection. This additional distance is that covered by the source fields to the reflector minus the extent to which the incident fields have already travelled in the direction of the output point (extent of colinearity). If R and vector $\hat{\rho}$ are colinear, the pathlength is, by definition, zero. A measure of the colinearity of two vectors is the cosine of the angle between them (dot product). Thus, the pathlength is

$$\gamma = \rho(1 - \hat{\rho} \cdot \hat{R}).$$

To evaluate $\hat{\rho} \cdot \hat{R}$, define the following:

$$\begin{array}{ll} x = \rho \sin \theta \cos \phi & X = R \sin \theta \cos \phi \\ y = \rho \sin \theta \sin \phi & Y = R \sin \theta \sin \phi \\ z = \rho \cos \theta & Z = R \cos \theta \end{array}$$

and

$$\hat{\rho} = (x\hat{i} + y\hat{j} + z\hat{k})/\rho, \quad \hat{R} = (X\hat{i} + Y\hat{j} + Z\hat{k})/R.$$

The dot product is then

$$\hat{\rho} \cdot \hat{R} = [\sin \theta \sin \Theta (\cos \phi \cos \Phi + \sin \phi \sin \Phi) + \cos \theta \cos \Theta].$$

The pathlength must be normalized to the wavelength. For an outward travelling spherical wave, the phase delay from source phase center to output point is

$$\text{Phase delay} = e^{-jk\gamma}.$$

At this point, our knowledge of the scattered fields can be summarized by

$$\bar{E}_s \propto \frac{e^{-jkR}}{R} \int_s \bar{K} e^{-jk\gamma} ds,$$

where \bar{E}_s and \bar{K} are expressed in volts/meter and amperes/meter, respectively. The variable pathlength is neglected in computing the spacial attenuation ($1/R$) of \bar{E}_s , since R goes to infinity. Therefore,

$$E_{\hat{\theta}} + E_{\hat{\phi}} \propto \int_s \bar{K} e^{-jk\gamma} ds. \quad (2)$$

To get the two sides of equation (2) compatible, we must convert the current-distance of the right side to voltage. Recalling the phase quadrature of associated currents and voltages, the constant of proportionality is $-2\pi j Z_0/\lambda$, where $Z_0 =$ the impedance of free space. The usual expression for this constant is $-j\omega\mu_0$, where μ_0 is the magnetic permeability of free space and $\omega = 2\pi$ times the frequency. Define

$$\bar{I}(\theta, \phi) = [E_{\hat{\theta}} + E_{\hat{\phi}}]/-j\omega\mu_0 \quad (\text{ampere-meters}).$$

Then

$$\bar{I} = \frac{1}{4\pi} \int_s \bar{K} e^{-jk\gamma} ds.$$

The $\frac{1}{4\pi}$ is a normalization constant which arises from integrating over a solid angle. A sphere has 4π square radians (steradians).

Since \bar{H}_i will often be evaluated within the source's near field, it will have a relatively strong radial (ρ) dependence. However, this dependence is predominantly of the form $1/\rho$. (The $e^{-jk\rho}$ phase term has already been included in the expression for γ .) Define:

$$\bar{H}_i = \frac{1}{\rho} [H_\rho \hat{\rho} + H_\theta \hat{\theta} + H_\phi \hat{\phi}],$$

where H_ρ , H_θ and H_ϕ are complex values, thus allowing for phase deviating from the $e^{-jk\rho}$ form. Since H_ρ , H_θ and H_ϕ are slowly varying with ρ , the numerical integration is easier.

Finally, some results from differential geometry concerning normals to surfaces must be considered. We have defined $\bar{\rho}$ as $\bar{\rho} = \bar{\rho}(\theta, \phi)$; specifying θ and ϕ describes the surface. Figure 6 shows the vectors $\bar{\rho}$, $\frac{\partial \bar{\rho}}{\partial \phi}$ and $\frac{\partial \bar{\rho}}{\partial \theta}$. The latter two are tangent to the surface at (θ, ϕ) and are normal to each other. As such, their cross product defines a normal to the surface at (θ, ϕ) . After some staring at Figure 6, one can see that $\frac{\partial \bar{\rho}}{\partial \theta}$ and $\frac{\partial \bar{\rho}}{\partial \phi}$ are the resultants of vectors in the directions $\hat{\rho}$ and $\hat{\theta}$ and the directions $\hat{\rho}$ and $\hat{\phi}$, respectively.

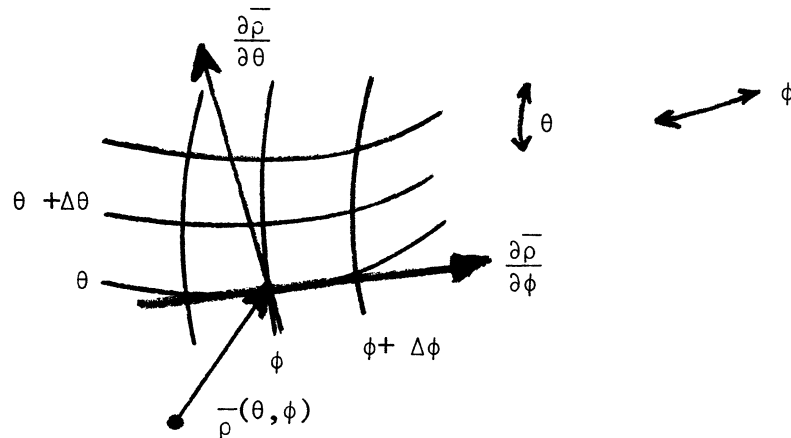


FIGURE 6

Evaluating these derivatives (Figure 1 is helpful), we find:

$$\frac{\partial \bar{\rho}}{\partial \theta} = \left(\frac{\partial \rho}{\partial \theta} \right) \hat{\rho} + \rho \left(\frac{\partial \hat{\rho}}{\partial \theta} \right) = \frac{\partial \rho}{\partial \theta} \hat{\rho} + \rho \hat{\theta}$$

$$\frac{\partial \bar{\rho}}{\partial \phi} = \left(\frac{\partial \rho}{\partial \phi} \right) \hat{\rho} + \rho \left(\frac{\partial \hat{\rho}}{\partial \phi} \right) = \frac{\partial \rho}{\partial \phi} \hat{\rho} + \rho \sin \theta \hat{\phi}$$

and

$$\frac{\partial \bar{\rho}}{\partial \phi} \times \frac{\partial \bar{\rho}}{\partial \theta} = \rho \frac{\partial \rho}{\partial \phi} \hat{\phi} + \rho \sin \theta \frac{\partial \rho}{\partial \theta} \hat{\theta} - \rho^2 \sin \theta \hat{\rho}.$$

The unit normal, \hat{n} , can be found from the relationship

$$\hat{n} dS = \frac{\partial \bar{\rho}}{\partial \phi} \times \frac{\partial \bar{\rho}}{\partial \theta} d\theta d\phi$$

where $dS = (\rho \sin \theta d\phi) (\rho d\theta) = \rho^2 \sin \theta d\phi d\theta$. Substituting for $\hat{n} dS \times \bar{H}_i$ in the equation for \bar{I} , we obtain

$$\bar{I} = \frac{1}{2\pi} \int_s \bar{F}(\theta, \phi) e^{-jkY} d\theta d\phi$$

where

$$\begin{aligned} \bar{F} = & \left(\frac{\partial \rho}{\partial \theta} \sin \theta H_\phi - \frac{\partial \rho}{\partial \phi} H_\theta \right) \hat{\rho} + \left(\frac{\partial \rho}{\partial \phi} H_\rho + \rho \sin \theta H_\phi \right) \hat{\theta} \\ & + \left(-\rho \sin \theta H_\theta - \sin \theta \frac{\partial \rho}{\partial \theta} H_\rho \right) \hat{\phi} . \end{aligned}$$

This is the form in which the radiation integral is evaluated.

B. Technique of Numerical Integration.

1. Form of integral:

Dropping all constants, the form of the integral over a given integration grid and for the output point θ_J, ϕ_K is

$$I(\theta_J, \phi_K) = \int_{\theta_1}^{\theta_M} \int_{\phi_1}^{\phi_N} F(\theta, \phi) \exp[jk\gamma(\theta, \phi, \theta_J, \phi_K)] d\theta d\phi$$

Define $\Delta\Omega_{mn}$ as the piece of solid angle bounded at its four corners by $\theta_m, \theta_{m+1}, \phi_m,$ and ϕ_{m+1} . Also, define

$$\begin{aligned}\Delta\theta_m &= \theta_{m+1} - \theta_m \\ \Delta\phi_n &= \phi_{n+1} - \phi_n\end{aligned}$$

$$\rho_{mn}, F_{mn}, \text{ etc.} = \rho(\theta_m, \phi_n), F(\theta_m, \phi_n), \text{ etc.}$$

The strong sinusoidal variation of the phase delay creates the need for rapid sampling of the scattering surface. Consider a $\Delta\Omega_{mn}$ with physical dimensions on the order of a wavelength; $\rho_{mn} \Delta\theta_m \sim \lambda$ and $\rho_{mn} \sin \theta_m \Delta\phi_n \sim \lambda$. (The dimension of one wavelength is chosen because electromagnetic fields rarely change abruptly over such distances.) For this situation the phase delay can vary up to one full cycle as $\Delta\Omega_{mn}$ is traversed. To successfully apply an integration technique such as Simpson's rule would require further subdivision of the scattering surface, thus creating a monstrous CPU time. Note that because γ is a nonlinear function of θ and ϕ , the fast Fourier transform cannot be applied.

2. The integration technique:

The sampling hassle can be alleviated through use of a linear representation of F and γ (especially) over each $\Delta\Omega_{mn}$. (The procedure to be described below was developed by Ludwig and is shown on page 13 of his JPL report.) Approximate F and γ by

$$\begin{aligned}F(\theta, \phi) &= a_{mn} + b_{mn} (\theta - \theta_m) + c_{mn} (\phi - \phi_n) \\ \gamma(\theta, \phi) &= \alpha_{mn} + \beta_{mn} (\theta - \theta_m) + \xi_{mn} (\phi - \phi_n)\end{aligned}$$

where θ and ϕ assume the corner values of $\Delta\Omega_{mn}$. Applying a least squares plane fit to F and γ at the corners of $\Delta\Omega_{mn}$, the following normal equations are obtained for F :

$$\Sigma F = a_{mn} \Sigma 1 + b_{mn} \Sigma (\theta - \theta_m) + c_{mn} \Sigma (\phi - \phi_n)$$

$$\Sigma F(\theta - \theta_m) = a_{mn} \Sigma (\theta - \theta_m) + b_{mn} \Sigma (\theta - \theta_m)^2 + c_{mn} \Sigma (\phi - \phi_n) (\theta - \theta_m)$$

$$\Sigma F(\phi - \phi_n) = a_{mn} \Sigma (\phi - \phi_n) + b_{mn} \Sigma (\theta - \theta_m) (\phi - \phi_n) + c_{mn} \Sigma (\phi - \phi_n)^2$$

where

$$\Sigma = \sum_n^{n+1} \sum_m^{m+1}$$

Solving for the coefficients yields the following:

$$a_{mn} = \frac{1}{4} [3 F_{mn} - F_{m+1 \ n+1} + F_{m+1 \ n} + F_{m \ n+1}]$$

$$b_{mn} = \frac{1}{2\Delta\theta_m} [F_{m+1 \ n} - F_{mn} + F_{m+1 \ n+1} - F_{m \ n+1}]$$

$$c_{mn} = \frac{1}{2\Delta\phi_n} [F_{m \ n+1} - F_{mn} + F_{m+1 \ n+1} - F_{m+1 \ n}]$$

The results are the same for γ and its coefficients.

Substituting the approximations for F and γ into equation (2) and integrating over $\Delta\Omega_{mn}$, the scattered field contribution ΔI_{mn} can analytically be determined. The expression for ΔI_{mn} is

$$\Delta I_{mn} = \exp jk\alpha_{mn} \left\{ a_{mn} \begin{bmatrix} \frac{e_m - 1}{jk\beta_{mn}} \\ \frac{e_n - 1}{jk\xi_{mn}} \end{bmatrix} \right. \\
+ b_{mn} \left[\frac{\Delta\theta_m}{jk\beta_{mn}} e_m - \left(\frac{e_m - 1}{(jk\beta_{mn})^2} \right) \right] \begin{bmatrix} \frac{e_n - 1}{jk\xi_{mn}} \end{bmatrix} \\
\left. + c_{mn} \begin{bmatrix} \frac{e_m - 1}{jk\beta_{mn}} \end{bmatrix} \left[\frac{\Delta\phi_n}{jk\xi_{mn}} e_n - \left(\frac{e_n - 1}{(jk\xi_{mn})^2} \right) \right] \right\}$$

where

$$e_m = \exp jk\beta_{mn} \Delta\theta_m$$

$$e_n = \exp jk\xi_{mn} \Delta\theta_n$$

For each θ_J, ϕ_K of the output grid, the integration subroutine (FINT) computes I from all θ_m, ϕ_n on the integration grid. At every integration point the subroutine computes the least squares coefficients and evaluates the above expression for ΔI_{mn} . To avoid numerical catastrophe an altered expression of ΔI_{mn} is used for β_{mn} and/or ξ_{mn} near zero.

3. Some comments on the method:

Analytic evaluation of the radiation integral can be thought of as summing the (relatively simple) patterns due to infinitesimal current dipoles. The numerical technique sums the patterns, given by ΔI_{mn} , for surface elements a wavelength on a side.

Ludwig claims that for scattering from a hyperboloid, a $\Delta\Omega_{mn}$ 2/3 square wavelengths in size results in errors 40 dB below the scattered pattern maximum.

C. A Few Details of Spherical Wave Expansions.

1. The SWE representation of an electromagnetic field

Define the following variables:

$\bar{E}, \bar{H}(\rho, \theta, \phi)$	=	description everywhere in a source-free region V of an electromagnetic field.
TE_{mn}, TM_{mn}	=	transverse electric and transverse magnetic fields used to describe \bar{E} and \bar{H} .
m	=	order of azimuthal variation of TE, TM fields.
n	=	mode order of TE, TM fields.
\bar{m}, \bar{n}	=	spherical wave solutions to Maxwell's equations; define TE and TM fields.
a_{mn}, b_{mn}	=	TE and TM expansion coefficients.
$z_n(k\rho)$	=	any solution to spherical Bessel (differential) equation.
$h_n^{(2)}(k\rho)$	=	spherical Hankel function (a particular z_n).
$P_n^m(\cos \theta)$	=	associated Legendre function (solution to a form of the Legendre differential equation).

The SWE is used (in this program) to represent a far field input pattern, \bar{E} and \bar{H} , in both the near and far fields. An expansion in TE and TM spherical waves is used as shown:

$$\bar{E} \text{ or } \bar{H} = \sum_m \sum_n a_{mn} TE_{mn} + b_{mn} TM_{mn}.$$

For \bar{E} , TE_{mn} has no radial components and for \bar{H} , TM_{mn} has no radial components.

As discussed earlier, the SWE program calculates the complex-valued wave coefficients. Input for \bar{H} is the far-zone magnetic field of, for example, a feed; $\bar{E} = 0$. The program reduces the flexibility of the expansion in two ways: one, by requiring that m assume a single integer value (usually $m = 1$) and eliminating the summation over m; and, two, by discarding the odd solutions of \bar{m}_{mn} and \bar{n}_{mn} , thus requiring that \bar{H} be linearly polarized. (The existence of even and odd solutions has been, for convenience, left unrecognized in the notation.)

The general spherical wave expansion is as follows:

Let a sphere of radius ρ_0 contain all field sources. Then the electromagnetic field in the (source-free) region V that includes all space outside of the sphere of radius ρ_0 is

$$\bar{E}(\rho, \theta, \phi) = -\sum_m \sum_n a_{mn} \bar{m}_{mn} + b_{mn} \bar{n}_{mn}$$

$$\bar{H}(\rho, \theta, \phi) = \frac{k}{j\omega\mu} \sum_m \sum_n a_{mn} \bar{n}_{mn} + b_{mn} \bar{m}_{mn}$$

where $\omega\mu/k$ is Z_0 = free space impedance and where

$$\begin{aligned} \bar{m}_{mn} &= z_n(k\rho) \frac{P_n^m(\cos \theta)}{\sin \theta} \sin m\phi \hat{\theta} \\ &\quad - z_n(k\rho) \frac{\partial}{\partial \theta} P_n^m(\cos \theta) \cos m\phi \hat{\phi} \\ \bar{n}_{mn} &= n(n+1) \frac{z_n(k\rho)}{k\rho} P_n^m(\cos \theta) \sin m\phi \hat{\rho} \\ &\quad + \frac{1}{k\rho} \frac{\partial}{\partial \rho} [\rho z_n(k\rho)] \frac{\partial}{\partial \theta} P_n^m(\cos \theta) \cos m\phi \hat{\theta} \\ &\quad + \frac{1}{k\rho} \frac{\partial}{\partial \rho} [\rho z_n(k\rho)] \frac{P_n^m(\cos \theta)}{\sin \theta} \cos m\phi \hat{\phi} \end{aligned}$$

(The even solutions contain the upper of the two signs and the upper of the sin-cos $m\phi$ pair. Further discussion of the specifics of the SWE program expansion will be delayed until section C.)

Since we are dealing with travelling spherical waves, $z_n(k\rho)$ should describe such. The spherical Bessel function that describes the radial variation of an outward travelling spherical wave is $h_n^{(2)}(k\rho)$, the Hankel function of the second kind.

Note the roles played by the mode indices, m and n ; two indices describe variation in three (coordinate) variables. Theta pattern variation, as described by $P_n^m(\cos \theta)$ is dependent on both mode orders (m is not an exponent).

2. Spherical waves as solutions to Maxwell's equations:

In a source free medium, the curl relationships of Maxwell's equations completely describe an electromagnetic field, \bar{E} and \bar{H} . That is,

$$\bar{\nabla} \times \bar{E} = -\mu_0 \frac{\partial \bar{H}}{\partial t} = -j\omega\mu_0 \bar{H}$$

$$\bar{\nabla} \times \bar{H} = \epsilon_0 \frac{\partial \bar{E}}{\partial t} = j\omega\epsilon_0 \bar{E}$$

where

$$\bar{\nabla} = \frac{\partial}{\partial \rho} \hat{\rho} + \frac{1}{\rho} \frac{\partial}{\partial \theta} \hat{\theta} + \frac{1}{\rho \sin \theta} \frac{\partial}{\partial \phi} \hat{\phi}, \quad \text{and}$$

where an $e^{j\omega t}$ time variation has been assumed and where $\omega\mu_0 = kZ_0$ and $\omega\epsilon_0 = k/Z_0$. These equations state that the change with distance in directions perpendicular to \bar{E} (\bar{H}) is proportional to the time rate of change of \bar{H} (\bar{E}). The constants of proportionality, μ_0 and ϵ_0 , specify the extent per unit distance to which free space can transmit energy in magnetic and electric fields, respectively. (μ_0 and ϵ_0 are, as shown above, directly related to the impedance of free space, Z_0 .) The negative sign in the first curl equation is a result of the orientation of \bar{E} and \bar{H} in a right-hand coordinate system.

Another interpretation of the curl equations is that the existence of an electric field ensures the existence of a corresponding magnetic field, and vice versa, for non-static fields.

Eliminating each variable of the curl equations yields the following:

$$\bar{\nabla} \times (\bar{\nabla} \times \mathbf{H}) = k^2 \bar{\mathbf{H}} \quad (1)$$

$$\bar{\nabla} \times (\bar{\nabla} \times \bar{\mathbf{E}}) = k^2 \bar{\mathbf{E}}$$

Using vector identities in rectangular coordinates

$$(\bar{\nabla} = \frac{\partial}{\partial x} \hat{i} + \frac{\partial}{\partial y} \hat{j} + \frac{\partial}{\partial z} \hat{k})$$

and noting that the divergence of $\bar{\mathbf{E}}$ and $\bar{\mathbf{H}}$ are zero ($\bar{\nabla} \cdot \bar{\mathbf{E}}, \bar{\nabla} \cdot \bar{\mathbf{H}} = 0$ in region V since no sources exist there), these equations are transformed to

$$\bar{\nabla}^2 \bar{\mathbf{H}} + k^2 \bar{\mathbf{H}} = 0 \quad (\text{Same for } \bar{\mathbf{E}}.) \quad (2)$$

This is the vector wave or vector Helmholtz equation.

$\bar{\nabla}^2$ is often called the Laplacian. The (negative of the) Laplacian of a function is closely related to the difference between the value of the function at some point and the average value of the function at straddling points. That is, $\bar{\nabla}^2$ determines the concavity or lumpiness of a function. Thus, equation (2) states that the concavity of the magnitude or direction of $\bar{\mathbf{H}}$ is proportional to $\bar{\mathbf{H}}$.

Since the curl of a vector function yields the magnitude and direction of the rotation, or vorticity, of the vector field, equation (1) can be understood as requiring that the vorticity lines of $\bar{\mathbf{E}}$ and $\bar{\mathbf{H}}$ must themselves exhibit vorticity. Further, this second order vorticity is proportional in magnitude and direction to $\bar{\mathbf{E}}$ and $\bar{\mathbf{H}}$. (See Methods of Theoretical Physics, Morse and Feshbach; chapter 1.) Note that in a source-free space, equations (1) and (2) are identical.

The vector solutions \bar{m}_{mn} and \bar{n}_{mn} of equation (2) are most easily found by considering some physical attributes of \bar{E} and \bar{H} . Recall that the divergence of \bar{E} and \bar{H} is zero, allowing them to be represented by the curl of a vector potential \bar{A} . (To this can be added the gradient of a scalar function. However, since the curl of the gradient of a function is identically zero, this term is of no use.) If $\bar{H} = \frac{1}{\mu_0} \text{curl } \bar{A}$, then from Maxwell's curl equations,

$$\bar{E} = -\frac{\mu_0}{Z_0} \frac{\partial \bar{A}}{\partial t} = -jk\bar{A}.$$

Our choice of \bar{A} can greatly reduce the complexity of the problem. To ease the pain of applying boundary conditions in spherical coordinates, the direction of \bar{A} is (if possible) chosen to be normal to a boundary surface. This will create fields tangent to the surface boundary. The magnitude of \bar{A} is $\psi(\rho, \theta, \phi)$, a rectangular component of \bar{E} or \bar{H} , times a coordinate scale factor. For spherical coordinates the direction and magnitude of \bar{A} are $\hat{\rho}$ and $\rho\psi$, respectively. Thus, a solution to equation (2) is

$$\bar{H} = \bar{\nabla} \times \bar{\rho}\psi = \bar{m} \quad (3)$$

where $\bar{\nabla}$ is in spherical coordinates. Then, equation (3) has no radial components and, if ψ is a component of \bar{H} , \bar{m} is a TM field.

We must, of course, have radial electromagnetic field components in a general field representation. Another solution to (2) which will supply the needed components is simply the curl of \bar{m} . Including a $\frac{1}{k}$ constant of proportionality, a second solution to (2) is

$$\bar{H} = \frac{1}{k} \bar{\nabla} \times \bar{m} = \bar{n}.$$

In general, \hat{n} has ρ, θ and ϕ components. (For detailed mathematical accounts of the solution to (2), see Stratton, chapter 7, and Morse and Feshback, chapter 13.)

The problem has been reduced to determining ψ . Recall that ψ is a solution to

$$\bar{\nabla}^2 H_p, E_p + k^2 H_p, E_p = 0 \quad , \quad p = x, y, z.$$

where E_p and H_p are functions of ρ, θ and ϕ . $\bar{\nabla}^2$ can then be expressed in spherical coordinates. The resulting differential equation is evaluated by the separation of variables

$$\psi(\rho, \theta, \phi) = \psi_1(\rho) \psi_2(\theta) \psi_3(\phi).$$

The solutions to the three resulting ordinary differential equations are:

$$\psi_1(\rho) = z_n(k\rho) = h_n^{(2)}(k\rho)$$

$$\psi_2(\theta) = P_n^m(\cos \theta)$$

$$\psi_3(\phi) = \frac{\cos}{\sin} m\phi \quad (\text{See Stratton, chapter 7.})$$

The previously shown form of \bar{m} and \bar{n} can be found from $\text{curl } \hat{\psi}_\rho$ and $\frac{1}{k} \text{curl curl } \bar{\psi}_\rho$.

3. What are spherical waves?

To understand, physically, what spherical waves are, it will be helpful to be aware of some mathematical characteristics of the spherical Hankel function and the associated Legendre function.

The spherical Hankel function of the second kind is defined as

$$h_n^{(2)}(\chi) = j_n(\chi) - j Y_n(\chi), \quad j = \sqrt{-1}$$

At $\chi = 0$, j_n is finite and Y_n has a pole. As $\chi_1 = k\rho_1 \rightarrow \infty$

$$h_n^{(2)}(\chi_1) = j^{n+1} \frac{e^{-j\chi_1}}{\chi_1}$$

and

$$\frac{1}{\chi_1} \frac{\partial}{\partial \chi} \left[\chi h_n^{(2)}(\chi) \right]_{\chi=\chi_1} = j^n \frac{e^{-j\chi_1}}{\chi_1}$$

These are closely related to the $e^{-jk\rho/\rho}$ dependence of a far field pattern.

The associated Legendre function provides the polar angle field variation. The order of index n runs from 0 to ∞ , and m runs 0 to n . A few modes of the function are shown below. (See Stratton, chapter 7; Tables of Function, Johnke and Emde, pp. 112-113.)

$$P_0(Z) = 1$$

$$P_1(Z) = Z = \cos \theta$$

$$P_1^1(Z) = (1 - Z^2)^{1/2} = \sin \theta$$

$$P_2(Z) = \frac{1}{2} (3Z^2 - 1) = \frac{1}{4} (3 \cos 2\theta + 1)$$

$$P_2^1(Z) = 3(1 - Z^2)^{1/2} Z = \frac{3}{2} \sin 2\theta$$

$$P_2^2(Z) = 3(1 - Z^2) = \frac{3}{2} (1 - \cos 2\theta)$$

The functions $\cos m\phi P_n^m(\cos \theta)$ and $\sin m\phi P_n^m(\cos \theta)$ are periodic on the surface of a unit sphere. For $m > 0$, the functions are zero at the poles. The number of nodal lines parallel to the equator is $n - m$. These are

orthogonally intersected by the $2m$ longitudinal nodes. Since the surface is divided into rectangular sections within which the above functions are alternately positive and negative, these functions are referred to as tesseral harmonics of n^{th} degree and m^{th} order.

Evaluating \bar{m} and \bar{n} at $m = 0$ and $n = 1$, the simplest form of \bar{m} and \bar{n} , will provide the vital clue in determining the nature of these spherical wave solutions.

Recall that \bar{E} and \bar{H} are defined everywhere in space V which surrounds but does not include a sphere of sources. Let this sphere be of radius a . An entirely equivalent representation of the sources is to place surface currents on the sphere and remove all sources within the sphere. From the boundary conditions, the surface currents are

$$\bar{K} = \hat{\rho} \times \bar{H}(a, \theta, \phi).$$

From the expressions for \bar{E} and \bar{H} involving the vector potential \bar{A} ,

$$\bar{E} = -jk\bar{A} \quad (\text{for } e^{j\omega t} \text{ time variation})$$

$$\bar{H} = \frac{1}{\mu_0} \text{curl } \bar{A}$$

where $\bar{A} = \bar{m}$ and \bar{n} , we can determine the induced currents. For $m = 0$, $n = 1$,

$$\bar{m}_{01} = h_1^{(2)}(k\rho) \sin \theta \hat{\phi}$$

$$\begin{aligned} \bar{n}_{01} &= \frac{1}{k} \bar{\nabla} \times \bar{m}_{01} \\ &= 2 \cos \theta \frac{1}{k\rho} h_1(k\rho) \hat{\rho} - \sin \theta \frac{1}{k\rho} \frac{\partial}{\partial \rho} [\rho h_1(k\rho)] \hat{\theta}. \end{aligned}$$

First consider the case of $\bar{A} = \bar{m}_{01}$. Through the use of various recurrence relations (see Stratton, chapter 7) and the relation $h_n(Z) = -j(-1)^n \left(\frac{d}{ZdZ} \right)^n \left(\frac{e^{iZ}}{Z} \right)$, the terms involving Hankel functions can be reduced to algebraic expressions in $k\rho$ and $e^{jk\rho}$. Then the surface current creating the $m = 0, n = 1$ mode fields is

$$\bar{K}_{01} \propto j \frac{\omega}{c} \sin \theta \left[\frac{1 - (ka)^2 - jka}{(ka)^3} \right] e^{jka} \hat{\phi}$$

(See Morse and Feshbach, chapter 13, p. 1867.)

Thus, we have a current oscillating in time parallel to the equator. The current goes to maximum at the equator and goes to zero at the poles and travels in the same direction at all points on the sphere. There is no charge build up as can be seen from the lack of radial \bar{E} components.

A current of such form characterizes the magnetic dipole and \bar{m}_{01} represents the field from the dipole.

If we are to determine the radial components of \bar{E} , we must turn the problem around and set $\bar{A} = \bar{n}_{01}$. Then, $\bar{E} = -jk \bar{n}_{01}$ and has a radial component. The surface current is

$$\bar{K}_{01} \propto \omega \mu_0 \sin \theta \left[\frac{j + ka}{(ka)^2} \right] e^{jka} \hat{\theta}$$

which oscillates between poles, alternately depositing positive and negative charges at the poles. As can be seen from $\hat{\rho} \cdot \bar{E}$, the charge is concentrated at the poles and is 90° out of phase with the current.

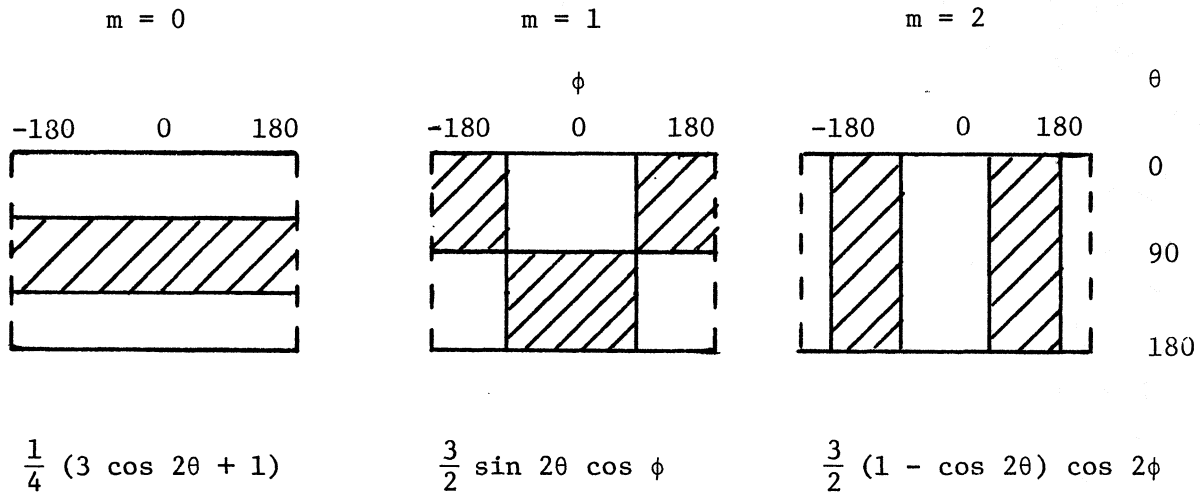
This oscillation of charge is merely a way of describing an electric dipole where \bar{n}_{01} represents the field from this dipole.

For $m = 0, n = 2$, the TE radiation is that of a magnetic quadrupole. The current circulates parallel to the equator going clockwise in one hemisphere

and counterclockwise in the other for half a cycle, then reversing for the next half cycle.

The corresponding TM radiation is that of an electric quadrupole. The current distribution describes charge moving alternately from the equator to the poles and then from the poles to the equator. Thus, free charge of one sign is accumulated at the poles and of the opposite sign at the equator, the signs changing at the next half cycle.

Among others, one important question that remains is the following: What is physically happening as m changes for a given n ? Plotting $\cos m\phi P_n^m(\cos\theta)$ on the unit sphere for various values of m seems to indicate that m determines two properties of the multipoles: the distribution of multipoles within the array and the orientation of the multipole array. As an example, consider the unit sphere plots for $n = 2$ and $m = 0, 1, 2$:



The shaded portions indicate negative value; nodal lines are solid. For $m = 0$, there is an overlap of 2 like (in sign) "monopoles". For $m = 1$ and 2, all "monopoles" (2 of each sign) are distinct but the quadrupole orientations are different.

4. Determining the SWE coefficients.

As mentioned previously, the summation over m in the spherical wave expansion is eliminated; usually an order of azimuthal variation of $m = 1$ is needed. The $m = 1$ corresponds to the situation in which a pattern may be divided into E and H planes. Since only the incident magnetic field contributes to the induced surface currents, only \bar{H}_i need be expanded. The expansion is:

$$\bar{H}_i(\rho, \theta, \phi) = \frac{k}{j\omega\mu} \sum_{n=1}^N a_n \bar{n}_n + b_n \bar{m}_n$$

where $\bar{n}_n \equiv \bar{n}_{mn}$ and is TE and $\bar{m}_n \equiv \bar{m}_{mn}$ and is TM. This expansion will determine \bar{H}_i , where \bar{H}_i is originally a far field pattern, anywhere in the space outside a sphere of radius ρ_0 which encloses the sources of \bar{H}_i .

Ludwig derives expressions for the expansion coefficients for the case in which the involved data are the tangential components of \bar{E} on a sphere of radius $\rho_1 > \rho_0$. The derivation will be outlined only; the mathematical details may be found in the JPL report.

The vector character of the expansion is eliminated (momentarily) by equating the tangential field components, $E_\theta(\theta, \phi)$ and $E_\phi(\theta, \phi)$, to the summation of the corresponding components of the expansion. The azimuthal terms of the summation are removed by taking the ordinary Fourier expansion of E_θ and E_ϕ , leaving a summation independent of ϕ . $E_\theta(\theta, 90^\circ)$ and $E_\phi(\theta, 0^\circ)$ are the usual far field patterns being expanded. The actual inputs to the program are $A_m(\theta)$ and $B_m(\theta)$, which represent the m^{th} Fourier component of the input pattern. (For $m = 1$, $A_m(\theta)$ and $B_m(\theta)$ are simply the pattern values.)

Using the integral (orthogonality) properties of Legendre functions, an integral (over θ) expression is found for the coefficients (JPL report; page 23,

equation 9). The far field value of the Hankel function and its derivative with respect to ρ is incorporated into the coefficients; that is, the coefficients actually computed are:

$$a_n' = a_n j^{n+1} e^{-jk\rho_1} / k\rho_1$$

$$b_n' = b_n j^n e^{-jk\rho_1} / k\rho_1$$

a_n' and b_n' are written on the computer disk where the SCAT program can gain access to them.

V. ACKNOWLEDGEMENTS

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(All but Ludwig's report are available in the Green Bank Library.)

VII.

PROGRAM CODE FOR BOTH PROGRAMS


```

C COMBINE SURFACE AND FIELD DATA TO DETERMINE VECTOR A
C A CONTAINS THREE COORDINATE COMPONENTS.
C
C 400 M=1,MMAX
C 400 N=1,NMAX
C T1=FT(M,N)*SIT(M)*HP(M,N)-EP(M,N)*HT(M,N)
C T2=FP(M,N)*HR(M,N)+F(M,N)*SIT(M)*HP(M,N)/PC
C T3=-FT(M,N)*SIT(M)*HR(M,N)-FM(N)*SIT(M)*HT(M,N)/PC
C A(M,N,1)=T1*SIT(M)*COP(N)+T2*COT(M)*COP(N)-T3*SIP(N)
C A(M,N,2)=T1*SIT(M)*SIP(N)+T2*COT(M)*SIP(N)+T3*COP(N)
C A(M,N,3)=T1*COT(M)-T2*SIT(M)
C 400 CONTINUE
C WRITE(6,2007)I
C
C BEGIN OUTPUT GRID LOOP FOR EACH POINT ON OUTPUT GRID DETERMINE
C SCATTERED FIELD CONTRIBUTION FROM ENTIRE INTEG GRID.
C
C WRITE(6,2006)I
C 500 K=1,KMAX
C TO=PP(K)/O,0.17453293
C WRITE(6,2011)O
C 500 J=1,JMAX
C
C ESTABLISH PATH LENGTH PARAMETER ON INTEGRATION GRID
C
C CALL PATHLEN(J,K,MMAX,NMAX,GAM)
C
C PERFORM INTEGRATION
C
C CALL FINI(I,P,A,GAM,MWAY-1,NMAX-1,STOT,I,J,K,PC,IFDGE,TFDGE)
C
C ASSIGN SCATTERED FIELDS AT OUTPUT POINT
C
C ETTO=COTT(J)*SIT(J)*COP(K)+STOT(2)*SIPP(K)-STOT(3)
C *SITT(J)
C EPPQ=STG(2)*COP(K)-STG(1)*SIPP(K)
C ETTO=-TO*O,1.01*PC/6.2831854*ETTO
C EPPQ=-TO*O,1.01*PC/6.2831854*EPPQ
C TO=TT(J)/O,0.17453293
C A1=REAL(ETTO)
C A2=AIMAG(ETTO)
C A3=REAL(EPPQ)
C A4=AIMAG(EPPQ)
C
C CONVERT TO POLAR COORDINATES
C
C CALL VECTOR(A1,A2,ETAMP,ETPHI)
C CALL VECTOR(A3,A4,EPAMP,EPPHI)
C WRITE(6,2012)O,ETAMP,ETPHI,EPAMP,EPPHI
C
C SUPERIMPOSE SCATTERED FIELD VALUE WITH VALUES CORRESPONDING
C TO OTHER INTEG GRIDS.
C
C ETIJ,K)=ETT(J,K)+ETTO
C EPIJ,K)=EPP(J,K)+EPPQ
C CONTINUE
C
C IF MORE INTEGRATION GRIDS REMAIN LOOP BACK
C
00910000 C
00911000 C
00912000 500
00920000 C
00930000 C
00940000 700
00950000 C
00960000 C
00990000 C
01000000 C
01010000 C
01040000 C
01050000 C
01090000 C
01090000 C
01091000 C
01092000 C
01120000 C
01130000 C
01140000 C
01150000 C
01170000 C
01180000 C
01190000 C
01191000 C
01200000 C
01270000 C
01280000 C
01281000 C
01290000 C
01320000 C
01330000 C
01340000 C
01350000 C
01360000 C
01370000 C
01380000 C
01390000 C
01400000 C
01410000 C
01420000 C
01430000 C
01440000 C
01450000 C
01451000 C
01460000 C
01470000 C
01480000 C
01481000 C
01482000 C
01483000 C
01484000 C
01490000 C
01500000 C
01530000 C
01540000 C
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01560000 C
01570000 C
01580000 C
01590000 C
01600000 C
01610000 C
01620000 C
01630000 C
01640000 C
01650000 C
01660000 C
01670000 C
01680000 C
01681000 C
01682000 C
01683000 C
01690000 C
01700000 C
01710000 C
01720000 C
01730000 C
01740000 C
01750000 C
01760000 C
01770000 C
01780000 C
01790000 C
01791000 C
01792000 C
01793000 C
01794000 C
01800000 C
01810000 C
01820000 C
01830000 C
01840000 C
01850000 C
01860000 C
01870000 C
01880000 C
01890000 C
01900000 C
01910000 C
01920000 C
01930000 C
01940000 C
01950000 C
01960000 C
01970000 C
01980000 C
01990000 C
02000000 C
02010000 C
02020000 C
02030000 C
02031000 C
02032000 C
02033000 C

```

```

IF(I=NGI)600,700,700
CALL RESET(I,MMAX,JMAX)
GO TO 100
I=I+1
ESTABLISH DIRECT RADIATION ON OUTPUT GRID
WRITE(6,2009)
DO 750 K=1,KMAX
TO=PP(K)/O,0.17453293
WRITE(6,2011)O
DO 750 J=1,JMAX
EVALUATE FAR-FIELD OF S.W.F
CALL FIELD2(I,J,K,ETTO,EPPQ)
TO=TT(J)/O,0.17453293
WRITE(10,2020)TO,ETTO,EPPQ
A1=REAL(ETTO)
A2=AIMAG(ETTO)
A3=REAL(EPPQ)
A4=AIMAG(EPPQ)
CALL VECTOR(A1,A2,ETAMP,ETPHI)
CALL VECTOR(A3,A4,EPAMP,EPPHI)
WRITE(6,2012)TO,ETAMP,ETPHI,EPAMP,EPPHI
ADD DIRECT AND SCATTERED FIELDS TO YIELD TOTAL(SCATTERED)
FIELDS.
ETT(J,K)=ETT(J,K)+ETTO
EPP(J,K)=EPP(J,K)+EPPQ
CONTINUE
TRANSLATE PHASE CENTER,SCALE FIELD AMPLITUDES,AND OUTPUT
TOTAL FIELDS
WRITE(6,2002)I,TITLE
WRITE(6,2010)XT,YT,ZT,SCALE
DO 760 K=1,KMAX
TO=PP(K)/O,0.17453293
WRITE(6,2011)O
DO 760 J=1,JMAX
TO=TT(J)/O,0.17453293
A1=REAL(ETT(J,K))
A2=AIMAG(ETT(J,K))
A3=REAL(EPP(J,K))
A4=AIMAG(EPP(J,K))
CALL VECTOR(A1,A2,ETAMP,ETPHI)
CALL VECTOR(A3,A4,EPAMP,EPPHI)
DP=XI*SITT(J)+COPP(K)+YT*SITT(J)*SIPP(K)+ZT*(COTT(J)+I,0)
DP=DP*CPC*57.29576
ETPHI=ETPHI-DP
EPPHI=EPPHI-DP
ADJUST PHASES TO -180,180 RANGE

```

```

02040000 C
02050000 C
02060000 C
02070000 C
02080000 C
02090000 C
02100000 C
02160000 C
02170000 C
02180000 C
02190000 C
02200000 C
02210000 C
02230000 C
02240000 C
02250000 C
02260000 C
02270000 C
02280000 C
02290000 C
02300000 C
02310000 C
02320000 C
02330000 C
02340000 C
02350000 C

CALL A2JUST(ETPHI)
CALL A2JUST(EPPHI)
ETAMP=ETAMP*SCALE
F2AMP=EPAMP*SCALE
WRITE(6,2012)TC,ETAMP,ETPHI,EPAMP,EPPHI
CONTINUE

260 C
STOP
FORMAT(18A4)
FORMAT(10F10.4)
2001 FORMAT(1H1,*,
2002 FORMAT(1H0,5X,18A4) NRAO SCATTERING PROGRAM,/,/5X,18A4)
2003 FORMAT(1O,*, PROPAGATION CONSTANT=,E14.8/)
2006 FORMAT(1H0,*, SCATTERED FIELDS FROM GRID*,I2)
2007 FORMAT(1H0,*, BEGIN INTEGRATION OVER GRID*,I2)
2009 FORMAT(1H1,*, DIRECT RADIATION FROM INCIDENT FIELDS*)
2010 FORMAT(1H0,*, SUPERPOSITION OF ALL GRID SCATTERED FIELDS AND DIRECT*,
E8 Z=,F10.4,*, Y=,F10.4,*,
E8 AMPLITUDE VALUES SCALED BY FACTOR OF*,E15.8)
2011 FORMAT(7H0 PHI=,F7.2,/,
E17X, THE THETA,15X,5HE PHI,/,
E9H THE TA,2I20H VOLTS PHASE)
2012 FORMAT(F9.2,F11.6,F9.2,F12.6,F8.2)
2013 FORMAT(F10.2,F10.6,F10.2,F10.6,F10.2)
END

02040000 C
02050000 C
02060000 C
02070000 C
02080000 C
02090000 C
02100000 C
02160000 C
02170000 C
02180000 C
02190000 C
02200000 C
02210000 C
02230000 C
02240000 C
02250000 C
02260000 C
02270000 C
02280000 C
02290000 C
02300000 C
02310000 C
02320000 C
02330000 C
02340000 C
02350000 C

SUBROUTINE SURFII,MMAX,NMAX,F,FT,FP,IEDGE,TEDGE)
C
C DIMENSION F(36,91),FT(36,91),FPI(36,91),TEDGE(91)
COMMON/GRIDI/SIT(36),COT(36),SIP( 91),COP( 91),T(36),P( 91)
C THIS SUB. PROVIDES MAIN WITH RHO,DRHO/DTHETA AND DRHO/DPHI
C INPUT RHOO,# INTEGRATION GRIDS,# THETA INTEG,GRIDS
C SOME VARIABLES:
C F,FT,FP = SURFACE PARAMETERS (RETURNED TO MAIN);F IS RHO AND
C FT AND FP ARE DRHO/DTHETA AND DRHO/DPHI.
C ISURF = INTEGER PARAMETER THAT SPECIFIES HOW SURFACE VALUES ARE
C OBTAINED.
READ(5,1001)ISURF
DT=(T(2)-T(1))/,01745329
DP=(P(2)-P(1))/,01745329
WRITE(6,2002)I,ISURF
FORMAT(1H0 *,FOR INTEGRATION GRID #,I3,*, ISURF =,I4,/)
1001 FORMAT(3I5)
WRITE(6,2004)
FORMAT(1H ,*, M N F FT FP*)
C ISURF CONSIDERS FOLLOWING CASES:
C ISURF<0,NOT=-2: F,FT,FP FOUND FROM EQUATIONS INSERTED BELOW.
C ISURF=-F FROM EQUATION;FT,FP NUMERICALLY DETERMINED.
C ISURF=0: F IS FUNCTION OF PHI ONLY AND IS IN TABULAR FORM.
C ISURF>0: FT,FP NUMERICALLY DETERMINED.
C F IS FUNCTION OF THETA AND PHI AND IS IN TABULAR
C FORM. FT,FP NUMERICALLY DETERMINED.
IF(1ISURF)20,40,30
CONTINUE
C CALCULATE F,FT,FP FROM EQUATION TO BE INSERTED HERE
DO 22 N=1,NMAX
DO 22 M=1,MMAX
DEN=COT(4)
F(M,N)=O.5 /DEN
IF(1ISURF .EQ. -2)GO TO 22
FT(M,N)=F(M,N)/COT(M)*SIT(M)
FP(M,N)=O.0
CONTINUE
22 IF(1ISURF .EQ. -2)GO TO 100
GO TO 110
CONTINUE
C CASE OF ZERO PHI VARIATION
40
C
C
C
C READ(5,1002)(=(M,1),M=1,MMAX)

```

```

MSTOP=0
DO 45 M=1,MMAX
  IF(MSTOP.EQ.1)GO TO 45
  IF (RHO=999.99999,EDGE HAS JUST BEEN PASSED;VALUES BEYOND EDGE
  WERE ASSIGNED ARBITRARILY.
  IF(IF(M,N).NE.999.99999)GO TO 49
  EXTRAPOLATE FOR RHO VALUES PAST REFLECTOR EDGE
  MSTOP=1
  IF(M.GE.4)F(M,N)=2.*F(M-1,N)-2.*F(M-2,N)+.5*F(M-3,N)
  IF(M.EQ.3)F(M,N)=2.*F(M-1,N)-F(M-2,N)
  IF(M.EQ.2)F(M,N)=F(M-1,N)
  CONTINUE
1002 FORMAT(1X,F9.5))
  CALCULATE FT USING FORWARD & BACKWARD DIFFERENCES
  CALL DIFF(I,-1, MMAX,1,F, FT,DT,DP)
  DO 43 M=1,MMAX
  DRHO/DPHI IDENTICALLY ZERO
  FP(M,1)=0.
  ASSIGNMENTS FOR N=2 TO NMAX
  DO 44 N=2,MMAX
  DO 44 M=1,MMAX
  F(M,N)=F(M,1)
  FT(M,N)=FT(M,1)
  FP(M,N)=FP(M,1)
  GO TO 110
  CONTINUE
  VARIATION IN THETA AND PHI
  CONTINUE
  READ(5,1002)((F(M,N),M=1,MMAX),N=1,NMAX)
  DO 31 N=1,NMAX
  MSTOP=C
  DO 32 M=1,MMAX
  IF(MSTOP.EQ.1)GO TO 35
  CHECK FOR RHO=999.99999
  IF(IF(M,N).NE.999.99999)GO TO 32
  EXTRAPOLATE RHO VALUES PAST REFLECTOR EDGE TO END OF INTEG GRID
  DO 31 N=1,NMAX
  MSTOP=1
  IF(M.GE.4)F(M,N)=2.*F(M-1,N)-2.*F(M-2,N)+.5*F(M-3,N)
  IF(M.EQ.3)F(M,N)=2.*F(M-1,N)-F(M-2,N)
  IF(M.EQ.2)F(M,N)=F(M-1,N)

```

```

32 CONTINUE
31 CONTINUE
  USE BACK & FORWARD DIFFERENCES TO COMPUTE FT & FP
  CALL DIFF(I,-1, MMAX,MMAX,F, FT,DT,DP)
  CALL DIFF(I,1, MMAX,MMAX,F, FP,DT,DP)
  CONTINUE
  WRITE SOME SURFACE PARAMETERS
  DO 120 N=1,MMAX,15
  DO 120 M=1,MMAX
  WRITE(6,2003) M,N,F(M,N),FT(M,N),FP(M,N)
  FORMAT(1X,2I3,3(1X,F10.4))
  RETURN
  END

```

```

00410000
00420000
-00430000
00431000
00432000
00433000
00434000
00440000
00450000
00460000
00470000
00480000
00490000
00500000
00510000
00520000
00530000
00540000
00550000
00560000
00570000
00580000
00590000
00591000
00592000
00593000
00600000
00610000
00620000
00630000
00640000
00650000
00660000
-00670000
00680000
-00690000
00700000
00710000
00720000
00730000
00740000
00750000
00760000
00770000
00780000
00790000
00800000
00801000
00802000
00803000
00810000
00811000
00812000
00813000
00820000
00830000
00840000
00850000

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CCCCCCCCCCCC
CCCCCCCCCCCC
C
C
SUBROUTINE EDGEED(NMAX,TEDEGE)
C
C   CALCULATE THETA AT EDGE OF REFLECTOR USING EQUATION IN PHI
C   CAN HAVE DIFFERENT EQUATIONS FOR SECTIONS IN PHI OF GRID
C
COMMON/GRID1/SIT(36),COT(36),SIP(91),COP(91),T(36),P(91)
DIMENSION TEDGE(91)

DTR=0.01745329
INSERT EQUATIONS HERE
NOTE: TEDGE MUST BE RETURNED IN DEGREES!!

DO 111 N=1,NMAX
IF((P(N)/DTR .GE. 315.) .AND. (P(N)/DTR .LE. 45.))GO TO 10
IF((P(N)/DTR .GE. 45.) .AND. (P(N)/DTR .LE. 135.))GO TO 20
IF((P(N)/DTR .GE. 135.) .AND. (P(N)/DTR .LE. 225.))GO TO 30
IF((P(N)/DTR .GE. 225.) .AND. (P(N)/DTR .LE. 315.))GO TO 40
10 TEDGE(N)=ATAN(0.069927/ABS(COP(N)))
GO TO 111
20 TEDGE(N)=ATAN(0.069927/ABS(SIP(N)))
GO TO 111
30 TEDGE(N)=ATAN(0.069927/ABS(COP(N)))
GO TO 111
40 TEDGE(N)=ATAN(0.069927/ABS(SIP(N)))
111 RETURN
END

SUBROUTINE FINI(X,Y,F,R,MMAX,NMAX,STOT,I,J,K,K,PC,TEDEGE,TEDEG)
THIS SUB. NUMERICALLY INTEGRATES THE RADIATION INTEGRAL
DIMENSION X(36),Y( 91), R(36,91),TEDEGE( 91), RTEMP(12)
COMPLEX F(36,91,31),STOT(31),T1,T2,T3,A,B,C,F12,F23,F14,SUM,TD
COMPLEX FVTEMP(2,3),AT,DELTA
SOME VARIABLE DEFINITIONS:
X,Y = THETA,PHI
R = PATHLENGTH TERM
F = VECTOR RELATED TO SURFACE CURRENTS
STOT = 3 COMPONENTS OF EVALUATED INTEGRAL (RETURNED TO MAIN)
TEDEGE = THETA VALUES THAT SPECIFY REFLECTOR EDGE AS A FUNCTION OF PHI
EDGE = INTEGER PARAMETER THAT SIGNIFIES PRESENCE OF EDGE AND SPECIFIES HOW TEDGE IS OBTAINED.
RTEMP,FVTEMP = TEMPORARY STORAGE FOR R,F
NCE = PARAMETER THAT INDICATES FOR EACH PHI,WHEN AN EDGE THETA HAS BEEN REACHED.
TEAVG = AVERAGE OF EDGE THETA VALUES CORRESPONDING TO PHI AND PHI+DELTA(PI). THIS COMPROMISE RESULTS IN A SLIGHTLY DISCONTINUOUS REFLECTOR EDGE.
NXP1=NMAX+1
MXP1=NMAX+1
INITIALIZE STOT
DO 10 L=1,3
STOT(L)=(0.0,0.0)
ENTER PHI LOOP
DO 200 N=1,NMAX
NCE=0
DY=0.5*(Y(N+1)-Y(N))
ENTER THETA LOOP (WITHIN PHI LOOP)
DO 201 M=1,MMAX
CHECK IF EDGE TO BE ENCOUNTERED ON THIS INTEG GRID
IF(TEDEGE .EQ. 11)GO TO 79
IF EDGE PRESENT,COMPARE EACH THETA (Y(M)) WITH THE EDGE VALUE OF THETA FOR THAT PHI. AT FIRST HINT OF EDGE,INTEGRATION PROCEDURE IS ALTERED.
X(1) MUST NOT BE > TEDGE(N) FOR ANY N
IF((TEDEGE(N) .GE. X(M+1)) .AND. (TEDEGE(N+1) .GE. X(M+1)))
GO TO 79
IF((X(M) .LE. TEDGE(N)) .AND. (TEDEGE(N) .LT. X(M+1)))
GO TO 30
IF((X(M) .LE. TEDGE(N+1)) .AND. (TEDEGE(N+1) .LT. X(M+1)))GO TO 30

```



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00350000    00320000
00360000    00624000
00370000    00626000
00380000    00630000
00390000    00631000
00395000    00631500
00400000    00632000
00405000    00633000
00410000    00640000
00415000    00650000
00420000    00660000
00430000    00670000
00435000    00671000
00440000    00672000
00450000    00673000
00460000    00680000
00470000    00690000
00475000    00700000
00480000    00710000
00485000    00720000
00500000    00730000
00505000    00731000
00510000    00732000
00520000    00733000
00530000    00734000
00540000    00740000
00550000    00750000
00560000    00760000
00570000    00770000
00575000    00780000
00576000    00790000
00576100    00800000
00576200    00810000
00576300    00820000
00580000    00830000
00590000    00840000
00595000    00850000
00600000    00860000
00601000    00870000
00601100    00880000
00601200    00890000
00601300    00900000
00601500    00910000
00602000    00920000
00603000    00930000
00604000    00940000
00606000    00950000
00610000    00960000
00611000    00970000
00612000    00980000
00613000    00990000
00615000    01000000
00615000    01000000
00615100    01010000
00615200    01020000
00615300    01030000
00616000    01031000
00620000    01032000
00621000    01033000

WRITE(6,2003)I,M,N,TEDEGE(N),X(M)
FORMAT(1H,'ERROR:---EDGE COORDINATES CHANGING TOO QUICKLY FOR',
&M,'I5,' & N='I5','THETA AT EDGE IS ',F8.2,' RAD. AND T(M) IS ',
&I4,F8.2)
CONTINUE

REPLACE X(M+1) WITH AVERAGE OF TEDGE(N) & TEDGE(N+1)
TEAVG=(TEDSE(N)+TEDGE(N+1))/2.

IF TEAVG > X(M+1),CONTINUE TO INTEGRATE OVER THIS SLIGHTLY
LARGER DELTA S (PHASE ERROR STILL SMALL)

TDIFF=X(M+1)-TEAVG
IF(TDIFF.LE.(X(M)-X(M+1))/4.))WRITE(6,2004)M,N
FORMAT(1H,'ERROR:---TOO MUCH PHASE WIND AT EDGE FOR M=',
&I5,' AND N='I5,' INTEGRATION CONTINUING.')
RECOMPUTE DX
DSE=DX*(TEAVG-X(M))
NCE=1
IF(J.NE.5)GO TO 40
TEAVG=TEAVG/0.01745329
WRITE(6,2002)TEAVG,M,N
FORMAT(1H,'THETA EDGE AVG IS ',F10.4,' FOR M,N = ',2I5)
CONTINUE
NPL=N+1
MNN=N
U= (TEAVG-X(M))/(X(M+1)-X(M))-1

FOR UPPER CORNERS OF (TRAPEZOIDAL) DELTA(S),INTERPOLATE R
AND F TO TEAVG. USE 3 TERM INTERPOLATION ;2 TERMS AT LEAST.
DO 51 NN=MNN,NPL
STORE PATHLENGTH VALUE OF ORIGINAL DELTA(S) CORNERS.
RTEMP(NN-MNN+1)=R(M+1,NN)
USE 2 TERMS IN INTERPOLATION IF MC3
IFM .LT. 3)GO TO 90
R(M+1,NN)=R(M+1,NN)+U*(R(M+1,NN)-R(M+1,NN+1))+S*U*(U-1)*(R(M+1,NN)
&-2)*R(M ,NN)+R(M-1,NN))
GO TO 91
R(M+1,NN)=R(M+1,NN)+U*(R(M+1,NN)-R(M,NN))
DO 54 L=1,3
STORE F VECTOR VALUE OF ORIGINAL DELTA(S) CORNERS.
FVTEMP(NN-MNN+1,L)=F(M+1,NN,L)
USE 2 TERMS IN INTERPOLATION IF MC3
IF(M .LT. 3)GO TO 92
F(M+1,NN,L)=F(M+1,NN,L)+U*(F(M+1,NN,L)-F(M,NN,L))+S*U*(U-1)*
&F(M+1,NN,L)-2.*F(M,NN,L)+F(M-1,NN,L))

```

```

GO TO 54
F(M+1,NN,L)=F(M+1,NN,L)+U*(F(M+1,NN,L)-F(M,NN,L))+S*U*(U-1)
CONTINUE
IF(J.NE.5)GO TO 53
WRITE(6,2007)M,NN,R(M-1,NN),S(M,NN),R*E*P(NN-MNN+1)*R(M+1,NN)
FORMAT(1H,'2I3+5(1X,E12.6)')
CONTINUE
GO TO 80
DSE=DX*(X(M+1)-X(M))
CONTINUE
A LEAST SQUARES BEST FIT IS APPLIED TO R UNTO DELTA(S)
R1=R(M+1,NN)-R(M,N)
R2=R(M,N+1)-R(M+1,N)
R3=R(M,N)*R(M+1,N)
BE=0.5*(R1+R2)
CE=0.5*(R1-R2)
AL=0.5*(R3-CE)
EVALUATE DELTA(I) EXPRESSION TO FIND SCATTERED FIELD
CONTRIBUTION FROM EACH DELTA(S).
IF(ABS(BE)-0.01)100,100,110
F1=BE*0.33333333
F1I=BE*0.5
F1R=L*0-F1I*F3I
F3R=0.5-BE*BE/8.0
GO TO 140
SINBE=SIN(BE)
COSBE=COS(BE)
F1R=F1R-COSBE/BE
F1I=SINBE/CE
F1=(1.0-COSBE)/EC
F3R=F1R-F1I/BE
F3I=(F1R-COSBE)/BE
IF(ABS(CE)-0.01)150,150*160
F4I=0.33333333*CE
F2I=0.5*CE
F2R=L*0-F2I*F4I
F4R=0.5-CE*CE/8.0
GO TO 170
SINCE=SIN(CE)
COSCE=COS(CE)
F2R=SINCE/CE
F2I=(1.0-COSCE)/CE
F4R=F2R-F2I/CE
F4I=(F2R-COSCE)/CE
SINAL=SIN(AL)
COSAL=COS(AL)
F1I=CPLX(F1R*F2R-F1I*F2I)+F1I*F2I*F1R
F1R=CPLX(F2R*F3R-F2I*F3I)+F2I*F3I*F2R
F1I=CPLX(F1R*F4R-F1I*F4I)+F1I*F4I*F1R
DO 202 L=1,3
A LEAST SQUARES BEST FIT IS APPLIED TO F UNTO DELTA(S)

```

```

T1=F(M+1,N+1,L)-F(M,N,L)
T2=F(M,N+1,L)-F(M,N,L)
T3=F(M,N,L)-F(M+1,N,L)
B=T1+T2
C=T1-T2
A=T3-0.5*C
SUM=A*F12+B*F23+C*F14
TOT=CMPLX(COSAL+SINAL)*SUM*DS
C
C      ADD PATTERN CONTRIBUTION FROM EACH DELTA(S) OF INTEG GRID.
C
C      202 STOT(L)=STOT(L)+TOT
C      IF(NCE.EQ. 0)GO TO 201
C
C      FOR NCE=1,RE-ASSIGN STORED R & F VALUES AND EXIT THETA LOOP
C
      DO 191 NN=NN,NPI
        R(M+1,NN)=RTEMP(NN-NNN+1)
        DO 191 L=1,3
          F(M+1,NN,L)=FVTEMP(NN-NNN+1,L)
        GO TO 200
      CONTINUE
      CONTINUE
      RETURN
      END
191
201
200
END

SUBROUTINE SETUP(NGL,I,JMAX,KMAX,MMAX,NMAX)
00010000
00020000
00030000
00040000
00050000
00060000
00070000
00080000
00090000
00100000
00110000
00120000
00130000
00140000
00150000
00160000
00170000
00180000
00190000
00200000
00210000
00220000
00230000
00240000
00250000
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00310000
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00490000
00500000
00510000
00520000
00530000
00540000
00550000
00560000
00570000
00580000

GIVEN THE # ANGLE VALUES,AN INITIAL VALUE AND AN INCREMENT,
THIS SUB. ESTABLISHES THE OUTPUT GRID(#2) AND UP TO 21
INTEGRATION GRIDS. IT ALSO PRCOMPUTES TRIG VALUES ON ALL
POINTS OF THE GRIDS.

DIMENSION MM(21),NN(21),T(36,21),P( 91,21)
COMMON/GRID1/SIT(36),COT(36),SIP( 91),COP( 91),TGI(36),PG1( 91)
COMMON/GRID2/SITT(181),COTT(181),STPP(51),COPP(51),TT(181),PP(51)

      SOME VARIABLE DEFINITIONS:

      T,P = THETA,PHI OF INTEG GRID
      MMAX,NMAX = # OF T,P VALUES
      TT,PP = SIGTETA,SIGPHI OF OUTPUT GRID
      JMAX,KMAX = # OF TT,PP VALUES
      NGL = # INTEGRATION GRIDS

      READ IN OUTPUT GRID
      READ(5,1001)JMAX,TT1,DTT
      IF(DTT)10,10,10
      DO 15 J=1,JMAX
        FJ=J-1
        TT(J)=TT1+DTT*FJ
        GO TO 30
      READ(5,1002)(TT(J),J=1,JMAX)
      READ(5,1001)KMAX,PP1,CPP
      IF(DPP)40,40,40
      DO 45 K=1,KMAX
        FK=K-1
        PP(K)=PP1+FK*CPP
        GO TO 60
      READ(5,1002)(PP(K),K=1,KMAX)
      CONTINUE

      INPUT INTEGRATION GRIDS
      READ(5,1001)NGI
      DO 100 I=1,NGI
        READ(5,1001)MM(I),T1,DT
        MMM=MM(I)
        IF(DT)110,110,110
        DO 115 M=1,MMM
          PM=M-1
          T(M,I)=T1+FM*DT
          GO TO 130
        READ(5,1002)(T(M,I),M=1,MMM)
        NNN=NN(I)
        IF(DP)140,140,140
        DO 145 N=1,NNN
          PN,N,I)=PI+FN*DP
          GO TO 100
      END

```

```

150 READ(5,1002)(P(I),I),N=1,NNN)
100 CONTINUE
C
C PRINT OUT GRID DATA
WRITE(6,2001)
WRITE(6,2006)(J,TT(J),J=1,JMAX)
WRITE(6,2002)
WRITE(6,2006)(K,PP(K),K=1,KMAX)
WRITE(6,2004)NG1
DO 155 I=1,NG1
WRITE(6,2005)I
:MM=MM(I)
NNN=NN(I)
WRITE(6,2006)(M,T(M),I),M=1,MM)
WRITE(6,2002)
WRITE(6,2006)(N,PI(N),I),N=1,NNN)
C
C COMPUTE GRID2 TABLE
DTR=0.017453293
DO 200 J=1,JMAX
TT(J)=TT(J)+DTR
SITT(J)=SIN(TT(J))
COTT(J)=COS(TT(J))
CONTINUE
DO 210 K=1,KMAX
PP(K)=PP(K)+DTR
COP(K)=COS(PP(K))
SIP(K)=SIN(PP(K))
CONTINUE
C
C COMPUTE GRID1 TABLES
I=0
FOR NG1 > 1 RE-ENTER SUB. HERE
ENTRY RESET(I,NMAX,NMAX)
I=I+1
MMAX=MM(I)
DO 400 M=1,MMAX
TGI(M)=DTR*(M-I)
SIT(M)=SIN(TGI(M))
COT(M)=COS(TGI(M))
CONTINUE
NMAX=NN(I)
DO 410 N=1,NMAX
PGI(N)=DTR*(N-I)
SIP(N)=SIN(PGI(N))
COP(N)=COS(PGI(N))
CONTINUE
RETURN
C
1001 FORMAT(15,F10.2)
1002 FORMAT(8F10.2)
2001 C THE FOLLOWING OUTPUT GRID HAS BEEN ESTABLISHED,/,
C ALL ANGLES IN DEGREES,/,/
2002 FORMAT(//)

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00590000
00600000
00610000
00620000
00630000
00640000
00650000
00660000
00670000
00680000
00690000
00700000
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00720000
00730000
00740000
00750000
00760000
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00780000
00790000
00800000
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00870000
00880000
00890000
00900000
00910000
00920000
00930000
00940000
00950000
00960000
00970000
00980000
00990000
01000000
01010000
01020000
01030000
01040000
01050000
01060000
01070000
01080000
01090000
01100000
01110000
01120000
01130000
01140000
01150000
01160000

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2004 FORMAT(, THE FOLLOWING, I3, INTEGRATION GRIDS HAVE BEEN ,
C ESTABLISHED, /)
2005 FORMAT(1H0, SEGMENT NUMBER, I3, /)
2006 FORMAT( 5(I4, F8.3))
END

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01170000
01180000
01190000
01200000
01210000

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SUBROUTINE FIELDS
C
C THIS SUB. FINDS THE MAGNETIC FIELD ON A SURFACE S BY
C EVALUATING A SPHERICAL WAVE EXPANSION. THE WAVE COEFFICIENTS
C AS CALCULATED BY ANOTHER PROGRAM ARE OBTAINED VIA A DISK
C DATA SET
C
COMMON/GRID1/SIT(36),COT(181),SIP( 91),COP( 91),T(36),PI( 91)
DIMENSION FI( 81),GI( 81)
DIMENSION NAME(18)
DIMENSION AI( 80,2),BI( 80,2)

C SOME VARIABLE DEFINITIONS:
C LMAX = MAXIMUM MODE ORDER
C F,G = VECTORS OF LEGENDRE POLYNOMIAL WEIGHTS EVALUATED FOR EACH OF
C MODE COMPONENT. THE ARGUMENT OF THE POLY IS THE COS OF
C A POLAR ANGLE.
C A,B = TE & TM WAVE COEFFICIENTS
C
C INPUT WAVE COEFFICIENTS FROM COMPUTER DISK
READ(9,1001)NAME
WRITE(6,2001)NAME
READ(9,1002)LMAX,MCOMP
READ(9,1003)NAME
FMC=MCOMP
WRITE(6,2003)MCOMP
WRITE(6,2002)NAME
WRITE(6,2004)(J,A(J,1),A(J,2),B(J,1),B(J,2),J=1,LMAX)
RETURN
C
CCCC
CCCC
C ENTRY POINT FOR MAGNETIC FIELDS AT FINITE R
ENTRY FIELD(MMAX,NMAX,HR,HT,HP,R)
DIMENSION R(36,91)
COMPLEX HR(36,91),HT(36,91),HP(36,91)
C
C MORE VARIABLES:
C R = RHO
C HR,HT,HP = INCIDENT MAGNETIC FIELD COMPONENTS (RETURNED
C TO FIELD)
C MMAX,NMAX = # OF THETA,PHI POINTS ON INTEG GRID
C
IENT=1
M=0
C
C ENTER THETA LOOP
C
M=N+1
SN=SIT(M)
Z=COT(M)
TOUT=T(M)
C
01220000 01230000 01240000 01250000 01260000 01270000 01280000 01290000 01300000 01310000 01320000 01330000 01340000 01350000 01360000 01370000 01380000 01390000 01400000 01410000 01420000 01430000 01440000 01450000 01460000 01470000 01480000 01490000 01500000 01510000 01520000 01530000 01540000 01550000 01560000 01570000 01580000 01590000 01600000 01610000 01620000 01630000 01640000 01650000 01660000 01670000 01680000 01690000 01700000 01710000 01720000 01730000 01740000 01750000 01760000 01770000 01780000 01790000
GO TO 99
C
C ENTRY POINT FOR ELECTRIC FIELDS AT INFINITE R
ENTRY FIELD2(JO,KO,ETTD,EPPD)
COMPLEX ETTD,EPPD
C
C STILL MORE VARIABLES:
C JO,KO = POINT ON OUTPUT GRID
C ETTD,EPPD = FIELD VALUES AT JO,KO (RETURNED TO MAIN)
C
IENT=2
SN=SIT(JO)
Z=COT(JO)
TOUT=T(JO)
C
C FORM LEGENDRE FUNCTION TABLES FOR M-TH THETA VALUE
*IF(ABS(SN)-.00001)200,100,100
DD 105 N=1,MCOMP
FINI=0
C
C OBTAIN F & G VECTORS FOR A GIVEN ORDER OF AZIMUTHAL VARIATION
NC=LMAX+1
CALL LEGEND(NC,MCOMP,Z,F)
DD 110 N=1,LMAX
T1=N-MCOMP+1
T2=N+1
GN1=T1*F(N+1)-T2*Z*F(N)
GN1=GN1/SN
DD 115 N=1,LMAX
FN1=F(N)/SN
C
C FOR NEAR FIELD CASE,MUST CONSIDER RADIAL DEPENDENCE OF FIELDS
C AS DEFINED BY THE SPHERICAL HANKEL FUNCTION.
GO TO (300,500),IENT
C
C SPECIAL EQUATIONS FOR THETA=0,180 DEG
IF(MCOMP-1)210,220,210
DD 215 N=1,LMAX
FINI=0
GN1=0
GO TO (300,500),IENT
DD 225 N=1,LMAX
FN=N*(N+1)
GN1=FN/2+0
GN1=FN/2+0
IF(TOUT-1,57)250,250,230
DD 235 N=1,LMAX,2
FN+1=-F(N+1)
GN1=-G(N)
GO TO (300,500),IENT
C
C FOR EACH PHI THE HANKEL FUNCTIONS ARE EVALUATED AT THE
01800000 01810000 01820000 01830000 01840000 01850000 01860000 01870000 01880000 01890000 01900000 01910000 01920000 01930000 01940000 01950000 01960000 01970000 01980000 01990000 02000000 02010000 02020000 02030000 02040000 02050000 02060000 02070000 02080000 02090000 02100000 02110000 02120000 02130000 02140000 02150000 02160000 02170000 02180000 02190000 02200000 02210000 02220000 02230000 02240000 02250000 02260000 02270000 02280000 02290000 02300000 02310000 02320000 02330000 02340000 02350000 02360000 02370000

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```

C CORRESPONDING VALUE OF R FOR EACH MODE COMPONENT.
C
C 300
DO 450 N=1,NMAX
HRR=0
HRI=0
HTR=0
HTI=0
HPR=0
HPI=0
CALL SPHANK(I,R(M,N),SOR,SOI)
DO 400 L=1,LMAX
FL=L
FN=(FL+1.0)/R(M,N)
NC=L+1
CALL SPHANK(NC,R(M,N),SIR,SII)
C
C TIME FOR SOME GLOP
C DETERMINE THE FIELD AT EACH THETA,PHI ON THE INTEG. GRID BY
C FIRST OBTAINING THE PRODUCT OF COEFFICIENTS AND THE SPHANK
C VALUES FOR EACH MODE COMPONENT.
C
C FIR=-AIL,I)*SIR+AIL,2)*SII
C FII=-AIL,I)*SII-AIL,2)*SIR
C F2R=FN*(AIL,1)*SOI+AIL,2)*SOR
C F2I=FN*(AIL,2)*SOI-AIL,1)*SOR
C F3R=-BIL,1)*SOR+BIL,2)*SOI
C F3I=-BIL,1)*SOI-BIL,2)*SOR
C
C THEN MULTIPLY BY THE CORRESPONDING MODE LEGENDRE POLY VALUES.
C FOR EACH THETA,PHI THE ABOVE VALUES ARE SUMMED OVER ALL MODE
C COMPONENTS.
C
C HRR=HRR+F2R*FII)*FL
C HRI=HRI+F2I*FII)*FL
C FII)=FII)*FMC
C HTR=HTR+GII)*FIR+F2R*G(L)+F3R*F(L)
C HTI=HTI+GII)*F2I*G(L)+F3I*F(L)
C HPR=HPR-FIR*F(L)-F2R*G(L)-F3R*G(L)
C HPI=HPI-FII*F(L)-F2I*G(L)-F3I*G(L)
C SOR=SIR
C SOI=SII
C
C 400
C CONTINUE
C
C FINALLY,THE PHI EXPANSION TERM IS INTRODUCED,COMPLETING THE
C EXPANSION.
C
C HRR=HRR*SM*COS(FMC*P(N))
C HRI=HRI*SM*COS(FMC*P(N))
C HRIM,N)=CMPLX(HRR,HRI)
C HTR=HTR*COS(FMC*P(N))
C HTI=HTI*COS(FMC*P(N))
C HTIM,N)=CMPLX(HTR,HTI)
C HPR=HPR*SIN(FMC*P(N))
C HPI=HPI*SIN(FMC*P(N))
C HRIM,N)=CMPLX(HPR,HPI)
C HRI TE(6,2005)M,N,HR(M,N),HT(M,N),HP(M,N)
C FORMAT(1H,'H-FIELDS',215,6F12.5)
C 2005
C 450
C CONTINUE

```

```

IF(M-VMAX*11,460,460
RETURN
C
C 460
C FAR FIELD CASE
C
C 500
C CONTINUE
C ETR=G*O
C ETI=O
C EPR=O
C EPI=O
C DD 550 L=1,LMAX
C F(L)=F(L)*FMC
C
C SUM PRODUCT OF COEFFICIENTS AND LEGENDRE WEIGHTS OVER ALL MODE
C COMPONENTS. THIS IS DONE FOR EACH OUTPUT GRID POINT.
C
C ETR=ETR+AIL,1)*F(L)+B(L,1)*G(L)
C ETI=ETI+AIL,2)*F(L)+B(L,2)*G(L)
C EPR=EPR+AIL,1)*G(L)+B(L,1)*F(L)
C EPI=EPI+AIL,2)*G(L)+B(L,2)*F(L)
C CONTINUE
C
C 550
C INTRODUCE BIGPHI VARIATION TO COMPLETE EXPANSION.
C
C ETR=ETR*SIN(FMC*P(KO))
C ETI=ETI*SIN(FMC*P(KO))
C ETO=CMPLX(ETR,ETI)
C EPR=EPR*COS(FMC*P(KO))
C EPI=EPI*COS(FMC*P(KO))
C EPO=CMPLX(EPR,EPI)
C RETURN
C
C 1001 FORMAT(118A4)
C 1002 FORMAT(15I5)
C 1003 FORMAT(15,2E17.8,2X,2E17.8)
C 2001 FORMAT(11H,' FIELD DATA INPUT IN THE FORM OF SPHERICAL WAVE ',
C *COEFFICIENTS,/,5X,18A4,/)
C 2002 FORMAT(1H05X,18A4,/,/,20X,4HAIN),22X,4HR(N),/,
C 5H N,7X,4HREAL,13X,4HIMAG,15X,4HREAL,13X,4HIMAG)
C 2003 FORMAT(1H,' AZIMUTHAL ORDER MCOMP=',I2)
C 2004 FORMAT(15,2E17.8,2X,2E17.8)
C END

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02380000
02390000
02400000
02410000
02420000
02430000
02440000
02450000
02460000
02470000
02480000
02490000
02500000
02510000
02520000
02530000
02540000
02550000
02560000
02570000
02580000
02590000
02600000
02610000
02620000
02630000
02640000
02650000
02660000
02670000
02680000
02690000
02700000
02710000
02720000
02730000
02740000
02750000
02760000
02770000
02780000
02790000
02800000
02810000
02820000
02830000
02840000
02850000
02860000
02870000
02880000
02890000
02900000
02910000
02920000
02930000
02940000
02950000

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02960000
02970000
02980000
02990000
03000000
03010000
03020000
03030000
03040000
03050000
03060000
03070000
03080000
03090000
03100000
03110000
03120000
03130000
03140000
03150000
03160000
03170000
03180000
03190000
03200000
03210000
03220000
03230000
03240000
03250000
03260000
03270000
03280000
03290000
03300000
03310000
03320000
03330000
03340000
03350000
03360000
03370000

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SUBROUTINE PATHL(RHO,J,K,M,N,M,NMAX,GAM)
C
C THIS SUB. COMPUTES PATH LENGTH FUNCTION GAMMA
C FOR EACH POINT ON THE INTEC GRID AND IN THE DIRECTION OF AN
C OUTPUT GRID POINT.
C I.E., THE PATHLENGTH IS A FUNCTION OF 4 VARIABLES.
C DIMENSION RHO(36,91),GAM(36,91)
C COMMON/GRID1/SIT(36),COT(36),SIP( 91),COP( 91),T(36),P( 91)
C COMMON/GRID2/SITT(181),COTT(181),SIPP(5),COPPI(5),TT(181),PP(5)
C
C VARIABLES:
C J,K = OUTPUT GRID POINT
C GAM = PATHLENGTH TERM (RETURNED TO MAIN)
C
C DO 10 M=1,M,NMAX
C T1=SIT(M)*SITT(J)*COPPI(K)
C T2=SIT(M)*SITT(J)*SIPP(K)
C T3=COT(M)*COTT(J)-1.0
C DO 10 N=1,N,NMAX
C GAM(M,N)=RHO(M,N)*(T1*COP(N)+T2*SIP(N)+T3)
C CONTINUE
C RETURN
C END

```

10

```

SUBROUTINE VECTOR(X,Y,AMP,PHI)
C
C THIS SUB. CONVERTS COMPLEX VALUES TO POLAR FORM
C AMP & PHI RETURNED
C
C=0.0
C IF(X)100,200,300
C IF(Y)110,120,120
C=360.0
C PHI=ATAN(Y/X)*57.29577951+180.0-C
C GO TO 400
C IF(Y)210,220,230
C PHI=-90.0
C GO TO 400
C PHI=0.0
C GO TO 400
C PHI=90.0
C GO TO 400
C PHI=ATAN(Y/X)*57.29577951
C AMP=SQRT(X*X+Y*Y)
C RETURN
C END

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03380000
03390000
03400000
03410000
03420000
03430000
03440000
03450000
03460000
03470000
03480000
03490000
03500000
03510000
03520000
03530000
03540000
03550000
03560000
03570000
03580000
03590000
03600000
03610000
03620000

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03630000
03640000
03650000
03660000
03670000
03680000
03690000
03700000
03710000
03720000
03730000
03740000
03750000
03760000
03770000
03780000
03790000
03800000
03810000
03820000
03830000
03840000

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```

SUBROUTINE SPHANK(N,R,X+Y)
C
C THIS SUB. COMPUTES VALUE OF SPHERICAL HANKEL FUNCTION TIMES
C THE FACTOR  $(-j)^{n+1} \cdot \rho^0 \cdot \exp(j \cdot \rho \cdot h)$ . THIS FACTOR IS THE
C RECIPROCAL OF THE FAR FIELD SPHERICAL HANKEL FUNCTION. THE
C FAR FIELD DEPENDENCE IS BUILT-IN TO THE SPHERICAL WAVE
C COEFFICIENTS AND MUST BE REMOVED FOR NEAR FIELD CASES.
C THE HANKEL FUNCTION PROVIDES THE NEAR FIELD RADIAL VARIATION
C OF THE PATTERN.
C
C DOUBLE PRECISION: TERM, AI, AR
C
C VARIABLES:
C N = MODE ORDER + 1
C R = ARGUMENT OF FUNCTION = PHASE CONSTANT * RHO
C X+Y = REAL & IMAG COMPUTED VALUES (RETURNED TO FIELD)
C
C AR=0
C AI=0
C PI=3.1415927
C K=0
C TERM=1
C GO TO 100
C K=K+1
C T1=N*K
C T2=N-K+1
C T3=2*K
C TERM=TERM*T1*T2/T3
C TERM=TERM/R
C
C THE NUMERICAL TECHNIQUE CAN BLOW-UP FOR R<50 (ROUGHLY) AND IS
C THEREFORE TRUNCATED BEFORE OVERFLOW OCCURS. THE IMAG COMPONENT
C SHOULD BLOW-UP (A MATH. PROPERTY OF THE FUNCTION) AND IS
C PROPERLY FOUND TO DO SO. THE REAL VALUE SHOULD GO TO ZERO AS
C THE IMAG VALUE BECOMES VERY LARGE. HOWEVER, THIS IS NOT THE CASE
C IT TOO BLOWS-UP.
C
C IF (DABS(TERM) .LT. .000001) GO TO 1000
C GO TO(200,100),160
C AR =AR+TERM
C IGO=1
C IF(K-N)20,1000,1000
C AI=AI-TERM
C IGG=2
C TERM=-TERM
C IF(K-N)20,1000,1000
C X=AR
C Y=AI
C RETURN
C END
100
200
1000

```

```

SUBROUTINE ADJUST(PHI)
C
C THIS SUB. SHIFTS PHI UNTIL IT LIES IN THE RANGE -180,180
C
C 1 IF(PHI-180.)20,20,10
C PHI=PHI-360
C GO TO 1
C 20 IF(PHI+180.)30,40,40
C PHI=PHI+360
C GO TO 20
C 40 RETURN
C END

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03850000
03860000
03870000
03880000
03890000
03900000
03910000
03920000
03930000
03940000
03950000
03960000
03970000
03980000
03990000
04000000
04010000
04020000
04030000
04040000
04050000
04060000
04070000
04080000
04090000
04100000
04110000
04120000
04130000
04140000
04150000
04160000
04170000
04180000
04190000
04200000
04210000
04220000
04230000
04240000
04250000
04260000
04270000
04280000
04290000
04300000
04310000
04320000
04330000
04340000
04350000
04360000

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04370000
04380000
04390000
04400000
04410000
04420000
04430000
04440000
04450000
04460000
04470000
04480000

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SUBROUTINE LEGEND(NMAX,M,Z,VAL)
C
C THIS SUB. CALCULATES VALUES OF THE ASSOCIATED LEGENDRE
C FUNCTION WITH INDICES N=M TO N=NMAX. DO NOT USE NMAX LESS
C THAN M-1. VALUES CHECKED WITH TABULATED VALUES TO 5 PLACES
C THRU N=56 FOR M=1, AND N=10 FOR M=5.
C
DIMENSION VAL( 81)
DOUBLE PRECISION TERM1,TERM2,TERM3,7D
VARIABLES:
NMAX = MODE ORDER + 1
M = MCOMP = ORDER OF AZIMUTHAL VARIATION
Z = ARGUMENT OF FUNCTION = COS(THETA) OR COS(RIGHTTHETA)
VAL = VALUES OF LEGENDRE POLY FOR EACH MODE (RETURNED TO
FIELD1, FIELD2)
ZD=Z
FM=M
TERM1=0.0
IF(M)8,8,9
TERM2=1.0
GO TO 11
KMAX=2*M-1
TERM2=(1.0-ZD*ZD)**(FM/2.0)
DO 10 K=1,KMAX,2
FK=K
TERM2=TERM2*(2.0*FM-FK)
VAL(M)=TERM2
NN=NMAX-1
DO 20 N=M,NN
FN=N
TERM3=((2.0*FN+1.0)*ZD*TERM2-(FN*FM)*TERM1)/(FN*FM+1.0)
VAL(N+1)=TERM3
TERM1=TERM2
TERM2=TERM3
CONTINUE
RETURN
END
04890000
04900000
04910000
04920000
04930000
04940000
04950000
04960000
04970000
04980000
04990000
05000000
05010000
05020000
05030000
05040000
05050000
05060000
05070000
05080000
05090000
05100000
05110000
05120000
05130000
05140000
05150000
05160000
05170000
05180000
05190000
05200000
05210000
05220000
05230000
05240000
05250000
05260000
05270000
05280000
05290000
05300000
05310000
05320000
05330000
05340000
05350000
05360000
05370000
05380000
05390000
05400000
05410000
05420000
05430000
05440000
05450000
05460000
SUBROUTINE DIFF(I1,J, MMAX,NMAX,F, FK,DT,DP)
C
C THIS SUB. RETURNS TO SUPR THE VALUE OF FT OR FP. THESE DERIVS
C ARE COMPUTED USING FORWARD AND BACKWARD DIFFERENCES.
C
DIMENSION F(36*91),FX(36*91), ITERM(21*20)
SOME VARIABLES USED:
ITERM = COEFFICIENTS WITHIN A TERM OF THE DIFFERENCE EQUATION
NT = # TERM$ IN EQUATION
II = INTEG GRID #
J = INDICATOR OF WHETHER FT OR FP IS BEING SOUGHT
FX = EITHER FT OR FP
DTR=0.01745329
FOR II=1 CALCULATE TERM COEFFICIENT FOR NUMERICAL DERIVATIVE
FORMULA. THESE COEFFICIENTS ARE IDENTICAL TO THOSE OF A
SINOMIAL EXPANSION.
IF(II .NE. 1)GO TO 10
ITERM(1,1)=1
ITERM(2,1)=-1
DO 11 JJ=2,20
JJPL=JJ+1
DO 12 I=1,JJPL
IF(II .NE. 1)GO TO 13
ITERM(1,JJ)=1
GO TO 12
IF(II .NE. JJPL)GO TO 14
-ITERM(1,JJ-1)
GO TO 12
ITERM(1,JJ)=ITERM(1,JJ-1)-ITERM(1,JJ-1)
CONTINUE
DO 71 JJ=1,10
JJPL=JJ+1
WRITE(6,2001)(ITERM(I,JJ),I=1,JJPL)
FORMAT(1H, 'ZLI5)
CONTINUE
CONTINUE
IF J=-1, COMPUTE FP; OTHERWISE COMPUTE FT.
IF(J .EQ. -1)GO TO 60
CONTINUE
CALCULATE FP
DEN=DTR*0.0254
ENTER THETA & PHI LOOPS (PHI WITHIN THETA). PERFORM CALCULATION
FOR EACH POINT ON INTEG GRID.
DO 54 M=1,NMAX
DO 54 N=1,NMAX
FX(M,N)=0.

```



```

45 T(J)=DTR*(T(J))
IF(1C2)50,50,60
50 DO 55 J=1,JIN
TH=DR*EP(J)
EP(J)=E(J)*SIN(TH)
E(J)=E(J)*COS(TH)
TH=DR*HP(J)
HP(J)=H(J)*SIN(TH)
H(J)=H(J)*COS(TH)
CONTINUE
C
C ESTABLISH DIFFERENTIAL PATTERN VALUES AT EACH INPUT POINT OF
C PATTERN. VALUES ARE STORED IN VECTORS A & B AND ARE USED IN
C EVALUATING THE COEFFICIENT INTEGRALS.
C
A(1,1)=E(1)*T(1)
A(1,2)=EP(1)*T(1)
B(1,1)=H(1)*T(1)
B(1,2)=HP(1)*T(1)
DO 65 J=2,JIN
DTH=T(J)-T(J-1)
A(J,1)=DTH*(J)
A(J,2)=DTH*EP(J)
B(J,1)=DTH*H(J)
B(J,2)=DTH*HP(J)
CONTINUE
C
C FMC=MCOMP
DO 80 J=1,JIN
Z=COS(T(J))
DO 85 N=1,MCOMP
PM(N)=0.
NC=NMAX+1
CALL LEGEND(NC,MCOMP,Z,PM)
DO 80 I=1,NMAX
F(I,J)=FMC*PM(I)
T1=I-MCOMP+1
T2=I+1
G(I,J)=T1*PM(I+1)-T2*PM(I)
CONTINUE
C
C EVALUATE COEFFICIENT INTEGRALS BY MULTIPLICATION OF F & G
C MATRICES AND A & B MATRICES. ACOE & BCOE RETURNED FROM MULT.
C
CALL MULT(JIN,NMAX+2,F,A,ACOE,0,80,121,121,2,80,2)
CALL MULT(JIN,NMAX+2,G,B,ACOE,1,80,121,121,2,80,2)
CALL MULT(JIN,NMAX+2,F,B,BCOE,0,80,121,121,2,80,2)
CALL MULT(JIN,NMAX+2,G,A,BCOE,1,80,121,121,2,80,2)
C
C NORMALIZE COEFFICIENTS AND COMPUTE POWER CONTRIBUTED FROM EACH
C MODE.
C
PIOZZ=(3.1415927/2.0)*0.002655
PTOT=0.

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00740000
00750000
00760000
00770000
00780000
00790000
00800000
00810000
00820000
00830000
00840000
00850000
00852000
00860000
00870000
00880000
00890000
00900000
00910000
00920000
00930000
00940000
00950000
00960000
00970000
00980000
00990000
00991000
00992000
01000000
01010000
01020000
01030000
01040000
01050000
01060000
01070000
01080000
01090000
01100000
01110000
01120000
01130000
01140000
01150000
01151000
01160000
01170000
01180000
01190000
01200000
01210000
01220000
01221000
01230000
01240000
01250000

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```

DO 95 N=1,NMAX
PA(N)=0.
PB(N)=0.
FACT=1.
IF(MCOMP)92,92,90
DO 91 M=1,MCOMP
FF=(N-M+1)*(N+M)
FACT=FACT*FF
FF=2*N+1
FACT=FACT/FF
FF=2*N*(N+1)
FACT=FACT*FF
K=0
BT=BCOE(N,K)
AT=ACOE(N,K)
BCOE(N,K)=BCOE(N,K)/FACT
ACOE(N,K)=ACOE(N,K)/FACT
PA(N)=PA(N)+AT*ACOE(N,K)*PIOZZ
PB(N)=PB(N)+BT*BCOE(N,K)*PIOZZ
IF(K-1)93,93,94
CONTINUE
PTOT=PTOT+PA(N)+PB(N)
CONTINUE
C
C WRITE SOME INPUTS ON DISK DATA SET
WRITE(9,2223)NAMEJ
WRITE(9,2224)NMAX,MCOMP
WRITE(9,2223)NAME
OUTPUT COEFFICIENTS
WRITE(6,2001)NAMEJ
WRITE(6,2004)MCOMP
PSUM=0
DO 100 J=1,NMAX
PA(J)=PA(J)/PTOT
PB(J)=PB(J)/PTOT
PSUM=PSUM+PA(J)+PB(J)
C
C WRITE COEFFICIENTS ON DISK. COEFFICIENTS HAVE FAR-FIELD RADIAL
C DEPENDENCE BUILT IN.
WRITE(9,2222)J,ACOE(J,1),ACOE(J,2),BCOE(J,1),BCOE(J,2)
FORMAT(15,2E17.8,2X,2E17.8)
WRITE(6,2005)J,ACOE(J,1),ACOE(J,2),BCOE(J,1),BCOE(J,2),PA(J)
&PB(J),PSUM
WRITE(6,2007)PTOT
C
C COMPUTE FAR-FIELDS PATTERN OF SPHERICAL WAVE COEFF FOR
C COMPARISON WITH INPUT PATTERN.
READ(5,1002)JMAX,JJIN
FORMAT(1X,2I4)
JJ=0
POUTS=0.0
IC1=1
IC2=-1

```

```

01260000
01270000
01280000
01290000
01300000
01310000
01320000
01330000
01340000
01350000
01360000
01370000
01380000
01390000
01400000
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01700000
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01760000

```


NATIONAL RADIO ASTRONOMY OBSERVATORY
GREEN BANK, WEST VIRGINIA

ADDENDUM
TO
ELECTRONICS DIVISION INTERNAL REPORT No. 221

UPDATED DESCRIPTION OF USING THE JPL
PHYSICAL-OPTICS SCATTERING PROGRAM

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UPDATED DESCRIPTION OF USING THE
JPL PHYSICAL-OPTICS SCATTERING PROGRAM

James R. Lyons

I. Introduction

A field plotting routine, multiple-scattering capability, and several other features, have been added to the JPL program as described in EDIR No. 221 (September 1981). This memo is designed to briefly describe these additions and to provide the user with a convenient list of program inputs. Specifically, it is a revision of section III, Details of Using the Program, of report 221. It is suggested that the previous report be read before attempting to read and apply this memo to the program. Neither the basic structure nor the underlying theory of the program have been altered, but there are additional input parameters.

II. Added Program Features

A. Field plotting routine.

A subroutine, EPLOT, has been written which will plot the amplitude (dB) and phase of a particular component of an electric or magnetic field. The spherical wave expansion (SWE) program employs this subroutine to plot the E and H plane of the input field and of the far-field form of the spherical wave representation. The scattering (SCAT) program uses EPLOT to plot the polar, azimuthal, or radial component of the near-field of the incident pattern, and to plot the polar, azimuthal or Cartesian-Y component of the far-field form of the incident and scattered patterns. Only the first 30° of the far-field incident pattern can be plotted; for $\theta(1) > 30^\circ$, the plot and power of this pattern are skipped. Input parameters are used to choose the field component to be plotted.

In the output grid, the Y component of the electric field has the form:

$$E_Y(\theta, \phi) = E_\theta(\theta, \phi) \cos \theta \sin \phi + E_\phi(\theta, \phi) \cos \phi.$$

The $\cos \theta$ term induces an azimuthal variation on E_Y which would destroy any beam circularity. Since our currently relevant problems involve circularly symmetric beams, the amplitude of this term has been neglected. E_Y then has the form

$$E_Y(\theta, \phi) = E_\theta(\theta, \phi) \text{sign}(\cos \theta) \sin \phi + E_\phi(\theta, \phi) \cos \phi.$$

Note also that $|\cos \theta| \sim 1$ for most scattering situations.

B. Multiple scattering.

In some scattering problems, there exists several reflectors, so it becomes necessary to perform a spherical wave expansion on a scattered field and reflect this off yet another surface. SCAT can do this by writing the Y component of the far-field scattered field values it computes into a disk data set. The SWE program then reads in the values and determines the expansion. SCAT stores 31 values (for $\Delta\theta = 1.0^\circ$, this corresponds to 30°), either those on the "right side" of the beam or the mirror-image of those on the "left side" of the beam. In either case, the polar angle of the field values is translated such that the beam peak occurs at $\theta = 0^\circ$. Since only half the beam is stored, the entire beam must be assumed circular. If the scattered field is asymmetric with respect to its beam peak, individually expanding both beam halves will define the extremes within which the true beam must lie.

The following table shows the data flow for the two programs when scattering from multiple reflectors. The numbers are the data set reference numbers used in the READ and WRITE statements.

		<u>READ</u>	<u>WRITE</u>
First reflection	{SWE	Cards	9
	{SCAT	9	12
Subsequent reflections	{SWE	12	9
	{SCAT	9	12

Note: SCAT only stores the pattern values corresponding to the final ϕ of the output grid.

C. Power calculations.

Relative power calculations are made for the same patterns that can be plotted by EPLLOT. For the near-field of the incident pattern in SCAT, the power integral has the form

$$P = \frac{1}{2\pi} \int_0^{2\pi} \int_0^{\theta(\phi)} \left[|H_{\theta}(\theta, \phi)|^2 + |H_{\phi}(\theta, \phi)|^2 \right] \sin \theta \, d\theta \, d\phi.$$

For the far-field patterns in SCAT and the patterns in the SWE program, the integral is:

$$P = \int_0^{\theta_0} \left[|E_{\theta}(\theta, \phi_0)|^2 + |E_{\phi}(\theta, \phi_0)|^2 \right] \sin \theta \, d\theta.$$

In the near-field integral, the power is calculated out to the reflector edge associated with each azimuthal angle. The cumulative power thru each integration grid is output along with the average value of the edge-angle of the integration grid. Power is also calculated for any section of the integration grid which lies beyond the reflector edge, thus allowing the spillover efficiency to be determined, if desired.

The far-field power integral can only be evaluated out to a constant polar angle. For the far-field incident pattern of SCAT, this angle is the average of the average edge-angles computed in the near-field case. For the scattered

pattern in SCAT and the patterns in the SWE program, the angle thru which the power is calculated is READ in by the programs.

The power contribution between two polar angles on the grid is evaluated at the midpoint of the angles:

$$\Delta P = \frac{1}{2} [|E(\theta_1)|^2 + |E(\theta_2)|^2] \cdot \sin \left(\frac{\theta_1 + \theta_2}{2} \right) \cdot \Delta\theta.$$

When the integration limit lies between grid values, ΔP is linearly scaled to the limit value. For the incident and SWE patterns, ΔP is doubled for each annulus because the pattern is actually that of a half-beam.

D. Field interpolation.

To accurately represent an input pattern, the SWE program requires values to be input at least every 0.25° or 0.5° in polar angle. Since most patterns are reasonably "smooth" thru the first 20 dB or so, points are input to the program every 1.0° or 2.0° and a linear interpolation is used to include points every 0.25° or 0.5° , respectively. To better represent the rounded beam peak, the first three interpolated values (i.e., .25, .5, .75, or .5, 1.0, 1.5) have a perturbation added to the interpolation.

The maximum number of points MAINOR can handle in computations is 121. Therefore, the maximum number of input points is 31 (e.g., 0° thru 30° , every 1.0°).

III. Input Variables and Their Formats

For both programs this section lists the input variables, gives a brief description of each variable, provides a concise table of inputs and their formats, and lists the output produced by the program.

A. SWE program.

1. Input variables descriptions:

Angles and phase are in degrees. All inputs read-in by MAIN.

- | | | |
|--------------|---|--|
| NAMEJ | - | Alphanumeric: Use to identify program;
≤ 72 characters. |
| MCOMP | - | Order of azimuthal variation; usually = 1. |
| LMAX | - | Maximum mode order, ≤ 80; usually ≈ 70. |
| NAME | - | Alphanumeric: Use to identify reflector, pattern,
etc. |
| JIN | - | Number of pattern points used in program, ≤ 121;
= 4* (actual number of card inputs - 1) + 1. |
| IC1 | - | If ≤ 0, convert input field from dB to volts. |
| IC2 | - | If > 0, set input field phase to zero. |
| IC3 | - | If > 0, compute pattern from equations inserted |
| IC4 | - | If > 0, E and H plane patterns are identical;
input only E plane. |
| NDISK | - | Pattern parameter: = 1 if values are to be input
from cards; = 2 if values are read from disk data
set (#12). |
| IPLANE | - | Specifies plane: = 1 if E plane power is calculated
and E plane pattern is plotted; = 2 for H-plane. |
| ILOT1 | - | Input pattern print-out parameter: = 1 for a plot;
= -1 for numerical values; = -11 for both; = 0 for
neither. |
| ILOT2 | - | Output (SWE) pattern print-out parameter: Same as
ILOT1. |
| TEDGE | - | Polar angle thru which power flow is calculated. |
| T(J) | - | Polar angle of input pattern values. |
| E(J), H(J) | - | E, H-plane input pattern amplitude. |
| EP(J), HP(J) | - | E, H-plane input pattern phase. |
| JMAX0 | - | = 180/Δθ + 1, where Δθ is desired output increment
of the output (SWE) pattern; typically = 181. |
| JOUT | - | Number of output points starting with θ = 0°;
typically = 91. |

2. Table of inputs and their formats.

	Card No.	Variables	Format	
	1	NAMEJ	18A4	
	2	MCOMP, LMAX	2I5	
	3	NAME	18A4	
	4	JIN, IC1, IC2, IC3, IC4, NDISK	6I5	
	5	IPLANE, IPLOT1, IPLOT2	3I5	
	6	TEDGE	F10.5	
If NDISK = 1	}	IC4≤0: T(1), E(1), EP(1), H(1), HP(1)	5F10.5	
		7	IC4>0: T(1), E(1), EP(1)	3F10.5
		⋮		
		⋮		
		JW + 6	IC4≤0: T(JIN), etc.	5F10.5
			IC4>0: T(JIN), etc.	3F10.5
	JIN + 7	JMAXO, JOUT	2I5	
If NDISK = 2	}	No cards:	NAMEJ (18A4)	
		Data read from computer disk (#12)	T(1), E(1), EP(1) (3F10.5)	
			⋮	
			T(JIN), etc. (3F10.5)	
	7	JMAXO, JOUT	2I5	

3. SWE program print-out.

- Two alphanumeric statements.
- Input pattern: plot or numerical values.
- Power in input pattern.
- Real and imaginary values of SWE coefficients.
- Fraction of total mode power in the coefficients for each mode order.
- Total coefficient mode power (not related to computed power in pattern).
- Output pattern (far-field of SWE): plot or numerical values.
- Power in output pattern.

B. SCAT Program.

1. Input variable descriptions.

All linear measures are in meters; all angles are in degrees.

Parentheses indicate which subroutine is reading in data.

(MAIN)

- TITLE - Alphanumeric: describe reflector, wavelength, etc.,
 ≤ 72 characters.
- NDISK - disk writing parameter: = 1 if 1st half of scattered
 values are to be stored in a disk data set; = 2 for
 2nd half of pattern; = 0 if no values are to be stored.
- PC - phase constant = $2\pi/\lambda$
- XT, YT, ZT - translation to expected phase center of scattered
 pattern.
- ALPHA - used as a reflector rotation parameter when surface
 is specified analytically; set to 0.0 if not needed.
- RHO \emptyset - distance to the reflector at $\theta = 0^\circ$.
- THETA \emptyset - used as an edge parameter when theta-edge is specified
 analytically; set to 0.0 if not needed.

- - - - -

(SETUP)

- JMAX - number of θ values in output grid; ≤ 361 .
- TT1 - initial θ value.
- DTT - θ increment, usually = 1.0° .
- KMAX - number of ϕ values in output grid; ≤ 46 .
- PP1 - initial ϕ value.
- DPP - ϕ increment.
- NG1 - number of integration grids, ≤ 21 .
- MM(I) - number of θ values on Ith integration grid; ≤ 36 .
- T1 - initial θ value.
- DT - θ increment, typically between 0.2° and 1.0° .

- NN(I) - number of ϕ values on Ith integration grid; ≤ 91 .
 P1 - initial ϕ value.
 DP - ϕ increment, typically between 2° and 10° .

(Note: Subroutine EPL0T and the power calculations assume that DT and DP are constant within an integration grid, and that DTT is constant for the output grid.)

- - - - -
(FIELDS)

- NAME - alphanumeric: identify SWE program.
 LMAX - maximum mode order.
 MCOMP - order of azimuthal variation.
 NAME - alphanumeric: identify reflector, pattern, etc.
 J - 1 to LMAX.
 A(J,1), A(J,2) - real and imaginary components of SWE "A" coefficients.
 B(J,1), B(J,2) - real and imaginary components of SWE "B" coefficients.

- - - - -
(MAIN)

- IEDGE - reflector parameter used to determine if and how program obtains the reflector edge-values of θ . Cases:
 < 0 - calculate values using equations inserted in EDGEEQ. READ in 1 value, TEDGE1, to be used in equations.
 = 0 - θ at edge is a constant. READ 1 value, TEDGE(1).
 = 11 - edge will not be encountered on present integration grid. READ no values.
 > 0 - θ at edge is in tabular form and is read in for each ϕ on the integration grid. READ in NMAX values.
 TEDGE(N) - edge value of θ . Read in either NMAX, 1, or no values.

- - - - -

(SURF)

- ISURF - reflector parameter used to determine how $\rho(\theta, \phi)$, $\partial\rho/\partial\theta(\theta, \phi)$, and $\partial\rho/\partial\phi(\theta, \phi)$ are obtained. Cases:
- < 0 - ρ , $\partial\rho/\partial\theta$, $\partial\rho/\partial\phi$ all determined analytically ($\neq -2$) from equations in SURF. READ no values.
 - = -2 - ρ determined from equations; derivatives numerically computed. READ no values.
 - = 0 - ρ is in tabular form and is a function of θ ; derivatives numerically computed. READ MMAX values.
 - 0 - ρ is in tabular form and is a function of θ and ϕ ; derivatives numerically computed. READ MMAX x NMAX values.
- F(M, N) - $\rho(\theta, \phi)$. READ either MMAX x NMAX, MMAX, or no values.

- - - - -

(MAIN)

- NIPLLOT - near-field incident pattern parameter:
- = 1, 2, 3 - plot $H_\theta(\theta, \phi_0)$, $H_\theta(\theta, \phi_0)$, $H_\phi(\theta, \phi_0)$, respectively.
 - = 0 - no plot.
- PHIPLT - ϕ_0 used with NIPLLOT: ϕ_0 must be on present integration grid segment.
- FIPLLOT - far-field incident pattern parameter:
- = 1 - plot $E_Y(\theta, \phi_0)$ = Y component.
 - = 2, 3 - plot $E_\theta(\theta, \phi_0)$, $E_\phi(\theta, \phi_0)$, respectively.
 - = -1 - numerical values listed.
 - = 0 - neither plot nor numerical values.
- FSPLLOT - far-field scattered pattern parameter.
(Same as FIPLLOT.)
- TEDGEA - half-angle from beam center, thru which power in scattered pattern is computed.

2. Table of inputs and their formats.

A dash (-) in the Card No. column indicates that the card number is not known or is difficult to determine.

Card No.	Variables	Format
1	TITLE	18A4
2	NDISK	I5
3	PC, XT, YT, ZT, ALPHA, RHO \emptyset , THETA \emptyset	7F10.4
4	JMAX, TT1, DTT	I5, 2F10.2
5	KMAX, PP1, DPP	I5, 2F10.2
6	NG1	I5
7	MM(1), T1, DT	I5, 2F10.2
8	NN(1), P1, DP	I5, 2F10.2
⋮	⋮	⋮
2NG1 + 5	MM(NG1), T1, DT	I5, 2F10.2
2NG1 + 6	NN(NG1), P1, DP	I5, 2F10.2
NO CARDS:	NAME	(18A4)
Data	LMAX, MCOMP	(I5)
read	NAME	(18A4)
from	1, A(1,1), A(1,2), B(1,1), B(1,2)	(I5, 2E17.8, 2X, 2E17.8)
computer	⋮	⋮
disk	LMAX, A(LMAX,1), A(LMAX,2)	(I5, 2E17.8, 2X, 2E17.8)
(#9).	B(LMAX,1), B(LMAX,2)	
Read for each integration grid. Repeat NG1 times.	IEDGE	I5
	TEDGE(N) (NMAX, 1, or no cards)	7(1X,F9.5)
	ISURF	I5
	F(M,N) (MMAX x NMAX, MMAX or no cards)	7(1X,F9.5)
	NIPLLOT, PHIPLT, FIPLLOT, FSPLLOT	I5,F7.2,2I5
	TEDGEA	F10.4

3. SCAT program print-out.

```

-- 2 alphanumeric statements followed by propa-
-- gation constant.
--  $\theta$  and  $\phi$  for each integration grid segment.
For each integration grid. {
-- integration grid number and value of IEDGE.
-- if edge is to be encountered, list of  $\theta$ -edge
-- and  $\rho$ -edge as functions of  $\phi$ .
-- integration grid number and value of ISURF.
-- plot of near-field of incident pattern.
-- near-field incident pattern power.
For each  $\phi$ . {
-- far-field of incident pattern: plot or numbers.
-- far-field incident pattern power.
-- phase center translations, scale factor (set
-- to 1.0).
For each  $\phi$ . {
-- far-field of scattered pattern: plot or numbers.
-- far-field scattered pattern power.

```

C. Submitting a job and organization of JCL.

The SCAT and SWE programs have been run on the IBM 4341 computer located in Charlottesville. The operating system used is called Pandora; blocks of program code (e.g., subroutines) are stored under various Pandora member names (PMN^S).

To submit a job with the Pandora system, the SUBMIT command is used followed by the Pandora members making up the program and data. All the Pandora members (except EPLLOT, which is just the plotting routine) have already been described in EDIR #221.

The submit command for the SWE program is SUBMIT_JCLSWE_MAINOR_EPLOT_LVM_DATASWE.

The JCL and the actual subroutines submitted are shown below. The PMN for a subroutine or block of subroutines is shown in the left-hand column.

PMN	JCL and Subroutines
	<pre> //SWAVE__JOB_(userI.D.),user name,MSGLEVEL=(2,0), CLASS=L,TIME=1 /*ROUTE__PRINT_REMOTE1 //_EXEC_FORTGCLG,ERROR=E,PARM.FORT=ID //FORT.SYSPRINT_DD_DUMMY //FORT.SYSIN_DD_* </pre>
MAINOR	[MAIN
ELOT	[Subroutine ELOT
LVM	<pre> Subroutines LEGEND, VECTOR, MULT /* //LKED.SYSPRINT_DD_DUMMY //GO.FT09F001_DD_DSN=user name.DATA,DISP=SHR //GO.FT12F001_DD_DSN=user name.DSCAT,DISP=SHR //GO.SYSIN_DD_* </pre>
DATASWE	[DATASWE
	[/*

A CLASS = L job uses a 216 K byte memory partition, double the size of a standard partition. The SWE program could be reduced to a standard partition by (carefully) reducing array sizes. If this is done, the corresponding dimensions in the CALL MULT statements in MAINOR must be changed along with the array dimensions in several subroutines.

For both SWE and SCAT jobs the TIME parameter is in (CPU) minutes.

The submit command for the SCAT program is

```
SUBMIT_JCLSCAT_MAINISC_EPLOT_SSURF_FINT_SFPVSALD_DATASCAT
```

The JCL and subroutine structure is as follows:

PMN	JCL and Subroutines
JCLSCAT	<pre>//SCAT_ _ _ _JOB_(userI.D.),user name,MSGLEVEL=(2,0), CLASS=0,TIME=user set ----- ----- Same as last 4 lines of JCLSWE ----- -----</pre>
MAINISC	[MAIN
EPLOT	[Subroutine EPLOT
SSURF	[Subroutine SURF
FINT	[Subroutine FINT
SFPVSALD	<pre>Subroutines SETUP, FIELDS, PATHL, VECTOR, SPHANK, ADJUST, LEGEND, DIFF ----- ----- ----- Same 5 lines of JCL as at end of member LVM ----- -----</pre>
DATASCAT	<pre>DATASCAT /*</pre>

A CLASS = 0 job is an extra large partition of 880 K bytes. SCAT (between 300 and 400 K bytes) could be reduced to an L job by reducing the array dimensions of the output grid and the integration grid, if necessary, and by decreasing the number of allowable integration grids. Note that the array dimensions in several subroutines must be changed.

D. Creating and listing disk data sets.

The following JCL is used to create a sequential data set on the computer disk.

```
//CREATEDS__JOB_(userI.D.),user name,CLASS=Q,
      MSGLEVEL=1
/*ROUTE__PRINT_REMOTE1
//NEWFILE___EXEC_PGM=(,CATLG),DSN=user name.DATA
      or DSCAT,UNIT=3300,SPACE=(CYL,(1,1)),
      DCB=(RECFM=FB,LRECL=80,BLKSIZE=1600).
```

To list the contents of a data set:

```
//LIST_JOB_(userI.D.),user name,CLASS=Q,MSGLEVEL=1
/*ROUTE__PRINT_REMOTE1
//_EXEC_LIST
//SYSIN_DD_DSN=user name.DATA or DSCAT,DISP=SHR
```

Only the above JCL need be submitted to perform the desired task.

Acknowledgements

I thank Rick Fisher for his many helpful suggestions and constructive criticisms and Carolyn Dunkle for typing this report.

IV. Changed and Additional Program Code.

```

G LEVEL 21          MAIN PROGRAM          11/15/21          DATE = 82211          MAIN          DATE = 82211          11/15/21
C THIS IS A SPHERICAL WAVE EXPANSION PROGRAM-ORTHOGONALITY
C TECHNIQUE. THE SPHERICAL WAVE COEFFICIENTS FOR THE M-TH
C FURTHER COMPONENT OF A RADIATION PATTERN ARE COMPUTED.
C THE EXPANSION MATCHES THE FAR-FIELD INPUT PATTERN WITH
C THE FAR-FIELD FORM OF SPHERICAL WAVES.
C DIMENSION NAME(18),NAMEJ(18)
C DIMENSION PA(80),PB(80)
C DIMENSION F(121),E(121),EP(121),H(121),HP(121)
C DIMENSION F(80,121),G(80,121),PM(81)
C DIMENSION A(121,2),B(121,2)
C DIMENSION ACCE(80,2),BCOE(80,2)
C DIMENSION ACUT(1,2),BSUT(1,2)
C DTR=G*017453293
C DT2=OTR/2*0
C SCME VARIABLE DEFINITIONS:
C T*E*EP*H*HP=INPUT PATTERN PARAMETERS
C JIN(<=121)=# THETA INPUT VALUES
C NMAX(<=80) = MAXIMUM MODE ORDER
C ACOE,BCOE = REAL & IMAG COMPONENTS OF TE AND TM WAVE COEFFICIENTS
C WRITTEN ON DISK
C A*B = AMP & PHASE OF INPUT FIELDS VALUES * DELTA(THETA)
C F*G = MATRICES (NMAX BY JIN OR NMAX BY JOUT OF VALUES RELATED TO
C ASSOCIATED LEGENDRE POLYNOMIALS
C G PROP. TO DERIVATIVE WRT THETA OF N-TH POLY.
C F PRCP. TO N-TH POLY.
C INPUT ALPHANUMERIC STATEMENTS,PATTERN & PLOTTING PARAMETERS.
C READ(5,2223)NAMEJ
C WRITE(6,2001)NAMEJ
C READ(5,1002)MCCMP,NMAX
C WRITE(6,2001)NAME
C READ(5,1002)JIN,IC1,IC2,IC3,IC4,NDISK
C REAC(5,1002)PLANE,IPLCT1,IPLCT2
C READ(5,2013)TELCG
C IF IC3>0,CALCULATE INPUT PATTERN PARAMETERS FROM EQUATION
C IF IC3<=0 AND IC4>0,READ IN ONLY T,E,EP, I,E.,PATTERN IS CIRC.
C IF IC3)113,113,133
C IF IC4)153,153,163
C CONTINUE
C READ-IN PATTERN FROM DISK DATA SET (#12)
C IF INDISK .EQ. 1)GC TC 110
C READ(12,2223)NAMEJ
C WRITE(6,2014)NAMEJ
C READ(12,2013)(J),E(J),EP(J),J=1,JIN*4)
C GC TC 115
G LEVEL 21          MAIN          DATE = 82211          11/15/21
110 CONTINUE
C READ-IN PATTERN FROM CARDS
C READ(5,2013)(T(M),E(M),EP(M),M=1,JIN*4)
C CONTINUE
115 CIRCULARLY SYMMETRIC BEAM
C DC 120 J=1,JIN*4
C H(J)=E(J)
C HP(J)=EP(J)
C CONTINUE
C GO TC 130
153 READ(5,2015)(T(M),E(M),EP(M),H(M),HP(M),M=1,JIN*4)
C FOR ICI < OR = 0 CONVERT FROM DB TO VOLTS
C IF(I(1)10,10,20
C DO 15 J=1,JIN*4
C E(J)=10*0**E(J)/20*0)
C F(J)=10*0**H(J)/20*0)
C CONTINUE
C INTERPOLATE PATTERN VALUES
C JINM=JIN-4
C DO 167 J=1,JINM*4
C JP2=J+2
C DO 166 K=J,JP2
C FD4=FLCAT(K+1-J)/4.
C T(K+1)=T(J)+FD4*(T(J+4)-T(J))
C E(K+1)=E(J)+FD4*(E(J+4)-E(J))
C EP(K+1)=EP(J)+FD4*(EP(J+4)-EP(J))
C H(K+1)=H(J)+FD4*(H(J+4)-H(J))
C HP(K+1)=HP(J)+FD4*(HP(J+4)-HP(J))
C ADD PERTURBATION
C IF(J.GT. 1)GO TO 166
C PERT=.25/FLOAT(K+J)
C E(K+1)=E(K+1)+PERT*(E(J)-E(J+4))
C H(K+1)=H(K+1)+PERT*(H(J)-H(J+4))
C CONTINUE
C GO TO 143
C CONTINUE
C DELT=0*25
C DC 123 M=1,JIN
C INSERT EQUATIONS FOR INPUT PATTERN HERE. NOTE ABOVE
C INITIALIZATIONS.
C T(M)=(M-1)*DELT
C E(M)=EXP(-T(M)*T(M)/31*4)
C H(M)=E(M)
0010000
0020000
0030000
0040000
0050000
0060000
0070000
0080000
0090000
0100000
0110000
0120000
0130000
0140000
0150000
0160000
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00980000
00990000
01000000
01010000
01020000
01025000
01030000
01040000

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G LEVEL 21          MAIN          DATE = 82211          11/15/21          11/15/21
C
C   NORMALIZE COEFFICIENTS AND COMPUTE POWER CONTRIBUTED FROM EACH
C   MODE.
C
C   PTOZZ=(3.1415927/2.0)*0.002655
C   PTOZ=0.
C   DO 95 N=1,NMAX
C   PA(N)=0.
C   PB(N)=0.
C   FACT=1.
C   IF(MCCMP)92,92,90
C   DO 91 M=1,MCCMP
C   FE=(N-M+1)*(N+M)
C   FACT=FACT*FE
C   FE=2*N+1
C   FACT=FACT/FE
C   FE=2*M*(N+1)
C   FACT=FACT*FE
C   K=0
C   K=K+1
C   BI=BCOE(N,K)
C   AT=ACOE(N,K)
C   BCOE(N,K)=BCOE(N,K)/FACT
C   ACOE(N,K)=ACOE(N,K)/FACT
C   PA(N)=PA(N)+AT*ACOE(N,K)*PIOZZ
C   PB(N)=PB(N)+BT*BCOE(N,K)*PIOZZ
C   IF(K-1)93,93,94
C   CONTINUE
C   PTOZ=PTOZ+PA(N)+PB(N)
C   CONTINUE
C
C   WRITE A FEW PROGRAM INPUTS INTO A DISK DATA SET (#9)
C
C   WRITE(9,2223)NAMEJ
C   WRITE(9,1002)NMAX,MCCMP
C   WRITE(9,2223)NAME
C
C   OUTPUT COEFFICIENTS
C
C   WRITE(6,2001)NAMEJ
C   WRITE(6,2004)MCOMP
C   PSUM=0
C   DO 100 J=1,NMAX
C   PA(J)=PA(J)/PTOT
C   PB(J)=PB(J)/PTOT
C   PSUM=PSUM+PA(J)+PB(J)
C
C   WRITE COEFFICIENTS ON DISK. COEFFICIENTS HAVE FAR-FIELD RADIAL
C   DEPENDENCE BUILT IN.
C
C   WRITE(9,2222)J,ACOE(J,1),ACOE(J,2),BCOE(J,1),BCOE(J,2)
C
C   WRITE(6,2005)J,ACOE(J,1),ACOE(J,2),BCOE(J,1),BCOE(J,2),PA(J)
C   &,PB(J),PSUM
C   WRITE(6,2007)PTOT
C
C   COMPUTE FAR-FIELDS PATTERN OF SPHERICAL WAVE COEFF FOR
C   COMPARISON WITH INPUT PATTERN.
C
G LEVEL 21          MAIN          DATE = 82211          11/15/21          11/15/21
C
C   READ(5,1002)JMAX,JIN
C   JC=0
C   PCUTS=0*0
C   IC1=1
C   IC2=-1
C   CT=JMAX-1
C   DT=180*0/DT
C   JPM=INT(TEDEGE/DT) +1
C   J=1
C   TH=0
C   IF(MCCMP-1)170,180,170
C   DO 175 N=1,NMAX
C   F(L,N)=0
C   G(L,N)=0
C   GO TO 215
C   DO 180 N=1,NMAX
C   FN=(N+1)
C   F(L,N)=FN/2*0
C   G(L,N)=FN/2*0
C   GO TO 215
C
C   ESTABLISH F & G MATRICES FOR OUTPUT THETAS:NMAX BY JOUT MATS.
C
C   FJ=J-1
C   TH=FJ*DT
C   Z=COS(DTR*TH)
C   S=SIN(DTR*TH)
C   DO 205 N=1,MCCMP
C   PMIN)=0
C   CALL LEGEND(FC,MCOMP,Z,PM)
C   DO 210 N=1,NMAX
C   F(L,N)=FMC*PM(N)/S
C   T2=N+1
C   G(L,N)=G(L,N)/S
C
C   MULTIPLY F & G MATS AND ACOE & BCOE TO DETERMINE FIELD VALUES
C   ACUT & BCUT AT EACH OUTPUT ANGLE.
C
C   CALL MULT(NMAX,1,2,F,ACOE,ADUT,0,80,121,80,2,1,2)
C   CALL MULT(NMAX,1,2,G,BCOE,ABUT,1,80,121,80,2,1,2)
C   CALL MULT(NMAX,1,2,F,BCOE,ABUT,0,80,121,80,2,1,2)
C   CALL MULT(NMAX,1,2,F,BCOE,ABUT,1,80,121,80,2,1,2)
C
C   CALL VECTOR(ACUT(1,1),ADUT(1,2),EAMP,EPHI)
C   CALL VECTOR(BCUT(1,1),BOUT(1,2),HAMP,HPHI)
C
C   IF(J .GT. JPLT)GO TO 270
C   IF((IPLOT2 .EQ. 0) .OR. (IPLCT2 .EQ. -1))GO TO 270
C
C   LABEL & PLCT OUTPUT PATTERN
C
C   IF(IJ .EQ. 1) .AND. (IPLANE .EQ. 1))WRITE(6,2018)
C   IF(IJ .EQ. 1) .AND. (IPLANE .EQ. -1))WRITE(6,2019)
C   IF(IPLANE .EQ. 1)CALL EPLDT(J,EAMP,EPHI,JPLT,DT,C,0,TH)
C   IF(IPLANE .EQ. -1)CALL EPLDT(J,HAMP,HPHI,JPLT,DT,C,0,TH)

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G LEVEL 21          EPL0T          DATE = 82211          11/15/21          G LEVEL 21          MULT          DATE = 82211          11/15/21
870 CONTINUE          01170000          C          SUBROUTINE MULT(M,N,K,A,B,C,IC,NA,MA,NB,MC,PC)
C          01180000          C          THIS SUB. MULTIPLIES 2 FULL MATRICES
C          01190000          C          DIMENSION A(NA,MA),B(NB,MB),C(INC,MC)
C          01200000          C          IF(I,C)10,20,30
C          01210000          10          DO 15 I=1,N
C          01220000          15          DO 15 J=1,M
C          01230000          20          C(I,J)=C(I,J)-A(I,L)*B(L,J)
C          01240000          25          RETURN
C          01250000          30          DO 25 I=1,N
C          01260000          35          DO 25 J=1,M
C          01270000          C          C(I,J)=C(I,J)+A(I,L)*B(L,J)
C          01280000          C          RETURN
C          01290000          C          DC 35 I=1,N
C          01300000          C          DC 35 J=1,M
C          01310000          C          DO 35 L=1,M
C          01320000          C          C(I,J)=C(I,J)+A(I,L)*B(L,J)
C          01330000          C          RETURN
C          01340000          C          DC 35 I=1,N
C          01350000          C          DC 35 J=1,M
C          01360000          C          DO 35 L=1,M
C          01370000          C          C(I,J)=C(I,J)+A(I,L)*B(L,J)
C          01380000          C          RETURN
C          01390000          C          END
C          01400000          C          END
C          01410000          C          END
C          01420000          C          END
C          01430000          C          END
C          01440000          C          END
C          01450000          C          END
C          01460000          C          END
C          01470000          C          END
C          01480000          C          END
C          01490000          C          END
C          01500000          C          END
C          01510000          C          END
C          01520000          C          END

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THIS SUBROUTINE WAS INADVERTENTLY LEFT OUT OF THE FIRST REPORT.


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G LEVEL 21          MAIN          DATE = 82211          11/19/25          G LEVEL 21          MAIN          DATE = 82211          11/19/25
C          IF(TTL .GE. 29.9)GO TO 755          03330000          DP=XT*SITT(J)*COPP(K)+YT*SITT(J)*SIPP(K)+ZT*(COTT(J)+1.0)          03720000
C          PERC=100.*EPW/TPW          03335000          DP=DP*PC*57.29578          03730000
C          WRITE(6,2301)EDGEA,EPW,PERC          03340000          ETPHI=ETPHI-DP          03740000
C          RAT=POWI/EPW          03350000          EPPHI=EPPHI-DP          03750000
C          PRINT-CUT RATIO OF NEAR TC FAR-FIELD INCIDENT POWER          03360000          EYPHI=EYPHI-DP          03760000
C          WRITE(6,2306)TEDGEA,RAT          03361000          C          ADJUST PHASES TO -180.180 RANGE          03770000
C          CONTINUE          03362000          C          CALL ADJUST(ETPHI)          03780000
C          TRANSLATE PHASE CENTER,SCALE FIELD AMPLITUDES,AND OUTPUT          03363000          C          CALL ADJUST(EPPHI)          03790000
C          TOTAL FIELDS          03370000          C          CALL ADJUST(EYPHI)          03800000
C          WRITE(6,2002)TITLE          03380000          C          SCALE FIELD AMPLITUDES          03810000
C          WRITE(6,2010)XT,YT,ZT,SCALE          03390000          C          ETAMP=ETAMP*SCALE          03820000
C          PARAMETERS USED IN WRITING SCATTERED FIELDS ON DISK          03400000          C          EPAMP=EPAMP*SCALE          03821000
C          TTPEAK=180.0-2.*ALPHA          03445000          C          EYAMP=EYAMP*SCALE          03822000
C          JPEAK=INT((TTPEAK-TTL)/DTT *.1) +1          03450000          C          IF(FSPLOT .EQ. -1)GO TO 780          03823000
C          READ-IN HALF-ANGLE FOR POWER CALCULATION          03460000          C          IF(IJ .NE. 1)GC TO 770          03830000
C          READ(5,1002)TEDGEA          03465000          C          IF(FSPLOT .EQ. 1)WRITE(6,2106)PO          03840000
C          DO 795 K=1,KMAX          03470000          C          IF(FSPLOT .EQ. 2)WRITE(6,2107)PO          03850000
C          PG=PP(K)/O.017453293          03475000          C          IF(FSPLOT .EQ. 3)WRITE(6,2108)PO          03860000
C          IF(FSPLOT .EQ. -1)WRITE(6,2011)PO          03480000          C          CONTINUE          03870000
C          IF(FSPLOT .EQ. 0)WRITE(6,2004)PO          03490000          C          IF(FSPLOT .EQ. 1)CALL EPLT(IJ,EYAMP,EYPHI,JMAX,DTT,TT1,TO)          03880000
C          TPOW=0.          03510000          C          IF(FSPLOT .EQ. 2)CALL EPLT(IJ,ETAMP,ETPHI,JMAX,DTT,TT1,TO)          03890000
C          EPW=0.          03520000          C          IF(FSPLOT .EQ. 3)CALL EPLT(IJ,EFAMP,EPPHI,JMAX,DTT,TT1,TO)          03900000
C          TRANSLATE ANGLES TOWARDS 0.0 DEG AT BEAM PEAK          03530000          C          CONTINUE          03910000
C          TDIFF=(180.-2.*ALPHA+.5*DTT)*DTR          03532000          C          STORE 31 SCATTERED PATTERN VALUES IN THETA,AMP,& PHASE ARRAYS.          03920000
C          DO 760 J=1,JMAX          03533000          C          CHOOSE EITHER LEFT OR RIGHT SIDE OF BEAM. TRANSLATE ANGLES          03930000
C          TO=TT(J)/O.017453293          03540000          C          IF(INDISK .EQ. 0)GC TO 790          03940000
C          SIGNTT=C.0          03550000          C          IF(IJ .LT. (JPEAK-30)) .OR. (J .GT. JPEAK)GO TO 790          03950000
C          IF(COTT(J) .EQ. 0.00000)GO TO 765          03560000          C          L=JPEAK-J+1          03960000
C          SIGNTT=COTT(J)/ABS(COTT(J))          03570000          C          TTOUT(L)=TTPEAK-TO          03970000
C          CONTINUE          03580000          C          TAMP(L)=EYAMP          03980000
C          EVALUATE Y COMPONENT OF SCATTERED FIELD          03590000          C          TPHI(L)=EYPHI          03990000
C          EPLANE=EIT(J,K)*SIPP(K)*SIGNTT + EPP(J,K)*COPP(K)          03600000          C          GC TO 790          03995000
C          A1=REAL(EIT(J,K))          03601000          C          CONTINUE          04000000
C          A2=AIMAG(EIT(J,K))          03602000          C          IF(IJ .LT. (JPEAK+30)) .OR. (J .GT. (JPEAK+30))GO TO 790          04010000
C          A3=REAL(EPP(J,K))          03603000          C          TTOUT(L)=TO-TTPEAK          04020000
C          A4=AIMAG(EPP(J,K))          03610000          C          TAMP(L)=EYAMP          04021000
C          A5=REAL(EPLANE)          03620000          C          TPHI(L)=EYPHI          04021200
C          A6=AIMAG(EPLANE)          03630000          C          CONTINUE          04021300
C          CALL VECTOR(A1,A2,ETAMP,ETPHI)          03640000          C          IF(IJ .LT. JPEAK) .OR. (J .GT. (JPEAK+30))GO TO 790          04021400
C          CALL VECTOR(A3,A4,EPAMP,EPPHI)          03650000          C          L=J-JPEAK+1          04021500
C          CALL VECTOR(A5,A6,EYAMP,EYPHI)          03660000          C          TTOUT(L)=TO-TTPEAK          04021600
C          TRANSLATE PHASE CENTER OF SCATTERED PATTERN          03670000          C          TAMP(L)=EYAMP          04021700
C          03680000          C          TPHI(L)=EYPHI          04021800
C          03690000          C          CONTINUE          04021900
C          03700000          C          PRINT-OUT SCATTERED PATTERN VALUES          04030000
C          03710000          C          IF(FSPLOT .EQ. -1)WRITE(6,2012)TO,ETAMP          04032000
C          03712000          C          E,ETPHI,EPAMP,EPPHI          04040000
C          03713000          C          CALCULATE POWER IN SCATTERED PATTERN          04050000
C          03714000          C          EEJ=ETAMP*ETAMP + EPAMP*EPAMP          04060000
C          03715000          C          IF(IJ .EQ. 1)GC TO 785          04070000
C          03716000          C          04080000
C          03717000          C          04090000
C          03718000          C          04100000
C          03719000          C          04101000

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